

Void vs Dark Energy

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Light propagation
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Fitting the data

SNIa Hubble diagram
WMAP
Putting things together

Other cosmological observations

Conclusions

Can an Inhomogeneous Universe mimic Dark Energy?

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CERN

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¹In collaboration with:

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- In Standard Cosmology we use the Friedmann-Lemaître-Robertson-Walker model.

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- In Standard Cosmology we use the Friedmann-Lemaître-Robertson-Walker model.
- We compute D_L (or D_A) and z

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- We use this to interpret several observations (SNIIa, Hubble constant, CMB, Baryon Acoustic Oscillations,...)

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- We compute D_L (or D_A) and z
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- To fit the observations we need a $p < 0$ term (“Dark Energy”).

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- We use this to interpret several observations (SNIIa, Hubble constant, CMB, Baryon Acoustic Oscillations,...)
- To fit the observations we need a $p < 0$ term (“Dark Energy”).
- **Problem:** We do not understand
 - the amount (why of the same amount as Matter today)?
 - its nature (is it vacuum energy?)

Two main pieces of evidence

- **SNIA** is incompatible with deceleration (independently on other observations)

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- **SN^{Ia}** is incompatible with deceleration (independently on other observations)
 - Assuming them as standard candles.

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- **CMB**: best-fit with power-law (k^{n_s}) primordial spectrum has $\Omega_\Lambda \sim 0.7$.

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But good fit² also with $\Omega_\Lambda = 0$ (*flat*)

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 - low- h (0.45)
 - non-standard primordial spectrum

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- The two dataset,
 - **SNIA**
 - **CMB together with measured h** : $0.55 \lesssim h \lesssim 0.8$are strong evidence for $\Omega_\Lambda \sim 0.7$.

²Blanchard et al. '03, Sarkar and Hunt '04, '07

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- The two dataset,
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 - **CMB together with measured h** : $0.55 \lesssim h \lesssim 0.8$are strong evidence for $\Omega_\Lambda \sim 0.7$.
- Other observations (**BAO** and LSS...) fit consistently

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Is there any alternative?

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- Look for some interesting critical point of view and other logical possibilities

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- Look for some interesting critical point of view and other logical possibilities
- What happens to observations when we have departure from a *homogeneous* model?

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- Look for some interesting critical point of view and other logical possibilities
- What happens to observations when we have departure from a *homogeneous* model?
- Can we accomodate for *all* this evidence if we relax (to some degree) homogeneity?

Homogenous Universe: a good approximation?

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- At $z \gg 1$ (CMB epoch, for example) tiny density fluctuations on all observed scales.

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- ..at late times $\delta \equiv \frac{\delta\rho}{\rho} > 1$ for all scales
 $L \lesssim \mathcal{O}(10)/h \text{ Mpc}$ (1% of Hubble radius)

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- ..at late times $\delta \equiv \frac{\delta\rho}{\rho} > 1$ for all scales
 $L \lesssim \mathcal{O}(10)/h \text{Mpc}$ (1% of Hubble radius)
- Superclusters upto few hundreds of Mpc (10% of Hubble radius), nonlinear objects ("cosmic web")

SDSS data ("The cosmic web")

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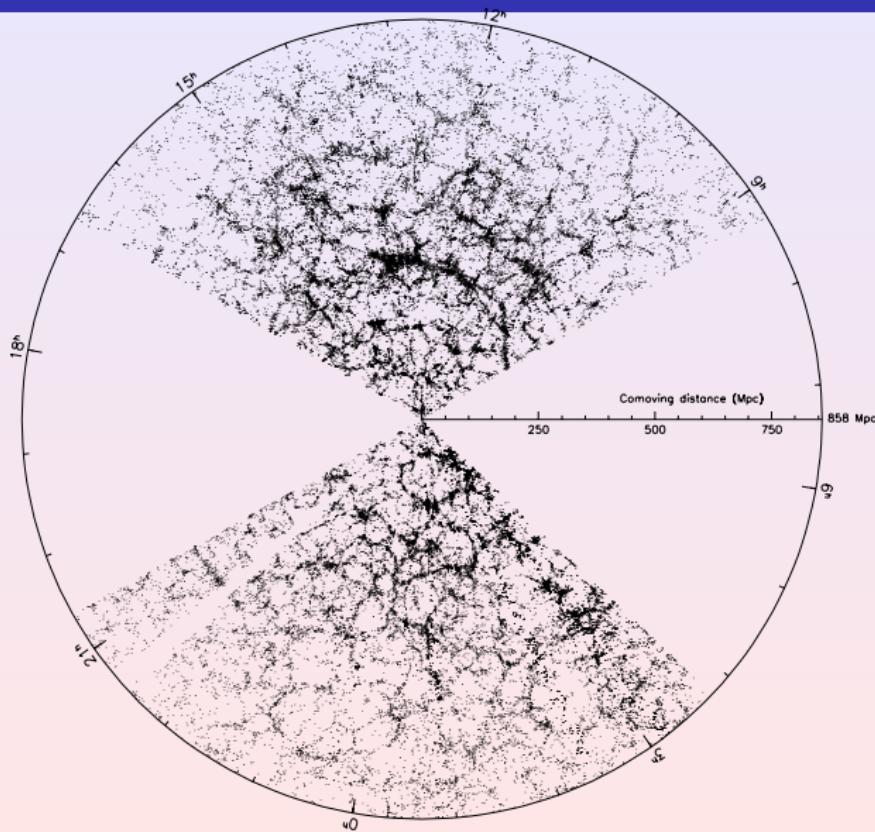
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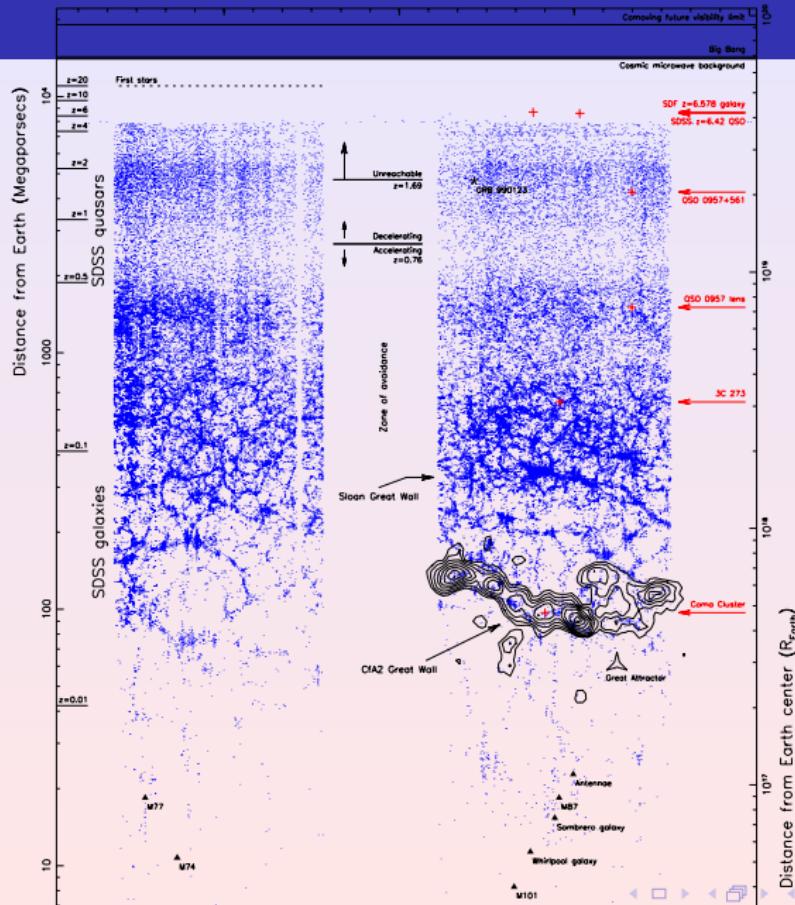


SDSS data

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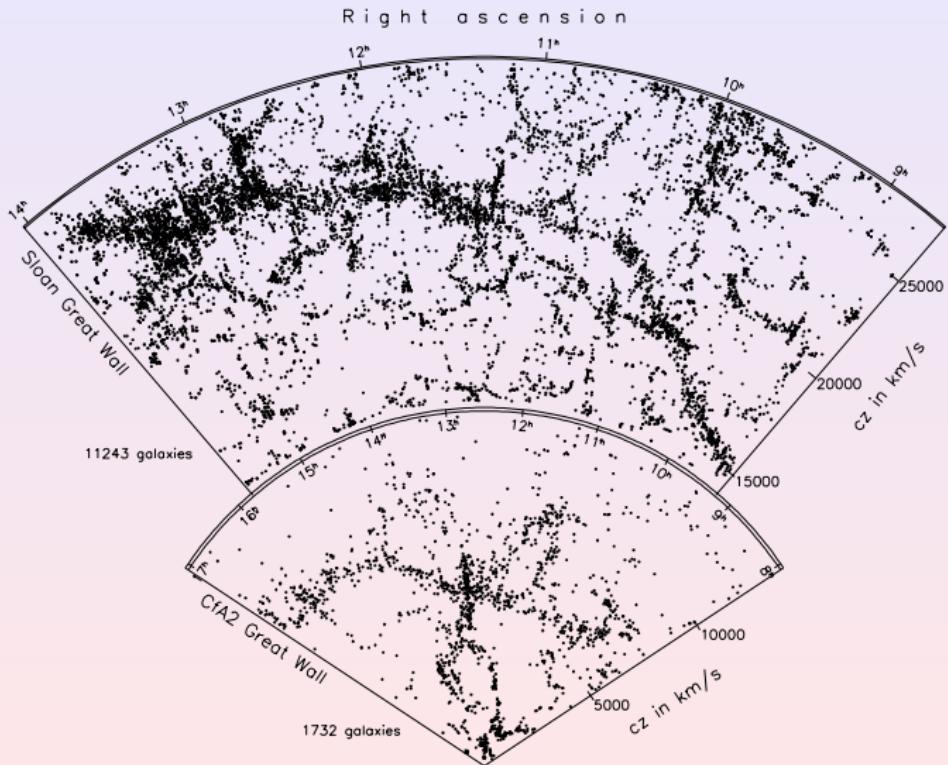
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Three physical effects of inhomogeneities

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Light meets voids and structures. Do they compensate?

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What if we live in a local void?

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- Averaging of Einstein's Equations

Nonlinearity \Rightarrow extra terms in Friedmann equations³

³G.Ellis-W.Stoeger '67, Futamase, Sakai et al, Buchert '95, S. Rasanen '03 ,

E.Kolb-S.Matarrese-A.Riotto-A. N. '04 , A. N. '06

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- Averaging of Einstein's Equations
Nonlinearity \Rightarrow extra terms in Friedmann equations³
- Consider a comoving **inhomogeneous** metric ($p=0$)

$$ds^2 = -dt^2 + a^2(t)dx^i dx^j h_{ij}(t, \mathbf{x})$$

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$$ds^2 = -dt^2 + a^2(t)dx^i dx^j h_{ij}(t, \mathbf{x})$$

- For a comoving domain D :

$$V_D = \int_D \sqrt{h} d^3x, \quad a_D(t) \equiv \left(\frac{V_D}{V_{D_0}} \right)^{1/3};$$

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- For a comoving domain D :

$$V_D = \int_D \sqrt{h} d^3x, \quad a_D(t) \equiv \left(\frac{V_D}{V_{D_0}} \right)^{1/3};$$

$$\frac{\ddot{a}_D}{a_D} = -\frac{4\pi G}{3} (\rho_{\text{eff}} + 3P_{\text{eff}}),$$

$$\left(\frac{\dot{a}_D}{a_D} \right)^2 = \frac{8\pi G}{3} \rho_{\text{eff}},$$

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The extra terms

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• Where

$$\begin{aligned}\rho_{\text{eff}} &= \langle \rho \rangle_D - \frac{Q_D}{16\pi G} - \frac{\langle R \rangle_D}{16\pi G} \\ P_{\text{eff}} &= -\frac{Q_D}{16\pi G} + \frac{\langle R \rangle_D}{48\pi G},\end{aligned}$$

The extra terms

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- The *real* question: **how large is it?**

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- On a large Domain the dominant term has the form⁴ :

$$\frac{\overline{H_D - H}}{H} = \frac{25}{54} \frac{1}{a^2 H^2} \overline{\langle \varphi \nabla^2 \varphi \rangle}$$

⁴

L. Hui-U. Seljak '95, S. Rasanen '03, E. W. Kolb-S. Matarrese-A.N.-A. Riotto '04

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$$= A^2 \frac{a}{a_0} \left(\frac{h \Gamma \text{Mpc}^{-1}}{H_0} \right)^2 \int_0^\infty dq q T^2(q)$$

where $A \sim 10^{-5}$, $\Gamma = \Omega_M h e^{-\Omega_B - \sqrt{2} h \Omega_B / \Omega_M}$.

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where $A \sim 10^{-5}$, $\Gamma = \Omega_M h e^{-\Omega_B - \sqrt{2} h \Omega_B / \Omega_M}$.

- Largest contribution from $\mathcal{O}(10 - 50) \text{Mpc}/h$

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$$\frac{\overline{H_D - H}}{H} \simeq 10^{-5}$$

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- Largest contribution from $\mathcal{O}(10 - 50) \text{Mpc}/h$

$$\frac{\overline{H_D - H}}{H} \simeq 10^{-5}$$

- Small, but not $10^{-10}!$

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- Largest contribution from $\mathcal{O}(10 - 50) \text{Mpc}/h$

$$\frac{\overline{H_D - H}}{H} \simeq 10^{-5}$$

- Small, but not 10^{-10} ! Enhanced by $\left(\frac{k_{EQ}}{H_0} \right)^2$

⁴

L. Hui-U. Seljak '95, S. Rasanen '03, E. W. Kolb-S. Matarrese-A.N.-A. Riotto '04

Power counting

- What about higher (n^{th}) orders ⁵?

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- What about higher (n^{th}) orders ⁵?
- They go as

$$\overline{\langle \varphi (\nabla^2 \varphi)^{n-1} \rangle}$$

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- What about higher (n^{th}) orders ⁵?
- They go as

$$\overline{\langle \varphi (\nabla^2 \varphi)^{n-1} \rangle}$$

- We can write the n^{th} order as

$$10^{-5} \epsilon^{n-1}$$

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- What about higher (n^{th}) orders ⁵?
- They go as

$$\overline{\langle \varphi (\nabla^2 \varphi)^{n-1} \rangle}$$

- We can write the n^{th} order as

$$10^{-5} \epsilon^{n-1}$$

where roughly

$$\epsilon \equiv \frac{A}{1+z} \left(\frac{h \Gamma \text{Mpc}^{-1}}{H_0} \right)^2 \times \text{Int}$$

with

$$\text{Int} = \int dq T^2(q) \approx 0.02$$

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- $\epsilon = \mathcal{O}(1)$ today

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- $\epsilon = \mathcal{O}(1)$ today
- ...each term in the series is of $\mathcal{O}(10^{-5})$!

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- $\epsilon = \mathcal{O}(1)$ today
- ...each term in the series is of $\mathcal{O}(10^{-5})$!
- Do they **sum** up to 10^{-5} or more??

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- Need non-perturbative treatment

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- $\epsilon = \mathcal{O}(1)$ today
- ...each term in the series is of $\mathcal{O}(10^{-5})$!
- Do they **sum** up to 10^{-5} or more??
- Need non-perturbative treatment
- Note: $\epsilon \ll 1$ at high z

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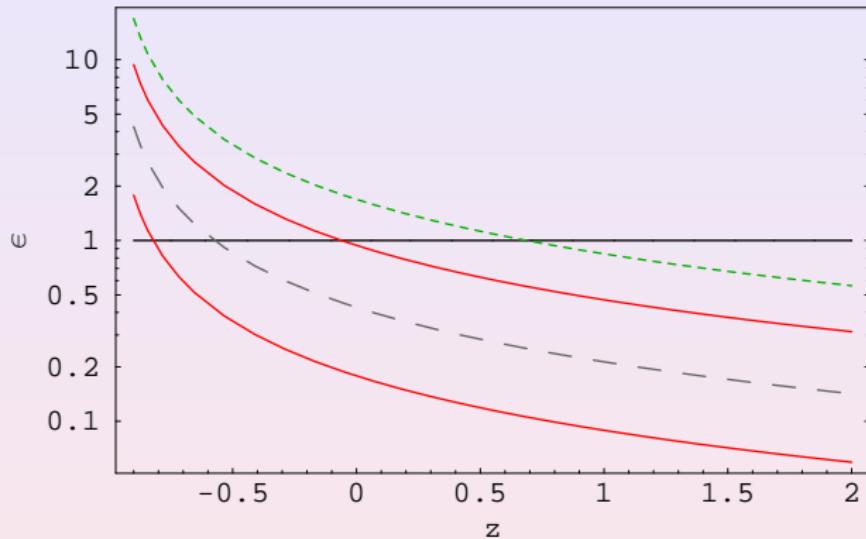


Figure: Grey dashed line: central value,
Red solid lines: 2σ ranges
(We used the growth factor as in matter domination. For comparison, green dotted line: $\Omega_M = 1$).

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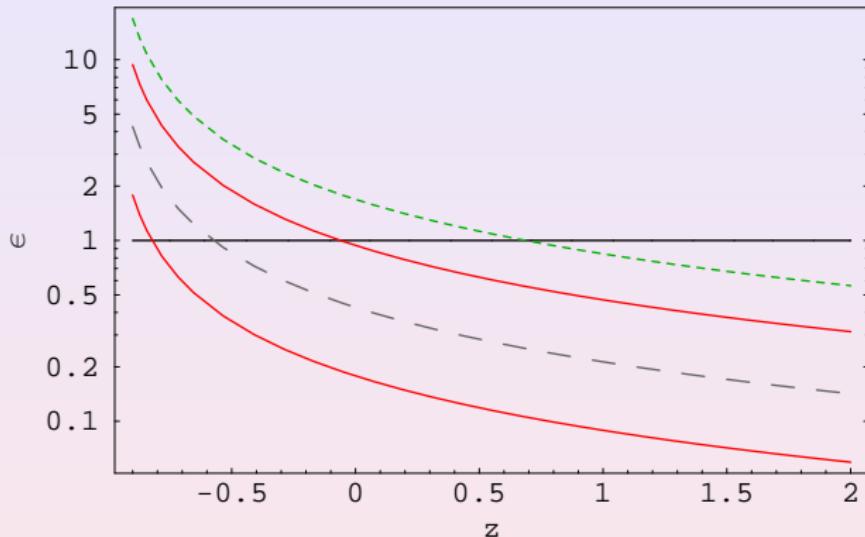


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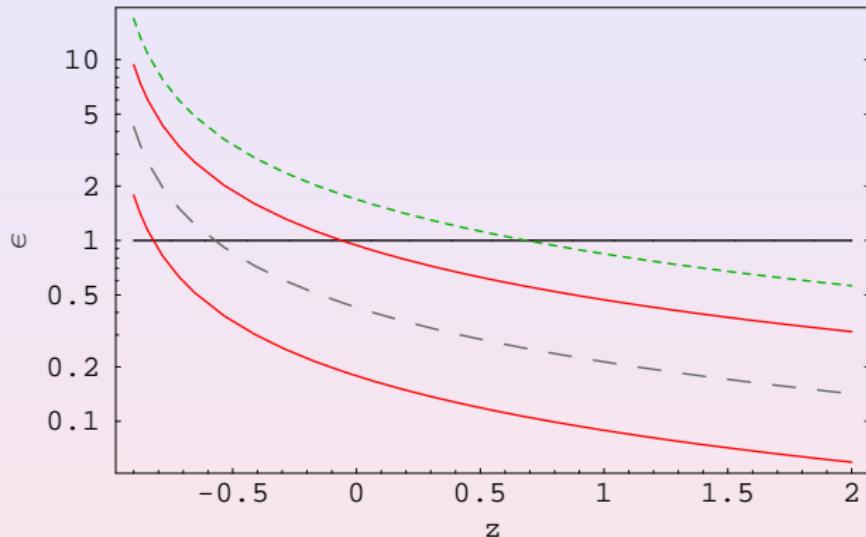


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- Even in absence of average effect on $H(z)$: corrections to photon trajectories

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- Even in absence of average effect on $H(z)$: corrections to photon trajectories
- In fact, actually we measure distances D and redshifts z

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- Even in absence of average effect on $H(z)$: corrections to photon trajectories
- In fact, actually we measure distances D and redshifts z
- All information from expansion comes from plots $D - z$

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- In fact, actually we measure distances D and redshifts z
- All information from expansion comes from plots $D - z$
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- Even in absence of average effect on $H(z)$: corrections to photon trajectories
- In fact, actually we measure distances D and redshifts z
- All information from expansion comes from plots $D - z$
- Cannot disentangle this from backreaction
- Compute $\frac{\Delta z}{1+z}$ and $\frac{\Delta D}{D}$ in the presence of structures

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- Consider Lemaître-Tolman-Bondi exact solutions of E.E. (with $p = 0$) which is

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- Consider Lemaître-Tolman-Bondi exact solutions of E.E. (with $p = 0$) which is
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- We consider two configurations:

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 - LTB spheres embedded in FLRW ("Swiss-Cheese")
 - LTB with shells of periodically varying density ("Onion")
- We study null geodesic in this metric

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- Net effect from one hole⁶ : $\frac{\Delta z}{1+z} \approx (L/r_H)^3 f(\delta)$

⁶T. Biswas-A. N. '06-'07

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- Still small (for late acceleration)

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- Still small (for late acceleration)
- Interesting in the CMB, as a Rees-Sciama effect.

⁶T. Biswas-A. N. '06-'07

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- Net effect scales as $\frac{\Delta z}{1+z} \approx (L/r_H)^2 f(\delta)^7$

⁷ Brouzakis-Tetradis-Tzavara '06, Kolb-Matarrese-Riotto '07, T. Biswas-A. N. '07

⁸ S. Weinberg '76, Brouzakis et al. '06-'07

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- Net effect scales as $\frac{\Delta z}{1+z} \approx (L/r_H)^2 f(\delta)$ ⁷
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- Tight packing: $N_{\text{holes}} \mathcal{O}(L/r_H)^3 = \mathcal{O}(L/r_H)$
- Not so small...
- But it should have zero angular average (unlike z)⁸

⁷ Brouzakis-Tetradis-Tzavara '06, Kolb-Matarrese-Riotto '07, T. Biswas-A. N. '07

⁸ S. Weinberg '76, Brouzakis et al. '06-'07

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- Reliable result or limited by the symmetries of the model?
- LTB model swiss-cheese: special case

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- Szekeres swiss-cheese model with asymmetric holes
(Bolejko '08)

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(Bolejko '08) *Effects of similar size*

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- The cheese feels *no backreaction* by construction
- What happens without spherical symmetry?
- Szekeres swiss-cheese model with asymmetric holes
(Bolejko '08) *Effects of similar size*
- But still special: the cheese feels *no backreaction* of the holes

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⁹Tomita '98, Tomita '00, Celerier '01, Wiltshire '05, Moffat '05, Alnes et al. '05, Mansouri et al. '06, Biswas & A.N.'07, Garcia-Bellido and Haugboelle '08, Zibin et al. '08 ...

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- Suppose that we live in a peculiar local region

⁹Tomita '98, Tomita '00, Celerier '01, Wiltshire '05, Moffat '05, Alnes et al. '05, Mansouri et al. '06, Biswas & A.N.'07, Garcia-Bellido and Haugboelle '08, Zibin et al. '08 ...

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- Suppose that we live in a peculiar local region
- \Rightarrow low z observations may be very different from average.

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Conclusions

- Suppose that we live in a peculiar local region
- \Rightarrow low z observations may be very different from average.
- One realizes that acceleration is inferred **comparing low z with high z ...**

⁹Tomita '98, Tomita '00, Celerier '01, Wiltshire '05, Moffat '05, Alnes et al. '05, Mansouri et al. '06, Biswas & A.N.'07, Garcia-Bellido and Haugboelle '08, Zibin et al. '08 ...

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- Suppose that we live in a peculiar local region
- \Rightarrow low z observations may be very different from average.
- One realizes that acceleration is inferred **comparing low z with high z ...**
- Can this mimic acceleration ⁹?

⁹Tomita '98, Tomita '00, Celerier '01, Wiltshire '05, Moffat '05, Alnes et al. '05, Mansouri et al. '06, Biswas & A.N.'07, Garcia-Bellido and Haugboelle '08, Zibin et al. '08 ...

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- Mustapha, Hellaby, Ellis '97: show that LTB can reproduce any observations

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- Mustapha, Hellaby, Ellis '97: show that LTB can reproduce any observations
- Celerier '99: showed that LTB can mimic Λ CDM

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- Mustapha, Hellaby, Ellis '97: show that LTB can reproduce any observations
- Celerier '99: showed that LTB can mimic Λ CDM
- Tomita '01: Compensated Void 200 – 300 Mpc/ h scale

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- Consider a “compensated Void” : a spherical Void plus an external shell of matter

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- Consider a “compensated Void” : a spherical Void plus an external shell of matter
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- Assumption: we live near the center
- A void expands faster than the “external” FLRW

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- So, nearby objects inside the void redshift more

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- This can mimic acceleration (as we will see...)

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- A void expands faster than the “external” FLRW
- So, nearby objects inside the void redshift more
- This can mimic acceleration (as we will see...)
- How much contrast δ and how large L is needed?

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- Before going to the quantitative analysis...
- Let's review some literature and observations on Voids

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- Inoue and Silk '06: some features of the low multipole **anomalies** in the CMB data could be explained by a pair of huge Voids ($L \sim 200 \text{ Mpc}/h$, $\delta \sim -0.3$)

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- The CMB has a **Cold Spot** (M. Cruz et al. ('06 and '07)): it could be explained by another similar Big Void (Inoue and Silk '06)

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- The Cold Spot in the CMB claimed to be correlated with an underdense region in the LSS (Rudnick, Brown and Williams '07, but see Huterer and ...)

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- The Cold Spot in the CMB claimed to be correlated with an underdense region in the LSS (Rudnick, Brown and Williams '07, but see Huterer and ...)
- It could be detected via **lensing** (S. Das and D. Spergel '08) and via **non-gaussian coupling Rees-Sciama effect - lensing** (I. Masina and A.N., in preparation)

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- Some observational evidence for a local large underdense region ($\sim 25\%$ less dense, $r \sim 200 \text{ Mpc}/h$) from number counts of galaxies (2MASS)
(Frith et al. Mon. Not. Roy. Astron. Soc. **345**, 1049 (2003))
- It would represent a 4σ fluctuation, at odds with Λ CDM.

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(Frith et al. Mon. Not. Roy. Astron. Soc. **345**, 1049 (2003))
- It would represent a 4σ fluctuation, at odds with Λ CDM.
- Many Large Voids identified via **ISW** effect in the SDSS LRG catalog (about $100 \text{ Mpc}/h$ radius) (Granett et al. '08)
- Also in contradiction with Λ CDM: $P < 10^{-8}$ (Sarkar & Hunt '08)

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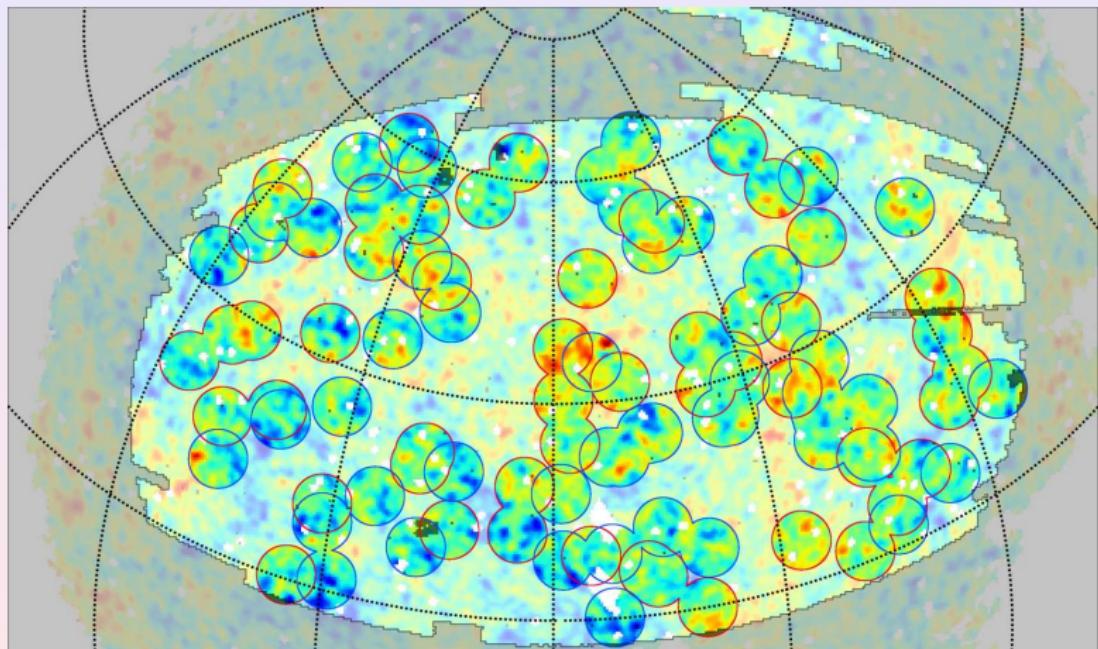


Figure: Granett, Neyrinck & Szapudi '08

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- Recent measurement (Kashlinsky et al.'08): **very large coherent motion** on $300\text{Mpc}/h$ scale, inconsistent with ΛCDM
- Could be due to very large scale inhomogeneous matter distribution

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- Recent measurement (Kashlinsky et al.'08): **very large coherent motion** on $300\text{Mpc}/h$ scale, inconsistent with ΛCDM
- Could be due to very large scale inhomogeneous matter distribution
- Watkins, Feldman & Hudson '08: use peculiar velocities of various (4500) objects in a $100\text{Mpc}/h$ radius.
Find 400km/sec (expected 100km/sec)

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- It will turn out that a Minimal Void needs at least the same size (*for Riess '07 SNIa and WMAP*)

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- What is the size we need to mimick DE?
- It will turn out that a Minimal Void needs at least the same size (*for Riess '07 SNIa and WMAP*)
- $r_{\text{Void}} \sim 200 - 250 \text{ Mpc}/h$ and $\delta \sim -0.4$

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 - Percolation of Voids?

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- Can we ever get huge Voids?
 - Percolation of Voids?
 - Non-standard structure formation?
 - Non-gaussianity?
 - Nucleation of primordial Bubbles

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- **Problem (II):** even larger size seems necessary for other observations
- For **UNION** (...) data **$500\text{Mpc}/h$** required

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- **Problem (II):** even larger size seems necessary for other observations
- For **UNION** (...) data **$500\text{Mpc}/h$** required and the fit is not very good: $\chi^2 = 343$ (306 d.o.f.)

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- **Problem (II):** even larger size seems necessary for other observations
- For **UNION** (...) data **$500\text{Mpc}/h$** required and the fit is not very good: $\chi^2 = 343$ (306 d.o.f.)
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- A very good fit obtained with **(2 Gpc)** (**Garcia-Bellido & Haugboelle**)
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- **Problem (II):** even larger size seems necessary for other observations
- For **UNION** (...) data **$500\text{Mpc}/h$** required and the fit is not very good: $\chi^2 = 343$ (306 d.o.f.)
- A very good fit obtained with **(2 Gpc)** (Garcia-Bellido & Haugboelle)
- Or adding open curvature
- Need for a Large Void also for **BAO**: a better fit is obtained with at least (1 Gpc) (Garcia-Bellido & Haugboelle).

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$$ds^2 = -dt^2 + \frac{R'^2(r, t)}{1 + 2r^2k(r)}dr^2 + R^2(r, t)(d\theta^2 + \sin^2\theta d\varphi^2)$$

with comoving coordinates (r, θ, φ) and proper time t .

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- Spherically symmetric.

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with comoving coordinates (r, θ, φ) and proper time t .

- Spherically symmetric.
- Einstein equations:

$$\frac{1}{2} \frac{\dot{R}^2(r, t)}{R^2(r, t)} - \frac{GM(r)}{R^3(r, t)} = \frac{r^2k(r)}{R^2(r, t)},$$
$$4\pi\rho(r, t) = \frac{M'(r)}{R'(r, t)R^2(r, t)},$$

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$$ds^2 = -dt^2 + \frac{R'^2(r, t)}{1 + 2r^2k(r)}dr^2 + R^2(r, t)(d\theta^2 + \sin^2\theta d\varphi^2)$$

It has the solutions:

- For $k(r) > 0$ ($k(r) < 0$),

$$R = \frac{GM(r)}{2r^2|k(r)|}[\cos h(u) - 1], \quad (4.1)$$

$$t - t_b(r) = \frac{GM(r)}{[2r^2|k(r)|]^{3/2}}[\sin h(u) - u].$$

- $k(r) = 0$,

$$R(r, t) = \left[\frac{9GM(r)}{2} \right]^{1/3} [t - t_b(r)]^{\frac{2}{3}}.$$

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- $t_b(r) = 0$ for our purposes, and “Gauge” choice:
 $M(r) \propto r^3$

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- $k(r)$ contains all the physical information about the profile.

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- $t_b(r) = 0$ for our purposes, and “Gauge” choice:
 $M(r) \propto r^3$
- $k(r)$ contains all the physical information about the profile.
- $k = 0$ flat FLRW, $k = \pm 1$ open/closed FLRW.

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- $t_b(r) = 0$ for our purposes, and “Gauge” choice:
 $M(r) \propto r^3$
- $k(r)$ contains all the physical information about the profile.
- $k = 0$ flat FLRW, $k = \pm 1$ open/closed FLRW.
- The idea is to describe structure formation
(start with $\delta(r, t_l) \ll 1$ and end up with $\delta(r, t_{\text{now}}) \gg 1$)

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- $t_b(r) = 0$ for our purposes, and “Gauge” choice:
 $M(r) \propto r^3$
- $k(r)$ contains all the physical information about the profile.
- $k = 0$ flat FLRW, $k = \pm 1$ open/closed FLRW.
- The idea is to describe structure formation
(start with $\delta(r, t_l) \ll 1$ and end up with $\delta(r, t_{\text{now}}) \gg 1$)
- We play with $k(r)$ to describe $\delta(r, t_l)$.

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- Matching of an LTB sphere (of radius L) to FLRW:

$$k'(0) = k'(L) = 0 ,$$

$$k(L) = \frac{4\pi}{3}\Omega_k, \quad \text{for } |\Omega_k| \ll 1 ,$$

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$$k(L) = \frac{4\pi}{3}\Omega_k, \quad \text{for } |\Omega_k| \ll 1 ,$$

We use:



$$k(r) = k_{max} \left[\left(\frac{r}{L} \right)^4 - 1 \right]^2 \quad (\text{for } r < L)$$

$$k(r) = 0 \text{ (flat)} \quad (\text{for } r > L)$$

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- Two parameters, L and k_{\max} .

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- Roughly:

$$\rho(r, t) \simeq \frac{\langle \rho \rangle(t)}{1 + (t/t_0)^{2/3} \epsilon(r)},$$

$$\text{where } \langle \rho \rangle(t) \equiv \frac{M_p^2}{6\pi t^2}, \quad \text{and } \epsilon(r) \equiv 3k(r) + rk'(r).$$

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- $\epsilon \ll 1$ linear growth

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- $\epsilon \ll 1$ linear growth $\propto a(t) \propto t^{2/3}$

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- $\epsilon \ll 1$ linear growth $\propto a(t) \propto t^{2/3}$
- ϵ not small: δ grows rapidly (as in Zel'dovich approx)

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$$\text{where } \langle \rho \rangle(t) \equiv \frac{M_p^2}{6\pi t^2}, \quad \text{and } \epsilon(r) \equiv 3k(r) + rk'(r).$$

- $\epsilon \ll 1$ linear growth $\propto a(t) \propto t^{2/3}$
- ϵ not small: δ grows rapidly (as in Zel'dovich approx)
- We work at most with $\delta \sim \mathcal{O}(1)$.

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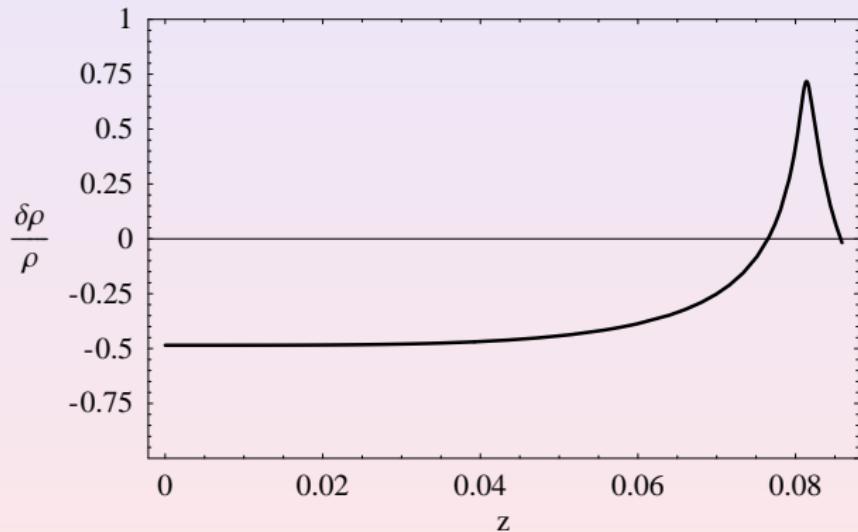
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- Solve for $t(r)$ along a $ds^2 = 0$ trajectory
- Then solve for

$$\frac{dz}{dr} = \frac{(1 + z(r))\dot{R}'(r, t(r))}{\sqrt{1 + 2r^2 k(r)}}.$$

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- Then solve for

$$\frac{dz}{dr} = \frac{(1 + z(r))\dot{R}'(r, t(r))}{\sqrt{1 + 2r^2 k(r)}}.$$

- The result $z(r)$ can be found numerically

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Conclusions

- Solve for $t(r)$ along a $ds^2 = 0$ trajectory
- Then solve for

$$\frac{dz}{dr} = \frac{(1 + z(r))\dot{R}'(r, t(r))}{\sqrt{1 + 2r^2 k(r)}}.$$

- The result $z(r)$ can be found numerically
- We also have some very good analytical approximations

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- Always in GR, luminosity distance and angular distance:

$$D_L = D_A(1 + z)^2.$$

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- Always in GR, luminosity distance and angular distance:

$$D_L = D_A(1 + z)^2.$$

-

$$D_A^2 \equiv \frac{dA}{d\Omega} = \frac{d\theta_S d\phi_S \sqrt{g_{\theta\theta} g_{\phi\phi}}}{d\bar{\theta}_O d\bar{\phi}_O}$$

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- If observer in the center:

$$D_A^2 = R^2|_S.$$

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- Always in GR, luminosity distance and angular distance:

$$D_L = D_A(1+z)^2.$$

-

$$D_A^2 \equiv \frac{dA}{d\Omega} = \frac{d\theta_S d\phi_S \sqrt{g_{\theta\theta} g_{\phi\phi}}}{d\bar{\theta}_O d\bar{\phi}_O} = \frac{d\theta_S d\phi_S}{d\bar{\theta}_O d\bar{\phi}_O} R^2|_S,$$

- If observer in the center:

$$D_A^2 = R^2|_S.$$

- For generic observer (but radial trajectory):

$$D_A = R_S \left(R_O \int_{r_O}^{r_S} \frac{R'(r, t(r))}{(1+2E(r))(1+z(r))R(r, t(r))^2} dr \right),$$

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$$f \equiv \frac{\sqrt[3]{2}(\cosh(u) - 1)}{3^{2/3}(\sinh(u) - u)^{2/3}} - 1 \quad (4.2)$$

$$u_0 = 6^{1/3}(\sinh(u) - u)^{1/3}. \quad (4.3)$$

Then, one can use this function in the following equations:

$$\tau(r) = \tau_0 - \frac{\pi}{9}\gamma^2 \bar{M}r[1 + f(\gamma^2 \tau_0^2 k(r))], \quad (4.4)$$

$$1 + z(r) = \left(\frac{\tau_0}{\tau(r)}\right)^2 \exp\left[\frac{4\pi\gamma^2 \bar{M}r}{9}f(\gamma^2 \tau_0^2 k(r))\right] \quad (4.5)$$

$$D_L(r) = \frac{\pi}{3}\gamma^2 r \tau(r)^2 [1 + f(\gamma^2 \tau_0^2 k(r))] [1 + z(r)]^2 \quad (4.6)$$

$$\tau_0 = \left(\frac{2\bar{M}}{3H_0}\right)^{1/3} \quad (4.7)$$

$$\gamma = \left(\frac{9\sqrt{2}}{\pi}\right)^{1/3} \quad (4.8)$$

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- Evidence for acceleration comes from mismatch between:
 - measurements at low redshift ($0.03 \lesssim z \lesssim 0.08$)
 - high- z SN (roughly $0.4 \lesssim z \lesssim 1$)

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 - measurements at low redshift ($0.03 \lesssim z \lesssim 0.08$)
 - high- z SN (roughly $0.4 \lesssim z \lesssim 1$)
- We choose large r_{Void} (at $z \approx 0.08 - 0.09$)

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 - high- z SN (roughly $0.4 \lesssim z \lesssim 1$)
- We choose large r_{Void} (at $z \approx 0.08 - 0.09$)
- \Rightarrow The Local Bubble is different from the average.

Outside just matter dominated (even if there are other Bubbles, their effect is small)

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- At high z ($z \gtrsim 0.1$), just matter dominated Universe:

$$D_{\text{EdS}}(z) \approx \frac{2}{H_{\text{out}}}(1 + z - \sqrt{1 + z}).$$

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- At high z ($z \gtrsim 0.1$), just matter dominated Universe:

$$D_{\text{EdS}}(z) \approx \frac{2}{H_{\text{out}}}(1 + z - \sqrt{1 + z}).$$

- At low z "open-like" Universe with a different H

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$$D_{\text{EdS}}(z) \approx \frac{2}{H_{\text{out}}}(1 + z - \sqrt{1 + z}).$$

- At low z "open-like" Universe with a different H
- Two (reduced) Hubble parameters: h and h_{out}

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- Rapid transition near the shell-like structure

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- At high z ($z \gtrsim 0.1$), just matter dominated Universe:

$$D_{\text{EdS}}(z) \approx \frac{2}{H_{\text{out}}}(1 + z - \sqrt{1 + z}).$$

- At low z "open-like" Universe with a different H
- Two (reduced) Hubble parameters: h and h_{out}
- Rapid transition near the shell-like structure
- h corresponds to what is measured

Δm for different models

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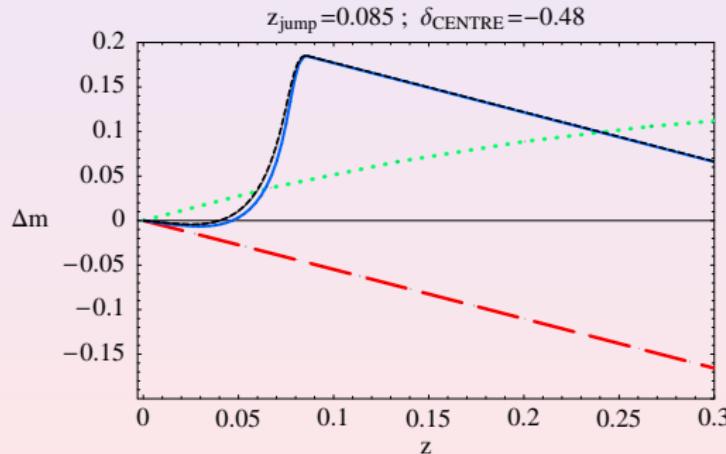
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- Magnitude is $m \equiv 5 \log_{10} D(z)$
- The open “empty” Universe is subtracted ($\Omega_K = -1$)



$m - z$ diagram: Riess data

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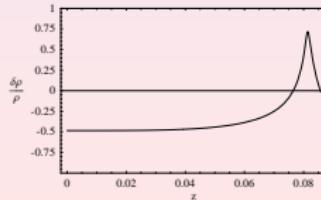
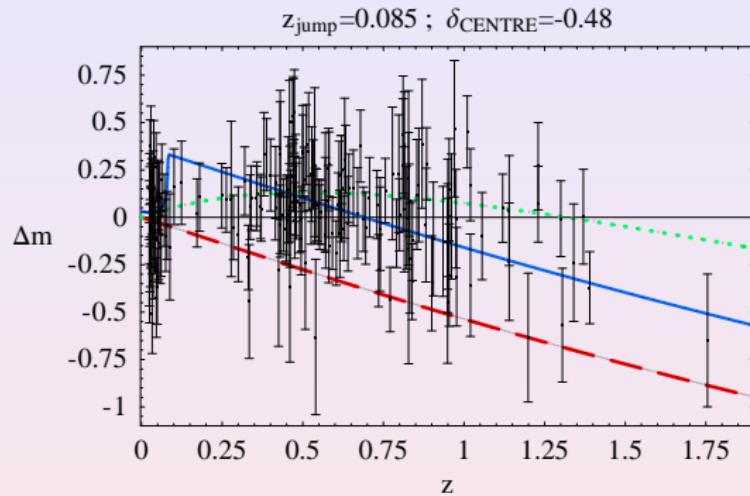
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- We fix several values of L

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- We fix several values of L
- What matters is just the Jump: $\mathcal{J} \equiv \frac{h}{h_{OUT}}$

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- We fix several values of L
- What matters is just the Jump: $\mathcal{J} \equiv \frac{h}{h_{OUT}}$
- This is also related to the central density contrast:
$$\mathcal{J} = 2 - (1 - \delta_0)^{1/3}$$

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- What matters is just the Jump: $\mathcal{J} \equiv \frac{h}{h_{OUT}}$
- This is also related to the central density contrast:
$$\mathcal{J} = 2 - (1 - \delta_0)^{1/3}$$
- We vary \mathcal{J} and compute the χ^2 .

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Riess et al. dataset, astro-ph/0611576 (182 SNIa)

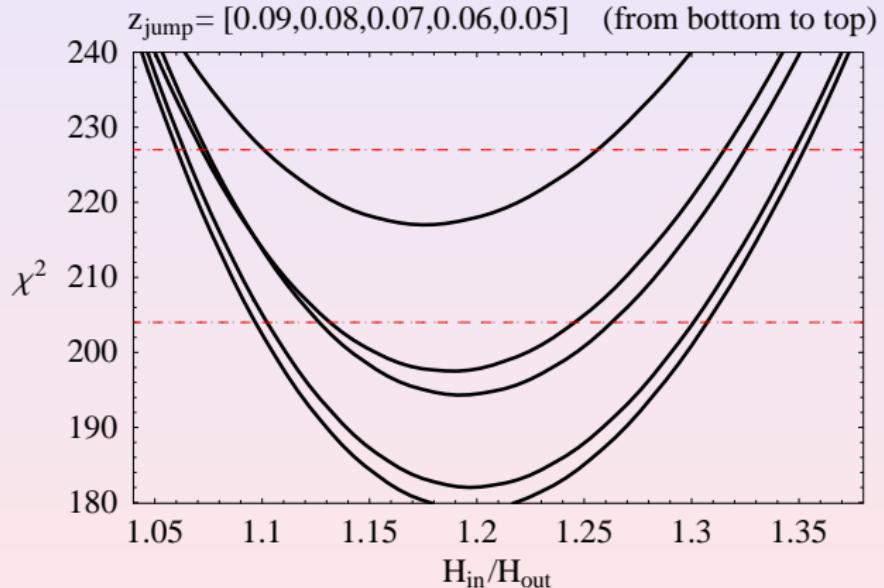


Figure: The red dashed lines are 10% and 1% goodness-of-fit (182 data points)

LTB Void fit

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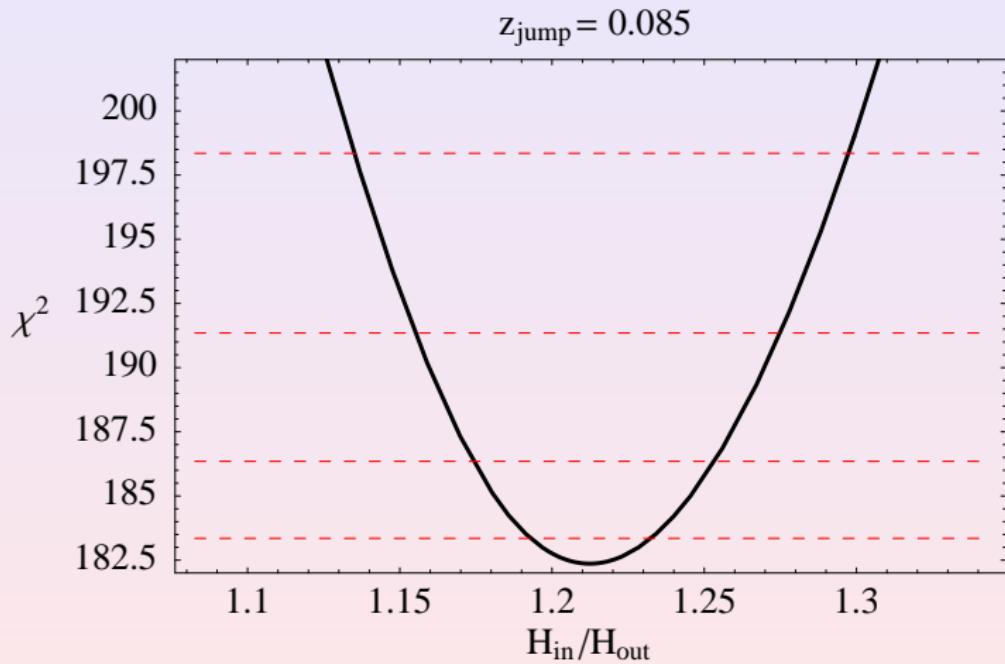


Figure: Here we use the full LTB model. We show 1σ , 2σ , 3σ and 4σ intervals (using likelihood $\propto e^{-\chi^2/2}$).

χ^2 : Riess data

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Table: Comparison with data (full data set of Riess et al.)

Model	χ^2 (181 d.o.f.)
Λ CDM (with $\Omega_M = 0.27, \Omega_\Lambda = 0.73$)	160
EdS (with $\Omega_M = 1, \Omega_\Lambda = 0$)	274
Void ($\sqrt{\langle \delta^2 \rangle} \approx 0.4$ on $L = 250/h\text{Mpc}$)	182

Remarks:

- With instrumental error only: no smooth curve can give a good fit
- Estimated error from intrinsic variability added in quadrature

χ^2 : Riess data

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Remarks:

- With instrumental error only: no smooth curve can give a good fit
- Estimated error from intrinsic variability added in quadrature
- Not as good as Λ CDM
- Becomes better including curvature Ω_k outside

χ^2 : UNION data

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Table: Comparison with data (Union data)

Model	χ^2 (307 d.o.f.)
Λ CDM (with $\Omega_M = 0.27, \Omega_\Lambda = 0.73$)	304
Void ($\sqrt{\langle \delta^2 \rangle} \approx 0.4$ on $L = 500/h\text{Mpc}$)	340
Void ($\sqrt{\langle \delta^2 \rangle} \approx 0.7$ on $L = 2000/h\text{Mpc}$)	304

Remarks:

- It seems necessary to consider a larger Void (Gpc scale)
- Or add curvature? (work in progress)

UNION fit with 2Gpc

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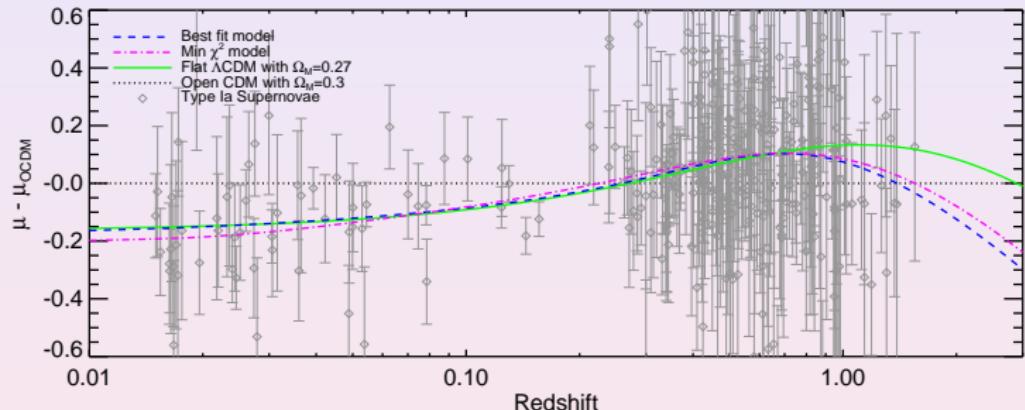


Figure: Taken from Garcia-Bellido & Haugboelle '08
(similar fits also in Zibin et al. '08)

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- We try to fit the WMAP 3-yr data

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- We try to fit the WMAP 3-yr data
- We look at TT and TE correlations, using CosmoMC

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- In principle: we should compute propagation in EdS from $z = 1100$ to $z \sim 0.1$, and then in the Bubble

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- In principle: we should compute propagation in EdS from $z = 1100$ to $z \sim 0.1$, and then in the Bubble
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Small because of compensation

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 - off-center location: dipole

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 - off-center location: dipole and integrated effect (low- l)
 - Non-sphericity (again effect on low- l)
- We do not consider them: just EdS with h_{out} , with some assumptions on the primordial spectrum:
 - n_s plus running α_s
 - Flat spectrum plus bump (as in P. Hunt and S. Sarkar, arXiv:0706.2443 [astro-ph]; A. Blanchard, M. Douspis, M. Rowan-Robinson and S. Sarkar, Astron. Astrophys. **412**, 35 (2003) [arXiv:astro-ph/0304237].)

Priors (Λ CDM)

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The usual prior set is:

- Allow for nonzero Ω_Λ .
- Power-law spectrum with index n_s .
- (eventually with running α_s)
- $P(k) \propto k^{n_s(k_0) + \frac{1}{2} \ln(k/k_0) \alpha_s}$

Priors: without Λ

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A different prior set, that we use:

- **Not** allow for Ω_Λ .
- Power-law spectrum with index n_s .
- with running α_s
- $P(k) \propto k^{n_s(k_0) + \frac{1}{2} \ln(k/k_0) \alpha_s}$

Priors: without Λ

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- **Not** allow for Ω_Λ .
- Power-law spectrum with index n_s .
- with running α_s
- $P(k) \propto k^{n_s(k_0) + \frac{1}{2} \ln(k/k_0) \alpha_s}$
- (we also may allow for some curvature)

Fit to WMAP3

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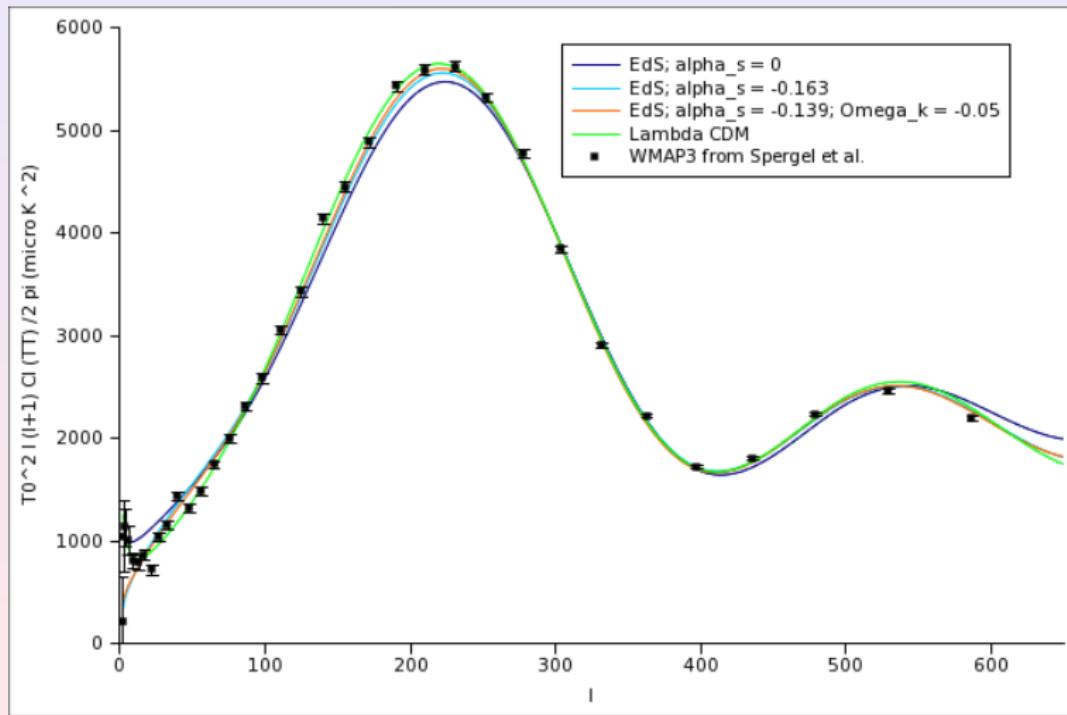
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Model	C_l^{TT}		$C_l^{TT} + C_l^{TE}$		Total	
	χ^2_{eff}	G.F.	χ^2_{eff}	G.F.	χ^2_{eff}	G.F.
Concordant Λ CDM	1038.9	4.7%	1455.2	11.3%	3538.6	41%
EdS $\alpha_S = 0$	1124.6	0%	1711.9	0%	3652.3	6%
EdS $\alpha_S \neq 0$	1057.8	1.9 %	1475.5	5.7%	3577.4	24.6%
EdS $\alpha_S, \Omega_k = -0.050$	1048.7	2.9%	1466	7.9%	3560.9	31.1%

Table:

1st column: high-/ TT ($31 \leq l \leq 1000$)

2nd column: high-/ TT ($31 \leq l \leq 1000$) and TE ($24 \leq l \leq 450$)

3rd column: total of TT ($2 \leq l \leq 1000$) and TE ($2 \leq l \leq 450$)

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The EdS model, with running, has:

- low h_{OUT} (about ~ 0.45)

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The EdS model, with running, has:

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It has to be consistent with the SNIa analysis and the local measurements of h

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The EdS model, with running, has:

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- low n_S (about ~ 0.73)

and large negative α_s (about ~ -0.16)

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Conclusions

The EdS model, with running, has:

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It has to be consistent with the SNIa analysis and the local measurements of h
- low n_S (about ~ 0.73)
and large negative α_s (about ~ -0.16)
- larger value of Ω_M/Ω_b (around 10 instead of 6)

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- low n_S (about ~ 0.73)
and large negative α_s (about ~ -0.16)
- larger value of Ω_M/Ω_b (around 10 instead of 6)
- $\Omega_b h_{\text{out}}^2$ ($\sim 0.018^{+0.001}_{-0.002}$) consistent with BBN constraint
(which is $0.017 \leq \Omega_b h_{\text{out}}^2 \leq 0.024$, at 95% C.L.)

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	Λ CDM	EdS, $\alpha_s = 0$	EdS, $\alpha_s \neq 0$	EdS, $\alpha_s, \Omega_k \neq 0$
$\Omega_b h_{\text{out}}^2$	$0.022^{+0.002}_{-0.002}$	$0.022^{+0.001}_{-0.001}$	$0.018^{+0.001}_{-0.002}$	$0.019^{+0.002}_{-0.001}$
$\Omega_m h_{\text{out}}^2$	$0.106^{+0.021}_{-0.013}$	$0.198^{+0.008}_{-0.011}$	$0.186^{+0.011}_{-0.009}$	$0.167^{+0.009}_{-0.007}$
Ω_Λ	$0.759^{+0.041}_{-0.103}$	0	0	0
z_{re}	$11.734^{+4.993}_{-7.619}$	$8.697^{+4.351}_{-6.694}$	$13.754^{+2.246}_{-5.752}$	$13.342^{+2.55}_{-5.011}$
Ω_k	0	0	0	-0.05
n_s	$0.96^{+0.04}_{-0.04}$	$0.94^{+0.021}_{-0.038}$	$0.732^{+0.07}_{-0.071}$	$0.761^{+0.069}_{-0.069}$
α_s	0	0	$-0.161^{+0.044}_{-0.044}$	$-0.13^{+0.037}_{-0.048}$
$10^{10} A_s$	$20.841^{+3.116}_{-3.442}$	$25.459^{+2.135}_{-2.766}$	$25.302^{+2.182}_{-2.968}$	$23.975^{+2.198}_{-2.448}$
Ω_m/Ω_b	$4.73^{+0.999}_{-0.485}$	$9.119^{+0.341}_{-0.357}$	$10.094^{+0.645}_{-0.489}$	$8.929^{+0.512}_{-0.541}$
h_{out}	$.72857^{+0.05137}_{-0.07393}$	$.46857^{+0.0888}_{-0.01307}$	$.4523^{+0.01291}_{-0.01129}$	$.42069^{+0.01107}_{-0.00919}$
Age/GYr	$13.733^{+0.389}_{-0.369}$	$13.908^{+0.399}_{-0.258}$	$14.408^{+0.369}_{-0.4}$	$15.338^{+0.342}_{-0.393}$
σ_8	$0.77^{+0.121}_{-0.109}$	$1.012^{+0.056}_{-0.081}$	$0.919^{+0.07}_{-0.075}$	$0.862^{+0.06}_{-0.063}$
τ	$0.095^{+0.072}_{-0.074}$	$0.047^{+0.037}_{-0.041}$	$0.079^{+0.023}_{-0.044}$	$0.081^{+0.024}_{-0.041}$

Table: Most likely parameter values with 1σ errors for the various COSMOMC Runs

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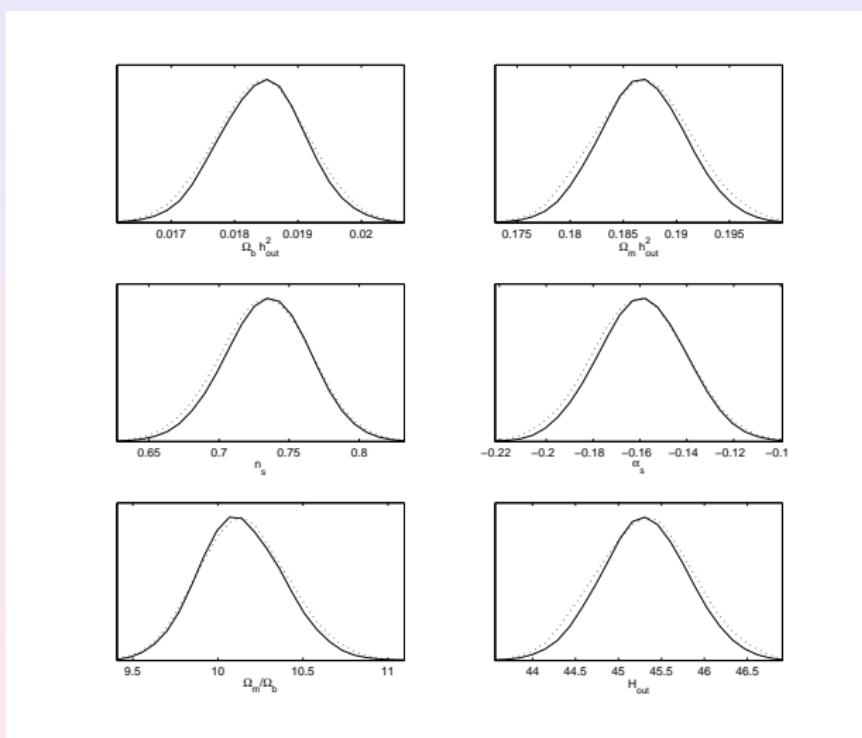


Figure: likelihoods to WMAP 3-yr for the run “EdS with α_s ”

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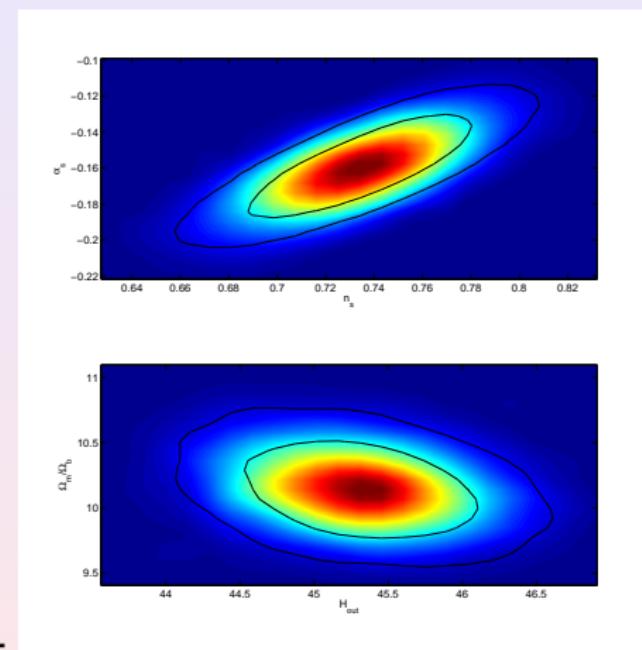


Figure: Contour likelihood plots to WMAP 3-yr for the run “EdS with α_s ”

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● A crucial point: we have

- a low h_{out}
- a constraint on $\mathcal{J} = h/h_{\text{out}}$

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● A crucial point: we have

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● We get a constraint on h

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- A crucial point: we have

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- a constraint on $\mathcal{J} = h/h_{\text{out}}$

- We get a constraint on h . Compatible with local observations?

- $h = 0.72 \pm 0.08$ from HST (W. L. Freedman *et al.*, *Astrophys. J.* **553**, 47 (2001))

- $h = 0.62 \pm 0.01 \pm 0.05$ from HST with corrected Cepheids (A. Sandage *et al.*, *Astrophys. J.* **653**, 843 (2006))

- $h = 0.59 \pm 0.04$ from Supernovae (Parodi, Saha, Sandage and Tammann, arXiv:astro-ph/0004063.)

- $h = 0.54_{-.04}^{-.03}$ SZ effect ($z \approx 1$) (E. D. Reese *et al.* *Astrophys. J.* **581**, 53 (2002))

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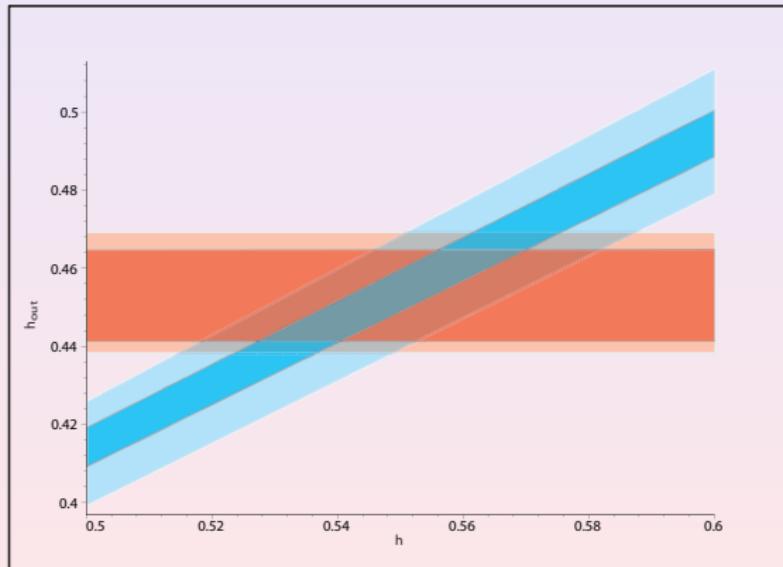


Figure: 1- σ and 2- σ Contour plots for h vs. h_{out} .

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At 95% C.L. we have (for $L \approx 250/h$ Mpc) :

- $1.17 \leq \mathcal{J} \leq 1.25 \Rightarrow 0.42 \leq |\delta_0| \leq 0.58$

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(but note that the average $\sqrt{\langle \delta^2 \rangle}$ is smaller)

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- $0.44 \leq h_{\text{out}} \leq 0.47$

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(but note that the average $\sqrt{\langle \delta^2 \rangle}$ is smaller)

- $0.44 \leq h_{\text{out}} \leq 0.47$

- $0.51 \leq h \leq 0.59$

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- The same happens for the model with a bump in the primordial spectrum (S. Sarkar et al. '03 and '07)
- The bump is at a scale of about $100/h$ Mpc

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- It gives a good fit to WMAP (better than Λ CDM)

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- The original proposal had too low h (~ 0.44)

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⇒ Combine with the Minimal Void scenario

h in the Bump model

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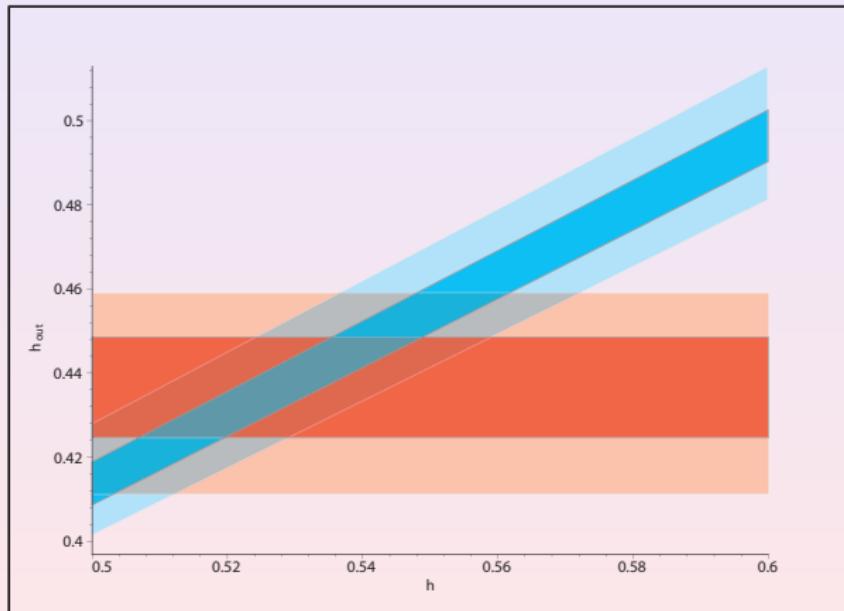


Figure: 1- σ and 2- σ Contour plots for h vs. h_{out} .

Baryon Acoustic Oscillations

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- Measurement of baryon acoustic peak in the galaxy distribution (Eisenstein et al., 2005).
- The position of the peak measures the ratio of the sound horizon at recombination **vs.** angular distance at $z = 0.35$

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- But it also depends on the spectral index n_s :

$$D_V = 1370 \pm 64 \text{ and } \Omega_m h^2 = 0.130 (n_s/0.98)^{-1.2} \pm 0.011$$

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$$D_V = 1370 \pm 64 \text{ and } \Omega_m h^2 = 0.130 (n_s/0.98)^{-1.2} \pm 0.011$$

- Caveat:
 - Constraints are derived *using* Λ CDM

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- Using $n_s \sim 0.73$ the constraint is:

$$\Omega_m h_{\text{out}}^2 = 0.185 \pm 0.011 , \quad (6.9)$$

- It agrees with our value (0.205 ± 0.01) within 2σ .

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- On the other hand:

$$D_A(0.35) = 1375 \text{ Mpc} \quad \text{for } \Lambda CDM$$

$$D_A(0.35) = 1850 \text{ Mpc} \quad \text{for EdS with } h_{\text{out}} \sim 0.45,$$

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$$D_A(0.35) = 1850 \text{ Mpc} \quad \text{for EdS with } h_{\text{out}} \sim 0.45 ,$$

- Not consistent with** Eisenstein et al., 2005:

$$D_V(0.35) = 1370 \pm 64 \text{ Mpc} ,$$

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- The problem is the low value of h_{out} from CMB!
($h_{\text{out}} \sim 0.56$ would work...)

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Flts well, but analysis with full CMB not done yet. It can also fit $D(0.35)/D(0.2)$ (Percival et al.)

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Fits well, but analysis with full CMB not done yet. It can also fit $D(0.35)/D(0.2)$ (Percival et al.)
 - Fit CMB with higher h
(Non-compensated Void?)

Radial BAO

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- It is possible to look (Gaztanaga et al.'08) for the BAO scale only for the radial direction as Δz (model-independent)

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Conclusions

- It is possible to look (Gaztanaga et al.'08) for the BAO scale only for the radial direction as Δz (model-independent)
- Zibin, Moss & Scott '08: it does not fit (*Gpc Void*) together with full CMB (which they fit with very low h and non-compensated Void)
- Garcia-Bellido & Haugboelle '08: it fits as well as Λ CDM(*Gpc Void*), but only first peak location and SN Union (no full CMB).

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- How much Observer can be off-center?
- Observer at Distance d_O
- $\frac{\delta T}{T} \sim v_O \sim \dot{d}_O$
- CMB dipole $\leq 10^{-3}$ if $d_O \sim 15 - 20$ Mpc (Tomita et al., Alnes et al.'06)

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¹⁰Alnes et al. '06

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- Higher multipoles go as higher powers of v_O : negligible¹⁰.
- Bulk dipole of the same size of our dipole 600km/s (Kashlinsky et al. '08: 600 – 1000km/s)

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- All objects inside the Void have some peculiar velocity

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- This gives rise to $\frac{\delta T}{T} \sim \frac{v}{c}$ and spectrum distortions (kinetic SZ effect)

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Conclusions

- All objects inside the Void have some peculiar velocity
- This gives rise to $\frac{\delta T}{T} \sim \frac{v}{c}$ and spectrum distortions (kinetic SZ effect)
- Goodman '95: $v/c \lesssim 0.01$ (at $z \sim 0.2$)
- Caldwell-Stebbins '07-'08: rule out Voids with $z_b > 0.9$

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- Garcia-Bellido & Haugboelle: using 9 clusters ($0.2 \leq z \leq 0.6$) with detection of spectral distortion one finds:
 $\bar{v} = 320 \text{ km/sec}$ and $\sigma = 1600 \text{ km/sec}$ (σ expected is only about 400 km/sec!)
- Exclude $L > 1.5 \text{ Gpc}$, with $\Omega_{IN} = 0.23$.

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- Exclude $L > 1.5 \text{ Gpc}$, with $\Omega_{IN} = 0.23$.
- But Kashlinsky et al. measure high $\frac{v}{c} \sim 1000 \text{ km/sec}$ on $300 \text{ Mpc}/h$ (they assume kSZ, but do not see spectral distortions).

Other ways to test the Copernican principle

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- Clarkson, Bassett and Lu '08: a consistency relation:
 $\mathcal{C}(z) = 0$ for FLRW, at *all* z ,

$$\mathcal{C}(z) \equiv 1 + H^2(DD'' - D'^2) + HH'DD'$$

- Uzan, Clarkson and Ellis '08: Time drift of the redshift (over 10 years)

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- Similarly the expansion is anisotropic if d_O nonzero¹¹.

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- Two papers claim significant anisotropy in H :
 - D.Schwarz & Weinhorst '07: in the SNIa dataset ($> 95\% C.L.$)
 - McClure & Dyer '07: in the *Hubble Key Project* data ($9 - 20 \text{ km/sec}$)

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- In addition this should be correlated with CMB dipole
- Also to be explored: non-sphericity of Void

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- It is well-known that linear Φ constant in Matter Dominated Universe

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- It is well-known that linear Φ constant in Matter Dominated Universe
- If they evolve instead \Rightarrow photon feels $\Delta\Phi$ inside structures \Rightarrow additional secondary CMB anisotropy
- Correlation of CMB with LSS

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- Detected with some significance by several groups at low- l
- Consistent with $\Omega_\Lambda \sim 0.7$

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- Consistent with $\Omega_\Lambda \sim 0.7$
- Can we get this in our scenario?
 - Inside the Void
 - If there is curvature
 - If there are other big Voids in the sky \Rightarrow nonlinear evolution of Φ
- Effect of order $(L/r_{\text{hor}})^3 \sim \mathcal{O}(10^{-5})$

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- A Void of at least $L \sim 200 - 250 \text{ Mpc}/h$ scale consistent with WMAP and SNIa (Riess data), and local h
- δ quite large (~ 0.4)
Incompatible with the expected value ($\delta \sim 0.04$).

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- But some observations seem to indicate such structures (need for more observations)
- Need for larger Void to fit Union data ($L \gtrsim 500 \text{ Mpc}/h$)
- More data will discriminate (especially SDSS-II for Supernovae)

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- Consistency with BAO problematic: H_0 too low
- Even larger Void? $\mathcal{O}(1 \text{ Gpc}/h)$ (Garcia-Bellido & Haugboelle '08)

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- Checking with curved models (*work in progress*)

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- Checking with curved models (*work in progress*)
- Non-compensated Voids? (monopole $T_0 = 2.73K$ gets large correction) Zibin-Moss-Scott '08

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- Observer has to sit **near the center** (10 – 20Mpc in radial position)

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- And would be consistent with recent Large Bulk flow measurements

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- Requires peculiar primordial spectrum: low tilt, large running.
- Analysis of LSS and Lyman- α forest to be included
- ISW effect to be included
- Check if the higher Ω_m/Ω_b is compatible with other data

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Backreaction:

- Theoretically challenging
- A quantitative realistic calculation still missing

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Backreaction:

- Theoretically challenging
- A quantitative realistic calculation still missing

Void:

- Many observations to reconcile together
- Before considering it as a valid alternative to Λ CDM
- More work to be done (and more data will soon discriminate)

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Conclusions

Can an Inhomogeneous Universe mimic Dark Energy?

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