

# Can an Inhomogeneous Universe mimic Dark Energy?

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CERN

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Void vs Dark  
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Backreaction

Light  
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SN Ia Hubble diagram

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# The need for Dark Energy

- In Standard Cosmology we use the Friedmann-Lemaître-Robertson-Walker model.

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- We compute  $D_L$  (or  $D_A$ ) and  $z$

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- We use this to interpret several observations (SN Ia, Hubble constant, CMB, Baryon Acoustic Oscillations,...)

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- To fit the observations we need a  $p < 0$  term (“Dark Energy”).

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- We use this to interpret several observations (SNIa, Hubble constant, CMB, Baryon Acoustic Oscillations,...)
- To fit the observations we need a  $p < 0$  term (“Dark Energy”).
- **Problem:** We do not understand
  - the amount (why of the same amount as Matter today)?
  - its nature (is it vacuum energy?)

# Two main pieces of evidence

- **SN Ia** is incompatible with deceleration (independently on other observations)

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# Two main pieces of evidence

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  - Assuming them as standard candles.

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But good fit<sup>2</sup> also with  $\Omega_\Lambda = 0$  (*flat*)

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- low- $h$  (0.45)
- non-standard primordial spectrum

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- The two dataset,

- **SN Ia**
- **CMB together with measured  $h$** :  $0.55 \lesssim h \lesssim 0.8$

are strong evidence for  $\Omega_\Lambda \sim 0.7$ .

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- Other observations (**BAO** and LSS...) fit consistently

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# Is there any alternative?

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- Look for some interesting critical point of view and other logical possibilities

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- Look for some interesting critical point of view and other logical possibilities
- What happens to observations when we have departure from a *homogeneous* model?

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- Look for some interesting critical point of view and other logical possibilities
- What happens to observations when we have departure from a *homogeneous* model?
- Can we accomodate for *all* this evidence if we relax (to some degree) homogeneity?

# Homogenous Universe: a good approximation?

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- At  $z \gg 1$  (CMB epoch, for example) tiny density fluctuations on all observed scales.

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- ..at late times  $\delta \equiv \frac{\delta\rho}{\rho} > 1$  for all scales  
 $L \lesssim \mathcal{O}(10)/h \text{ Mpc}$  (1% of Hubble radius)

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- ..at late times  $\delta \equiv \frac{\delta\rho}{\rho} > 1$  for all scales  
 $L \lesssim \mathcal{O}(10)/h \text{ Mpc}$  (1% of Hubble radius)
- Superclusters upto few hundreds of Mpc (10% of Hubble radius), nonlinear objects ("cosmic web")

# SDSS data ("The cosmic web")

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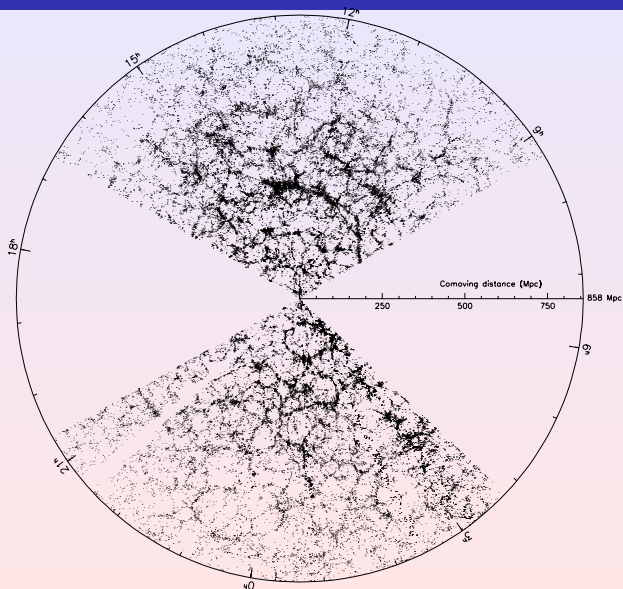
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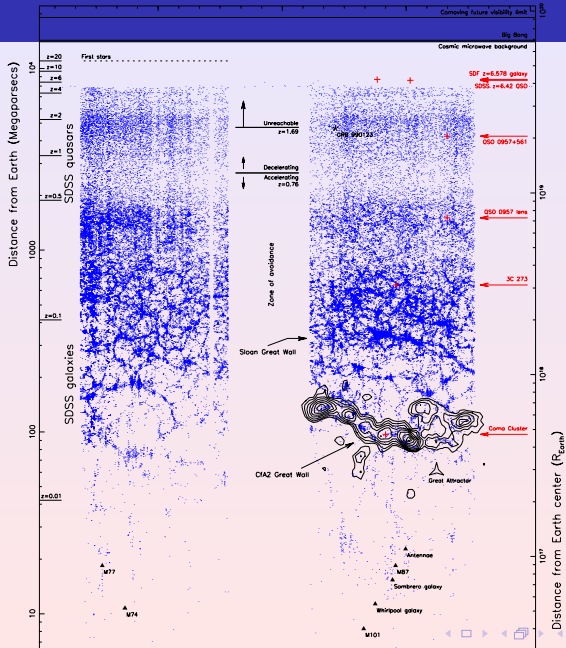
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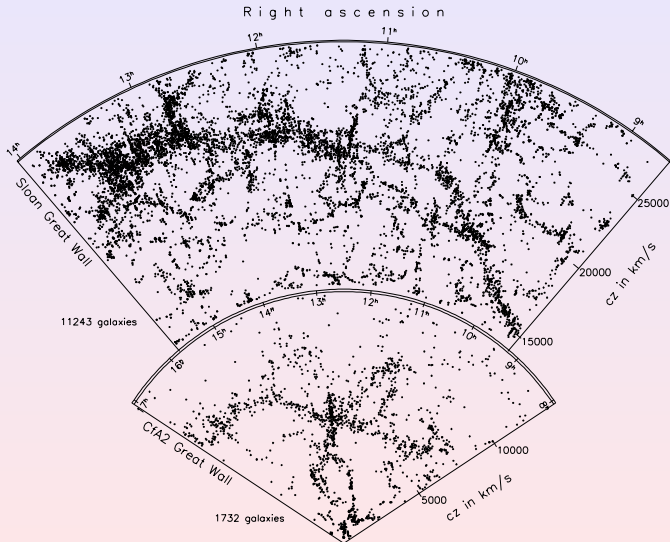
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# Three physical effects of inhomogeneities

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In general:

- Backreaction

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perturbations affect the background

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Light meets voids and structures. Do they compensate?

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Light meets voids and structures. Do they compensate?

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What if we live in a local void?



# A few words on backreaction

- Averaging of Einstein's Equations  
Nonlinearity  $\Rightarrow$  extra terms in Friedmann equations<sup>3</sup>

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<sup>3</sup>G.Ellis-W.Stoeger '67, Futamase, Sakai et al, Buchert '95, S. Rasanen '03 ,  
E.Kolb-S.Matarrese-A.Riotto-A. N. '04 , A. N. '06

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- Averaging of Einstein's Equations  
Nonlinearity  $\Rightarrow$  extra terms in Friedmann equations<sup>3</sup>
- Consider a comoving **inhomogeneous** metric ( $p=0$ )

$$ds^2 = -dt^2 + a^2(t) dx^i dx^j h_{ij}(t, \mathbf{x})$$

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$$V_D = \int_D \sqrt{h} d^3x, \quad a_D(t) \equiv \left( \frac{V_D}{V_{D_0}} \right)^{1/3};$$

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$$\frac{\ddot{a}_D}{a_D} = -\frac{4\pi G}{3}(\rho_{\text{eff}} + 3P_{\text{eff}}),$$

$$\left(\frac{\dot{a}_D}{a_D}\right)^2 = \frac{8\pi G}{3}\rho_{\text{eff}},$$

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# The extra terms

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## • Where

$$\begin{aligned}\rho_{\text{eff}} &= \langle \rho \rangle_D - \frac{Q_D}{16\pi G} - \frac{\langle R \rangle_D}{16\pi G} \\ P_{\text{eff}} &= -\frac{Q_D}{16\pi G} + \frac{\langle R \rangle_D}{48\pi G},\end{aligned}$$

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$$\rho_{\text{eff}} = \langle \rho \rangle_D - \frac{Q_D}{16\pi G} - \frac{\langle R \rangle_D}{16\pi G}$$

$$P_{\text{eff}} = -\frac{Q_D}{16\pi G} + \frac{\langle R \rangle_D}{48\pi G},$$

- The *real* question: how large is it?

# Perturbatively: 2<sup>nd</sup> order

- On a large Domain the dominant term has the form<sup>4</sup> :

$$\frac{\overline{H_D - H}}{H} = \frac{25}{54} \frac{1}{a^2 H^2} \overline{\langle \varphi \nabla^2 \varphi \rangle}$$

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$$= A^2 \frac{a}{a_0} \left( \frac{h \Gamma \text{Mpc}^{-1}}{H_0} \right)^2 \int_0^\infty dq q T^2(q)$$

where  $A \sim 10^{-5}$ ,  $\Gamma = \Omega_M h e^{-\Omega_B - \sqrt{2} h \Omega_B / \Omega_M}$ .



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where  $A \sim 10^{-5}$ ,  $\Gamma = \Omega_M h e^{-\Omega_B - \sqrt{2} h \Omega_B / \Omega_M}$ .

- Largest contribution from  $\mathcal{O}(10 - 50) \text{Mpc}/h$

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$$= A^2 \frac{a}{a_0} \left( \frac{h \Gamma \text{Mpc}^{-1}}{H_0} \right)^2 \int_0^\infty dq q T^2(q)$$

where  $A \sim 10^{-5}$ ,  $\Gamma = \Omega_M h e^{-\Omega_B - \sqrt{2} h \Omega_B / \Omega_M}$ .

- Largest contribution from  $\mathcal{O}(10 - 50) \text{Mpc}/h$

$$\frac{\overline{H_D - H}}{H} \simeq 10^{-5}$$

- Small, but not  $10^{-10}$ !

<sup>4</sup> L. Hui-U. Seljak '95, S. Rasanen'03, E. W. Kolb-S. Matarrese-A. N. A. Riotto '04, ...

# Perturbatively: 2<sup>nd</sup> order

- On a large Domain the dominant term has the form<sup>4</sup> :

$$\frac{\overline{H_D - H}}{H} = \frac{25}{54} \frac{1}{a^2 H^2} \overline{\langle \varphi \nabla^2 \varphi \rangle}$$
$$= A^2 \frac{a}{a_0} \left( \frac{h \Gamma \text{Mpc}^{-1}}{H_0} \right)^2 \int_0^\infty dq q T^2(q)$$

where  $A \sim 10^{-5}$ ,  $\Gamma = \Omega_M h e^{-\Omega_B - \sqrt{2} h \Omega_B / \Omega_M}$ .

- Largest contribution from  $\mathcal{O}(10 - 50) \text{Mpc}/h$

$$\frac{\overline{H_D - H}}{H} \simeq 10^{-5}$$

- Small, but not  $10^{-10}$ ! Enhanced by  $\left( \frac{k_{EQ}}{H_0} \right)^2$

<sup>4</sup> L. Hui-U. Seljak '95, S. Rasanen'03, E. W. Kolb-S. Matarrese-A.N.-A. Riotto '04, ...

# Power counting

- What about higher ( $n^{\text{th}}$ ) orders <sup>5</sup>?

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- What about higher ( $n^{\text{th}}$ ) orders <sup>5</sup>?
- They go as

$$\overline{\langle \varphi (\nabla^2 \varphi)^{n-1} \rangle}$$

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- They go as

$$\overline{\langle \varphi (\nabla^2 \varphi)^{n-1} \rangle}$$

- We can write the  $n^{\text{th}}$  order as

$$10^{-5} \epsilon^{n-1}$$

where roughly

$$\epsilon \equiv \frac{A}{1+z} \left( \frac{h \Gamma \text{Mpc}^{-1}}{H_0} \right)^2 \times \text{Int}$$

with

$$\text{Int} = \int dq T^2(q) \approx 0.02$$

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- $\epsilon = \mathcal{O}(1)$  today

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- ...each term in the series is of  $\mathcal{O}(10^{-5})$  !

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- Need non-perturbative treatment

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- ...each term in the series is of  $\mathcal{O}(10^{-5})$  !
- Do they **sum** up to  $10^{-5}$  or more??
- Need non-perturbative treatment
- Note:  $\epsilon \ll 1$  at high  $z$

$$\epsilon(z)$$

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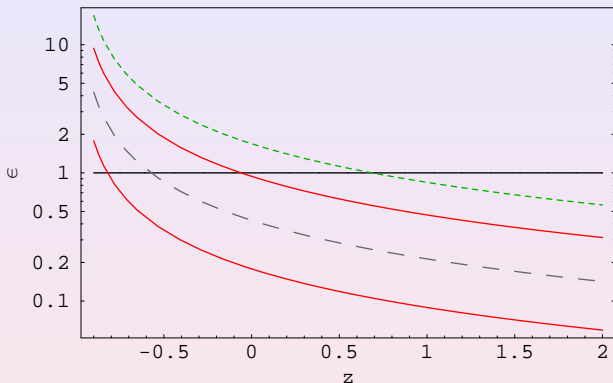
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**Figure:** Grey dashed line: central value,  
Red solid lines:  $2\sigma$  ranges  
(We used the growth factor as in matter domination. For comparison, green dotted line:  $\Omega_M = 1$ ).

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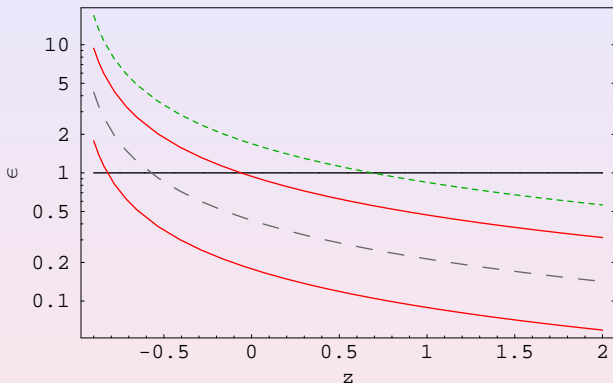
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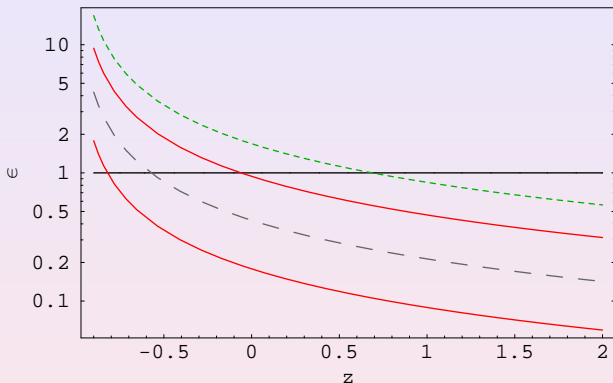
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- Even in absence of average effect on  $H(z)$  : corrections to photon trajectories

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- In fact, actually we measure distances  $D$  and redshifts  $z$

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- Even in absence of average effect on  $H(z)$  : corrections to photon trajectories
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- All information from expansion comes from plots  $D - z$

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- Even in absence of average effect on  $H(z)$  : corrections to photon trajectories
- In fact, actually we measure distances  $D$  and redshifts  $z$
- All information from expansion comes from plots  $D - z$
- Cannot disentangle this from backreaction
- Compute  $\frac{\Delta z}{1+z}$  and  $\frac{\Delta D}{D}$  in the presence of structures

# LTB exact solutions

- Consider Lemaître-Tolman-Bondi exact solutions of E.E. (with  $p = 0$ ) which is

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- Consider Lemaître-Tolman-Bondi exact solutions of E.E. (with  $p = 0$ ) which is
  - inhomogeneous
  - nonlinear

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- Consider Lemaître-Tolman-Bondi exact solutions of E.E. (with  $p = 0$ ) which is
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  - nonlinear
  - Spherically symmetric

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- Consider Lemaître-Tolman-Bondi exact solutions of E.E. (with  $p = 0$ ) which is
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- We consider two configurations:

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- Consider Lemaître-Tolman-Bondi exact solutions of E.E. (with  $p = 0$ ) which is
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  - LTB spheres embedded in FLRW ("Swiss-Cheese")

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  - Spherically symmetric
- We consider two configurations:
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  - LTB with shells of periodically varying density ("Onion")
- We study null geodesic in this metric

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- Net effect from one hole<sup>6</sup> :  $\frac{\Delta z}{1+z} \approx (L/r_H)^3 f(\delta)$

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<sup>6</sup>T. Biswas-A. N. '06-'07

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- At 2<sup>nd</sup> order usual Rees-Sciama effect  $(L/r_H)^3 \delta^2$

---

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- $f(\delta)$  does *not* compensate the suppression for  $\delta \gg 1$

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- Still small (for late acceleration)

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- Still small (for late acceleration)
- Interesting in the CMB, as a Rees-Sciama effect.

---

<sup>6</sup>T. Biswas-A. N. '06-'07



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- Net effect scales as  $\frac{\Delta z}{1+z} \approx (L/r_H)^2 f(\delta)$ <sup>7</sup>

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<sup>7</sup> Brouzakis-Tetradis-Tzavara '06, Kolb-Matarrese-Riotto '07, T. Biswas-A. N. '07

<sup>8</sup> S. Weinberg '76, Brouzakis et al. '06-'07

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- Tight packing:  $N_{\text{holes}} \mathcal{O}(L/r_H)^3 = \mathcal{O}(L/r_H)$
- Not so small...
- But it should have zero angular average (unlike  $z$ )<sup>8</sup>

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<sup>7</sup> Brouzakis-Tetradis-Tzavara '06, Kolb-Matarrese-Riotto '07, T. Biswas-A. N. '07

<sup>8</sup> S. Weinberg '76, Brouzakis et al. '06-'07

# Beyond LTB?

- Reliable result or limited by the symmetries of the model?

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- What happens without spherical symmetry?

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- Szekeres swiss-cheese model with asymmetric holes  
(Bolejko '08)

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(Bolejko '08) *Effects of similar size*

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- The cheese feels *no backreaction* by construction
- What happens without spherical symmetry?
- Szekeres swiss-cheese model with asymmetric holes  
(Bolejko '08) *Effects of similar size*
- But still special: the cheese feels *no backreaction* of the holes

# A local fluctuation?

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<sup>9</sup>Tomita '98, Tomita '00, Celerier '01, Wiltshire '05, Moffat '05, Alnes et al. '05, Mansouri et al. '06, Biswas & A.N.'07, Garcia-Bellido and Haugboelle '08, Zibin et al. '08 ...

# A local fluctuation?

- Suppose that we live in a peculiar local region

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Conclusions

- Suppose that we live in a peculiar local region
- $\Rightarrow$  low  $z$  observations may be very different from average.

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- One realizes that acceleration is inferred **comparing low  $z$  with high  $z$ ...**

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- $\Rightarrow$  low  $z$  observations may be very different from average.
- One realizes that acceleration is inferred **comparing low  $z$  with high  $z$ ...**
- Can this mimic acceleration <sup>9</sup>?

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- Mustapha, Hellaby, Ellis '97: show that LTB can reproduce any observations

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- Mustapha, Hellaby, Ellis '97: show that LTB can reproduce any observations

- Celerier '99: showed that LTB can mimic  $\Lambda$ CDM

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### Conclusions

- Mustapha, Hellaby, Ellis '97: show that LTB can reproduce any observations
- Celerier '99: showed that LTB can mimic  $\Lambda$ CDM
- Tomita '01: Compensated Void 200 – 300 Mpc/ $h$  scale

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- Consider a “compensated Void” : a spherical Void plus an external shell of matter

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- Consider a “compensated Void” : a spherical Void plus an external shell of matter  
(on average same density as “external” FLRW)

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- Assumption: we live near the center

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- A void expands faster than the “external” FLRW



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- So, nearby objects inside the void redshift more

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- This can mimic acceleration (as we will see...)

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- So, nearby objects inside the void redshift more
- This can mimic acceleration (as we will see...)
- How much contrast  $\delta$  and how large  $L$  is needed?

# About Voids

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- Before going to the quantitative analysis...

- Let's review some literature and observations on Voids

# Other uses of Voids

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Conclusions

- Inoue and Silk '06: some features of the low multipole **anomalies** in the CMB data could be explained by a pair of huge Voids ( $L \sim 200 \text{ Mpc}/h$ ,  $\delta \sim -0.3$ )

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- Inoue and Silk '06: some features of the low multipole **anomalies** in the CMB data could be explained by a pair of huge Voids ( $L \sim 200 \text{ Mpc}/h$ ,  $\delta \sim -0.3$ )
- The CMB has a **Cold Spot** (M. Cruz et al. ('06 and '07)): it could be explained by another similar Big Void (Inoue and Silk '06)

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- The Cold Spot in the CMB claimed to be correlated with an underdense region in the LSS (Rudnick, Brown and Williams '07, but see Huterer and ...)



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- The Cold Spot in the CMB claimed to be correlated with an underdense region in the LSS (Rudnick, Brown and Williams '07, but see Huterer and ...)
- It could be detected via **lensing** (S. Das and D. Spergel '08) and via **non-gaussian** coupling Rees-Sciama effect - lensing (I. Masina and A.N., in preparation)

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- Some observational evidence for a local large underdense region ( $\sim 25\%$  less dense,  $r \sim 200 \text{ Mpc}/h$ ) from number counts of galaxies (2MASS)  
(Frith et al. Mon. Not. Roy. Astron. Soc. **345**, 1049 (2003))
- It would represent a  $4\sigma$  fluctuation, at odds with  $\Lambda\text{CDM}$ .

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(Frith et al. Mon. Not. Roy. Astron. Soc. **345**, 1049 (2003))
- It would represent a  $4\sigma$  fluctuation, at odds with  $\Lambda\text{CDM}$ .
- Many Large Voids identified via **ISW** effect in the SDSS LRG catalog (about  $100 \text{ Mpc}/h$  radius) (Granett et al. '08)
- Also in contradiction with  $\Lambda\text{CDM}$ :  $P < 10^{-8}$  (Sarkar & Hunt '08)

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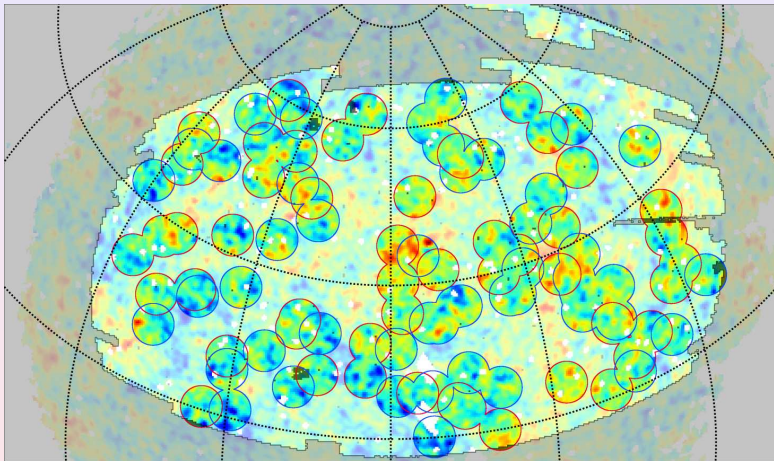


Figure: Granett, Neyrinck & Szapudi '08

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- Recent measurement (Kashlinsky et al.'08): **very large coherent motion** on  $300\text{Mpc}/h$  scale, inconsistent with  $\Lambda\text{CDM}$
- Could be due to very large scale inhomogeneous matter distribution

# Large bulk motion?

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Conclusions

- Recent measurement (Kashlinsky et al.'08): **very large coherent motion** on  $300\text{Mpc}/h$  scale, inconsistent with  $\Lambda\text{CDM}$
- Could be due to very large scale inhomogeneous matter distribution
- Watkins, Feldman & Hudson '08: use peculiar velocities of various (4500) objects in a  $100\text{Mpc}/h$  radius. Find  $400\text{km}/\text{sec}$  (expected  $100\text{km}/\text{sec}$ )

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# A “Minimal” Void ?

- What is the size we need to mimick DE?

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- What is the size we need to mimick DE?
- It will turn out that a Minimal Void needs at least the same size (*for Riess '07 SN Ia and WMAP*)

# A “Minimal” Void ?

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Conclusions

- What is the size we need to mimick DE?
- It will turn out that a Minimal Void needs at least the same size (*for Riess '07 SN Ia and WMAP*)
- $r_{\text{Void}} \sim 200 - 250 \text{ Mpc}/h$  and  $\delta \sim -0.4$

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 $\delta \sim 0.03 - 0.05$ , using *linear* and *Gaussian* spectrum

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  - Non-standard structure formation?

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- **Problem (I):** on this scale the typical contrast is:  
 $\delta \sim 0.03 - 0.05$ , using *linear* and *Gaussian* spectrum
- Can we ever get huge Voids?
  - Percolation of Voids?
  - Non-standard structure formation?
  - Non-gaussianity?
  - Nucleation of primordial Bubbles



# Need for a Larger Void

- **Problem (II):** even larger size seems necessary for other observations

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- **Problem (II)**: even larger size seems necessary for other observations
- For **UNION** (...) data **500Mpc/h** required

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- **Problem (II)**: even larger size seems necessary for other observations
- For **UNION** (...) data **500Mpc/h** required and the fit is not very good:  $\chi^2 = 343$  (306 d.o.f.)

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- Or adding open curvature

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- A very good fit obtained with (**2 Gpc**) (Garcia-Bellido & Haugboelle)
- Or adding open curvature
- Need for a Large Void also for **BAO**: a better fit is obtained with at least (1 Gpc) (Garcia-Bellido & Haugboelle).

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$$ds^2 = -dt^2 + \frac{R'^2(r, t)}{1 + 2r^2 k(r)} dr^2 + R^2(r, t)(d\theta^2 + \sin^2 \theta d\varphi^2)$$

with comoving coordinates  $(r, \theta, \varphi)$  and proper time  $t$ .



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with comoving coordinates  $(r, \theta, \varphi)$  and proper time  $t$ .

- Spherically symmetric.

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with comoving coordinates  $(r, \theta, \varphi)$  and proper time  $t$ .

- Spherically symmetric.
- Einstein equations:

$$\frac{1}{2} \frac{\dot{R}^2(r, t)}{R^2(r, t)} - \frac{GM(r)}{R^3(r, t)} = \frac{r^2 k(r)}{R^2(r, t)},$$

$$4\pi\rho(r, t) = \frac{M'(r)}{R'(r, t)R^2(r, t)},$$

# LTB metrics

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$$ds^2 = -dt^2 + \frac{R'^2(r, t)}{1 + 2r^2 k(r)} dr^2 + R^2(r, t)(d\theta^2 + \sin^2 \theta d\varphi^2)$$

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$$ds^2 = -dt^2 + \frac{R'^2(r, t)}{1 + 2r^2 k(r)} dr^2 + R^2(r, t)(d\theta^2 + \sin^2 \theta d\varphi^2)$$

It has the solutions:

- For  $k(r) > 0$  ( $k(r) < 0$ ),

$$R = \frac{GM(r)}{2r^2 |k(r)|} [\cos h(u) - 1], \quad (4.1)$$

$$t - t_b(r) = \frac{GM(r)}{[2r^2 |k(r)|]^{3/2}} [\sin h(u) - u].$$

- $k(r) = 0$ ,

$$R(r, t) = \left[ \frac{9GM(r)}{2} \right]^{1/3} [t - t_b(r)]^{2/3}.$$

# Choosing the functions

- $t_b(r) = 0$  for our purposes, and “Gauge” choice:  
 $M(r) \propto r^3$

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- $k(r)$  contains all the physical information about the profile.

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- $k(r)$  contains all the physical information about the profile.
- $k = 0$  flat FLRW,  $k = \pm 1$  open/closed FLRW.

# Choosing the functions

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- The idea is to describe structure formation  
(start with  $\delta(r, t_i) \ll 1$  and end up with  $\delta(r, t_{\text{now}}) \gg 1$ )



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 $M(r) \propto r^3$
- $k(r)$  contains all the physical information about the profile.
- $k = 0$  flat FLRW,  $k = \pm 1$  open/closed FLRW.
- The idea is to describe structure formation  
(start with  $\delta(r, t_i) \ll 1$  and end up with  $\delta(r, t_{\text{now}}) \gg 1$ )
- We play with  $k(r)$  to describe  $\delta(r, t_i)$ .

# LTB merged to FLRW

Void vs Dark  
Energy

- Matching of an LTB sphere (of radius  $L$ ) to FLRW:

$$k'(0) = k'(L) = 0,$$

$$k(L) = \frac{4\pi}{3}\Omega_k, \quad \text{for } |\Omega_k| \ll 1,$$

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We use:



$$k(r) = k_{max} \left[ \left( \frac{r}{L} \right)^4 - 1 \right]^2 \quad (\text{for } r < L)$$

$$k(r) = 0 \text{ (flat)} \quad (\text{for } r > L)$$

# LTB merged to FLRW

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$$k(r) = 0 \text{ (flat)} \quad (\text{for } r > L)$$

- Two parameters,  $L$  and  $k_{\max}$ .

# The density

- Roughly:

$$\rho(r, t) \simeq \frac{\langle \rho \rangle(t)}{1 + (t/t_0)^{2/3} \epsilon(r)},$$

$$\text{where } \langle \rho \rangle(t) \equiv \frac{M_p^2}{6\pi t^2}, \quad \text{and } \epsilon(r) \equiv 3k(r) + rk'(r).$$

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- $\epsilon \ll 1$  linear growth

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- $\epsilon \ll 1$  linear growth  $\propto a(t) \propto t^{2/3}$

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- $\epsilon \ll 1$  linear growth  $\propto a(t) \propto t^{2/3}$
- $\epsilon$  not small:  $\delta$  grows rapidly (as in Zel'dovich approx)

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# The density

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$$\text{where } \langle \rho \rangle(t) \equiv \frac{M_p^2}{6\pi t^2}, \quad \text{and } \epsilon(r) \equiv 3k(r) + rk'(r).$$

- $\epsilon \ll 1$  linear growth  $\propto a(t) \propto t^{2/3}$
- $\epsilon$  not small:  $\delta$  grows rapidly (as in Zel'dovich approx)
- We work at most with  $\delta \sim \mathcal{O}(1)$ .

# The density profile

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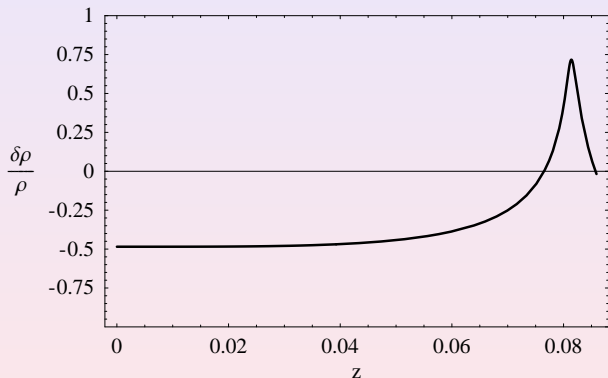
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- Solve for  $t(r)$  along a  $ds^2 = 0$  trajectory

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- Solve for  $t(r)$  along a  $ds^2 = 0$  trajectory
- Then solve for

$$\frac{dz}{dr} = \frac{(1+z(r))\dot{R}'(r, t(r))}{\sqrt{1+2r^2k(r)}}.$$

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- The result  $z(r)$  can be found numerically

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$$\frac{dz}{dr} = \frac{(1+z(r))\dot{R}'(r, t(r))}{\sqrt{1+2r^2k(r)}}.$$

- The result  $z(r)$  can be found numerically
- We also have some very good analytical approximations

# Luminosity (Angular) Distance

Void vs Dark  
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- Always in GR, luminosity distance and angular distance:

$$D_L = D_A(1 + z)^2.$$

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# Luminosity (Angular) Distance

- Always in GR, luminosity distance and angular distance:

$$D_L = D_A(1 + z)^2.$$



$$D_A^2 \equiv \frac{dA}{d\Omega} = \frac{d\theta_S d\phi_S \sqrt{g_{\theta\theta} g_{\phi\phi}}}{d\bar{\theta}_O d\bar{\phi}_O}$$

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$$D_A^2 \equiv \frac{dA}{d\Omega} = \frac{d\theta_S d\phi_S \sqrt{g_{\theta\theta} g_{\phi\phi}}}{d\bar{\theta}_O d\bar{\phi}_O} = \frac{d\theta_S d\phi_S}{d\bar{\theta}_O d\bar{\phi}_O} R^2|_S,$$

- If observer in the center:

$$D_A^2 = R^2|_S.$$

# Luminosity (Angular) Distance

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- Always in GR, luminosity distance and angular distance:

$$D_L = D_A(1 + z)^2.$$

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- $$D_A^2 \equiv \frac{dA}{d\Omega} = \frac{d\theta_S d\phi_S \sqrt{g_{\theta\theta} g_{\phi\phi}}}{d\bar{\theta}_O d\bar{\phi}_O} = \frac{d\theta_S d\phi_S}{d\bar{\theta}_O d\bar{\phi}_O} R^2|_S,$$

- If observer in the center:

$$D_A^2 = R^2|_S.$$

- For generic observer (but radial trajectory):

$$D_A = R_S \left( R_O \int_{r_O}^{r_S} \frac{R'(r, t(r))}{(1 + 2E(r))(1 + z(r))R(r, t(r))^2} dr \right),$$

# Analytical approximation

Void vs Dark  
Energy

$$f \equiv \frac{\sqrt[3]{2}(\cosh(u) - 1)}{3^{2/3}(\sinh(u) - u)^{2/3}} - 1 \quad (4.2)$$

$$u_0 = 6^{1/3}(\sinh(u) - u)^{1/3}. \quad (4.3)$$

Then, one can use this function in the following equations:

$$\tau(r) = \tau_0 - \frac{\pi}{9}\gamma^2\bar{M}r[1 + f(\gamma^2\tau_0^2k(r))], \quad (4.4)$$

$$1 + z(r) = \left(\frac{\tau_0}{\tau(r)}\right)^2 \exp\left[\frac{4\pi\gamma^2\bar{M}r}{9}f(\gamma^2\tau_0^2k(r))\right] \quad (4.5)$$

$$D_L(r) = \frac{\pi}{3}\gamma^2r\tau(r)^2[1 + f(\gamma^2\tau_0^2k(r))][1 + z(r)]^2 \quad (4.6)$$

$$\tau_0 = \left(\frac{2\bar{M}}{3H_0}\right)^{1/3} \quad (4.7)$$

$$\gamma = \left(\frac{9\sqrt{2}}{\pi}\right)^{1/3} \quad (4.8)$$

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- Evidence for acceleration comes from **mismatch** between:
  - measurements at low redshift (  $0.03 \lesssim z \lesssim 0.08$  )
  - high- $z$  SN (roughly  $0.4 \lesssim z \lesssim 1$  )

# High and low $z$

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### Conclusions

- Evidence for acceleration comes from **mismatch** between:
  - measurements at low redshift ( $0.03 \lesssim z \lesssim 0.08$ )
  - high- $z$  SN (roughly  $0.4 \lesssim z \lesssim 1$ )
- We choose large  $r_{\text{void}}$  (at  $z \approx 0.08 - 0.09$ )



# High and low $z$

## Void vs Dark Energy

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- Evidence for acceleration comes from **mismatch** between:

- measurements at low redshift (  $0.03 \lesssim z \lesssim 0.08$  )
- high- $z$  SN (roughly  $0.4 \lesssim z \lesssim 1$  )

- We choose large  $r_{\text{Void}}$  (at  $z \approx 0.08 - 0.09$ )

- $\Rightarrow$  The Local Bubble is different from the average.

Outside just matter dominated (even if there are other Bubbles, their effect is small)

# Roughly

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- At high  $z$  ( $z \gtrsim 0.1$ ), just matter dominated Universe:

$$D_{\text{EdS}}(z) \approx \frac{2}{H_{\text{out}}}(1 + z - \sqrt{1 + z}).$$

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- At low  $z$  "open-like" Universe with a different  $H$

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- Two (reduced) Hubble parameters:  $h$  and  $h_{\text{out}}$

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- Rapid transition near the shell-like structure

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- At low  $z$  "open-like" Universe with a different  $H$
- Two (reduced) Hubble parameters:  $h$  and  $h_{\text{out}}$
- Rapid transition near the shell-like structure
- $h$  corresponds to what is measured

# $\Delta m$ for different models

## Void vs Dark Energy

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### Fitting the data

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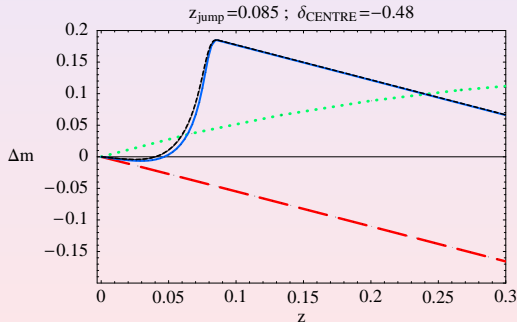
#### WMAP

#### Putting things together

### Other cosmological observations

### Conclusions

- Magnitude is  $m \equiv 5 \text{Log}_{10} D(z)$
- The open “empty” Universe is subtracted ( $\Omega_K = -1$ )



# $m - z$ diagram: Riess data

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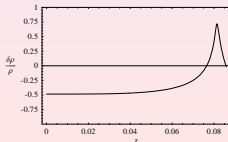
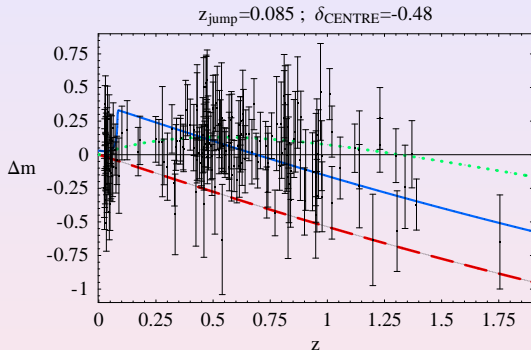
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# Finding the best fit

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- We fix several values of  $L$

# Finding the best fit

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### Conclusions

- We fix several values of  $L$

- What matters is just the Jump:  $\mathcal{J} \equiv \frac{h}{h_{OUT}}$

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Conclusions

- We fix several values of  $L$
- What matters is just the Jump:  $\mathcal{J} \equiv \frac{h}{h_{OUT}}$
- This is also related to the central density contrast:  
$$\mathcal{J} = 2 - (1 - \delta_0)^{1/3}$$

# Finding the best fit

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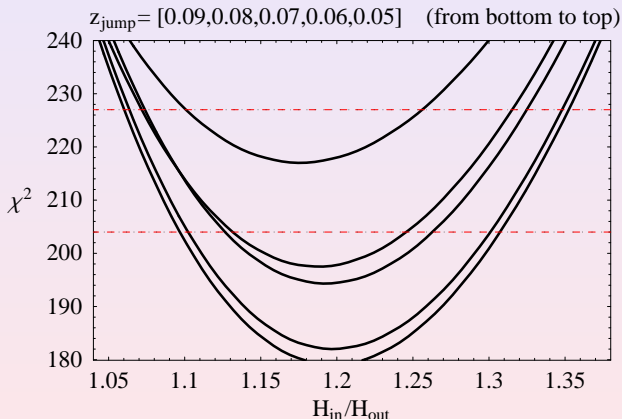
Other  
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Conclusions

- We fix several values of  $L$
- What matters is just the Jump:  $\mathcal{J} \equiv \frac{h}{h_{OUT}}$
- This is also related to the central density contrast:  
$$\mathcal{J} = 2 - (1 - \delta_0)^{1/3}$$
- We vary  $\mathcal{J}$  and compute the  $\chi^2$ .

# Fitting SNIa with a Jump

Riess et al. dataset, [astro-ph/0611576](#) (182 SNIa)



**Figure:** The red dashed lines are 10% and 1% goodness-of-fit (182 data points)

# LTB Void fit

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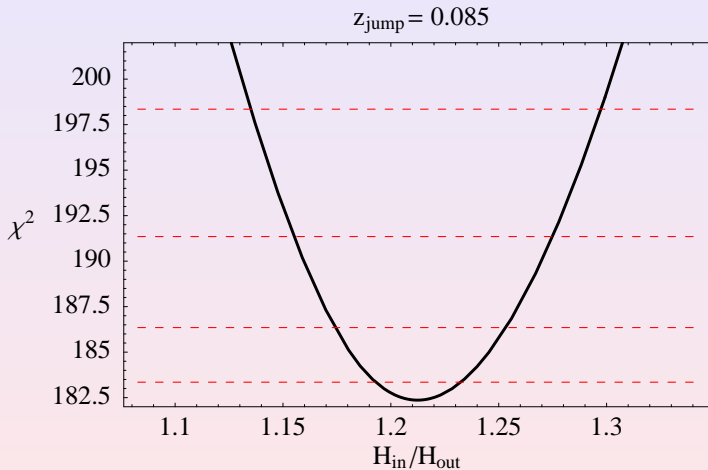
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**Figure:** Here we use the full LTB model. We show  $1\sigma$ ,  $2\sigma$ ,  $3\sigma$  and  $4\sigma$  intervals (using likelihood  $\propto e^{-\chi^2/2}$ ).

# $\chi^2$ : Riess data

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**Table:** Comparison with data (full data set of Riess et al.)

Model	$\chi^2$ (181 d.o.f.)
$\Lambda$ CDM (with $\Omega_M = 0.27, \Omega_\Lambda = 0.73$ )	160
EdS (with $\Omega_M = 1, \Omega_\Lambda = 0$ )	274
Void ( $\sqrt{\langle \delta^2 \rangle} \approx 0.4$ on $L = 250/h\text{Mpc}$ )	182

Remarks:

- With instrumental error only: no smooth curve can give a good fit
- Estimated error from intrinsic variability added in quadrature

# $\chi^2$ : Riess data

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Remarks:

- With instrumental error only: no smooth curve can give a good fit
- Estimated error from intrinsic variability added in quadrature
- Not as good as  $\Lambda$ CDM
- Becomes better including curvature  $\Omega_k$  *outside*



# $\chi^2$ : UNION data

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**Table:** Comparison with data (Union data)

Model	$\chi^2$ (307 d.o.f.)
$\Lambda$ CDM (with $\Omega_M = 0.27, \Omega_\Lambda = 0.73$ )	304
Void ( $\sqrt{\langle \delta^2 \rangle} \approx 0.4$ on $L = 500/h\text{Mpc}$ )	340
Void ( $\sqrt{\langle \delta^2 \rangle} \approx 0.7$ on $L = 2000/h\text{Mpc}$ )	304

Remarks:

- It seems necessary to consider a larger Void ( $Gpc$  scale)
- Or add curvature? (work in progress)

# UNION fit with 2Gpc

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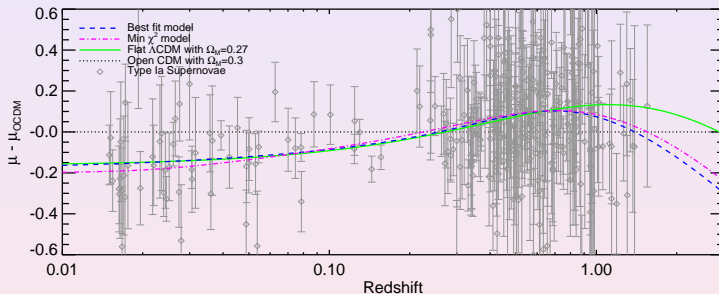
SN Ia Hubble diagram

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**Figure:** Taken from Garcia-Bellido & Haugboelle '08  
(similar fits also in Zibin et al. '08)

# Outline

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- 2 Backreaction
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# The $\Lambda$ CDM fit

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- We try to fit the WMAP 3-yr data

# The $\Lambda$ CDM fit

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Conclusions

- We try to fit the WMAP 3-yr data
- We look at TT and TE correlations, using CosmoMC

# How do we fit?

- In principle: we should compute propagation in EdS from  $z = 1100$  to  $z \sim 0.1$ , and then in the Bubble

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- In principle: we should compute propagation in EdS from  $z = 1100$  to  $z \sim 0.1$ , and then in the Bubble
- Possible “secondary” effects in the Bubble:

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  - Small offset to  $D_A$  and  $T_0$  of  $\mathcal{O}(r_{\text{Void}}/r_{\text{Hor}})^2$



# How do we fit?

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**Small because of compensation**

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**Small because of compensation**
  - off-center location: dipole

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  - $n_s$  plus running  $\alpha_s$

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  - off-center location: dipole and integrated effect (low- $l$ )
  - Non-sphericity (again effect on low- $l$ )
- We do not consider them: just EdS with  $h_{\text{out}}$ , with some assumptions on the primordial spectrum:
  - $n_s$  plus running  $\alpha_s$
  - Flat spectrum plus bump (as in P. Hunt and S. Sarkar, arXiv:0706.2443 [astro-ph]; A. Blanchard, M. Douspis, M. Rowan-Robinson and S. Sarkar, Astron. Astrophys. **412**, 35 (2003) [arXiv:astro-ph/0304237]. )



# Priors ( $\Lambda$ CDM)

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The usual prior set is:

- Allow for nonzero  $\Omega_\Lambda$ .
- Power-law spectrum with index  $n_s$ .
- (eventually with running  $\alpha_s$ )
- $P(k) \propto k^{n_s(k_0) + \frac{1}{2} \ln(k/k_0) \alpha_s}$

# Priors: without $\Lambda$

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Conclusions

A different prior set, that we use:

- **Not** allow for  $\Omega_\Lambda$ .
- Power-law spectrum with index  $n_s$ .
- with running  $\alpha_s$
- $P(k) \propto k^{n_s(k_0) + \frac{1}{2} \ln(k/k_0) \alpha_s}$

# Priors: without $\Lambda$

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A different prior set, that we use:

- **Not** allow for  $\Omega_\Lambda$ .
- Power-law spectrum with index  $n_s$ .
- with running  $\alpha_s$
- $P(k) \propto k^{n_s(k_0) + \frac{1}{2} \ln(k/k_0) \alpha_s}$
- (we also may allow for some curvature)

# Fit to WMAP3

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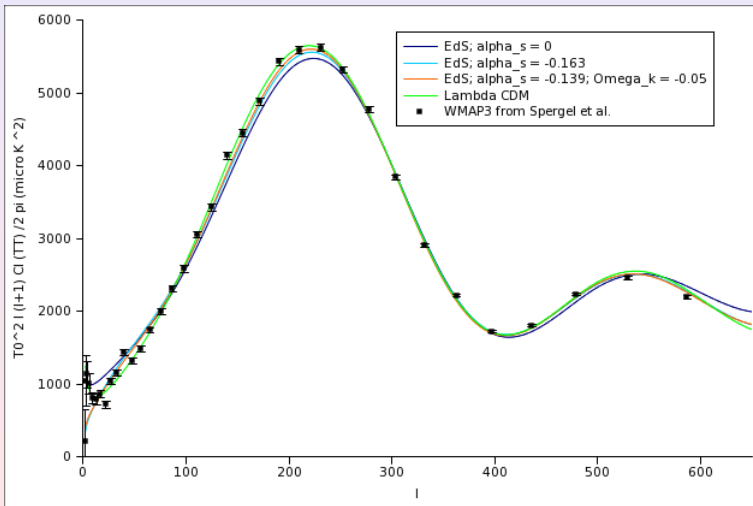
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# Goodness-of-fit

## Void vs Dark Energy

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Model	$C_l^{TT}$		$C_l^{TT} + C_l^{TE}$		Total	
	$\chi^2_{eff}$	G.F.	$\chi^2_{eff}$	G.F.	$\chi^2_{eff}$	G.F.
Concordant $\Lambda$ CDM	1038.9	4.7%	1455.2	11.3%	3538.6	41%
EdS $\alpha_S = 0$	1124.6	0%	1711.9	0%	3652.3	6%
EdS $\alpha_S \neq 0$	1057.8	1.9 %	1475.5	5.7%	3577.4	24.6%
EdS $\alpha_S, \Omega_k = -0.050$	1048.7	2.9%	1466	7.9%	3560.9	31.1%

### Table:

1<sup>st</sup> column: high- $l$  TT ( $31 \leq l \leq 1000$ )

2<sup>nd</sup> column: high- $l$  TT ( $31 \leq l \leq 1000$ ) and TE ( $24 \leq l \leq 450$ )

3<sup>rd</sup> column: total of TT ( $2 \leq l \leq 1000$ ) and TE ( $2 \leq l \leq 450$ )

# Result for parameters

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The EdS model, with running, has:

- low  $h_{\text{OUT}}$  (about  $\sim 0.45$ )

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- low  $n_S$  (about  $\sim 0.73$ )  
and large negative  $\alpha_S$  (about  $\sim -0.16$ )



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- low  $n_s$  (about  $\sim 0.73$ )  
and large negative  $\alpha_s$  (about  $\sim -0.16$ )

- larger value of  $\Omega_M/\Omega_b$  (around 10 instead of 6)

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The EdS model, with running, has:

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and large negative  $\alpha_s$  (about  $\sim -0.16$ )
- larger value of  $\Omega_M/\Omega_b$  (around 10 instead of 6)
- $\Omega_b h_{\text{out}}^2$  ( $\sim 0.018^{+0.001}_{-0.002}$ ) consistent with BBN constraint  
(which is  $0.017 \leq \Omega_b h_{\text{out}}^2 \leq 0.024$ , at 95% C.L.)

# Parameter values

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	$\Lambda$ CDM	EdS, $\alpha_s = 0$	Eds, $\alpha_s \neq 0$	Eds, $\alpha_s, \Omega_k \neq 0$
$\Omega_b h^2_{\text{out}}$	$0.022^{+0.002}_{-0.002}$	$0.022^{+0.001}_{-0.001}$	$0.018^{+0.001}_{-0.002}$	$0.019^{+0.002}_{-0.001}$
$\Omega_m h^2_{\text{out}}$	$0.106^{+0.021}_{-0.013}$	$0.198^{+0.008}_{-0.011}$	$0.186^{+0.011}_{-0.009}$	$0.167^{+0.009}_{-0.007}$
$\Omega_\Lambda$	$0.759^{+0.041}_{-0.103}$	0	0	0
$z_{re}$	$11.734^{+4.993}_{-7.619}$	$8.697^{+4.351}_{-6.694}$	$13.754^{+2.246}_{-5.752}$	$13.342^{+2.55}_{-5.011}$
$\Omega_k$	0	0	0	-0.05
$n_s$	$0.96^{+0.04}_{-0.04}$	$0.94^{+0.021}_{-0.038}$	$0.732^{+0.07}_{-0.071}$	$0.761^{+0.069}_{-0.069}$
$\alpha_s$	0	0	$-0.161^{+0.044}_{-0.044}$	$-0.13^{+0.037}_{-0.048}$
$10^{10} A_s$	$20.841^{+3.116}_{-3.442}$	$25.459^{+2.135}_{-2.766}$	$25.302^{+2.182}_{-2.968}$	$23.975^{+2.198}_{-2.448}$
$\Omega_m / \Omega_b$	$4.73^{+0.999}_{-0.485}$	$9.119^{+0.341}_{-0.357}$	$10.094^{+0.645}_{-0.489}$	$8.929^{+0.512}_{-0.541}$
$h_{\text{out}}$	$.72857^{+.05137}_{-.07393}$	$.46857^{+.00888}_{-.01307}$	$.4523^{+.01291}_{-.01129}$	$.42069^{+.01107}_{-.00919}$
Age / GYr	$13.733^{+0.389}_{-0.369}$	$13.908^{+0.399}_{-0.258}$	$14.408^{+0.369}_{-0.4}$	$15.338^{+0.342}_{-0.393}$
$\sigma_8$	$0.77^{+0.121}_{-0.109}$	$1.012^{+0.056}_{-0.081}$	$0.919^{+0.07}_{-0.075}$	$0.862^{+0.06}_{-0.063}$
$\tau$	$0.095^{+0.072}_{-0.074}$	$0.047^{+0.037}_{-0.041}$	$0.079^{+0.023}_{-0.044}$	$0.081^{+0.024}_{-0.041}$

**Table:** Most likely parameter values with  $1 \sigma$  errors for the various COSMOMC Runs

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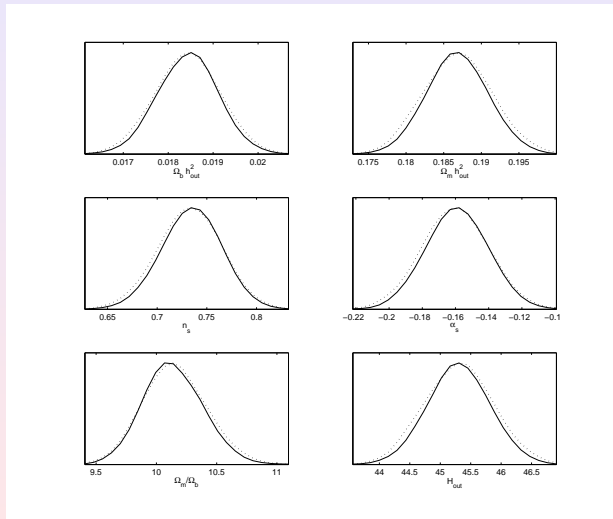


Figure: likelihoods to WMAP 3-yr for the run “EdS with  $\alpha_s$ ”

# Parameter likelihood

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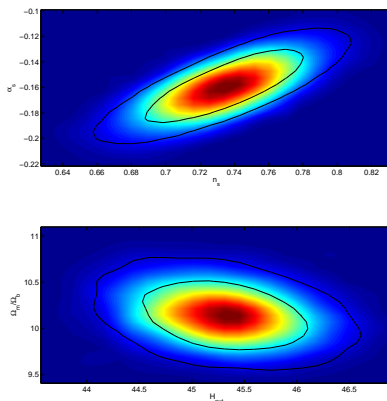
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**Figure:** Contour likelihood plots to WMAP 3-yr for the run “EdS with  $\alpha_s$ ”

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# Is this compatible with local $h$ ?

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- A crucial point: we have
  - a low  $h_{\text{out}}$
  - a constraint on  $\mathcal{J} = h/h_{\text{out}}$

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Conclusions

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  - a low  $h_{\text{out}}$
  - a constraint on  $\mathcal{J} = h/h_{\text{out}}$
- We get a constraint on  $h$



# Is this compatible with local $h$ ?

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- A crucial point: we have
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# Is this compatible with local $h$ ?

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- A crucial point: we have
  - a low  $h_{\text{out}}$
  - a constraint on  $\mathcal{J} = h/h_{\text{out}}$
- We get a constraint on  $h$ . Compatible with local observations?
  - $h = 0.72 \pm 0.08$  from HST (W. L. Freedman *et al.*, *Astrophys. J.* **553**, 47 (2001) )
  - $h = 0.62 \pm 0.01 \pm 0.05$  from HST with corrected Cepheids (A. Sandage *et al.*, *Astrophys. J.* **653**, 843 (2006))
  - $h = 0.59 \pm 0.04$  from Supernovae (Parodi, Saha, Sandage and Tammann, arXiv:astro-ph/0004063. )
  - $h = 0.54_{+0.04}^{-0.03}$  SZ effect ( $z \approx 1$ ) (E. D. Reese *et al.* *Astrophys. J.* **581**, 53 (2002) )

# Parameter Contours

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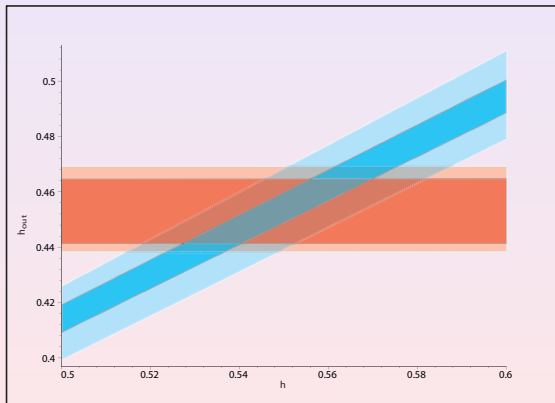


Figure: 1- $\sigma$  and 2- $\sigma$  Contour plots for  $h$  vs.  $h_{out}$ .

# Summarizing the constraints

Void vs Dark  
Energy

At 95% C.L. we have (for  $L \approx 250/h$  Mpc) :

$$\bullet \quad 1.17 \leq \mathcal{J} \leq 1.25 \Rightarrow 0.42 \leq |\delta_0| \leq 0.58$$

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(but note that the average  $\sqrt{\langle \delta^2 \rangle}$  is smaller)

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- $0.44 \leq h_{\text{out}} \leq 0.47$

- $0.51 \leq h \leq 0.59$

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### Conclusions

- The same happens for the model with a bump in the primordial spectrum (S. Sarkar et al. '03 and '07)
- The bump is at a scale of about  $100/h$  Mpc



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 $\Rightarrow$  Combine with the Minimal Void scenario

# $h$ in the Bump model

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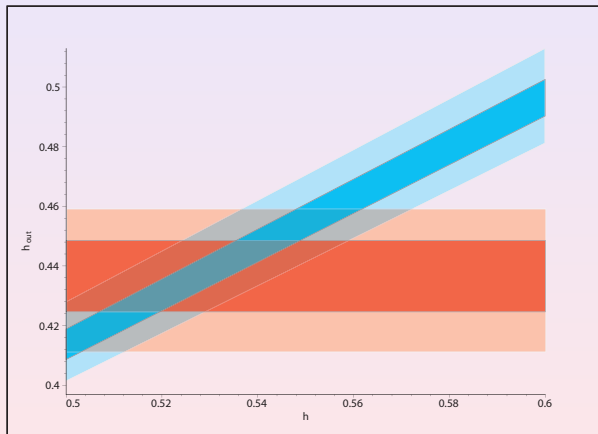


Figure: 1- $\sigma$  and 2- $\sigma$  Contour plots for  $h$  vs.  $h_{out}$ .

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- The position of the peak measures the ratio of the sound horizon at recombination **vs.** angular distance at  $z = 0.35$

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- It constrains two quantities:  $\Omega_m h^2$  and  $D_A(0.35)$
- But it also depends on the spectral index  $n_s$ :

$$D_V = 1370 \pm 64 \quad \text{and} \quad \Omega_m h^2 = 0.130 (n_s/0.98)^{-1.2} \pm 0.011$$



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$$D_V = 1370 \pm 64 \text{ and } \Omega_m h^2 = 0.130 (n_s/0.98)^{-1.2} \pm 0.011$$

- Caveat:
  - Constraints are derived *using*  $\Lambda$ CDM

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Conclusions

- Using  $n_s \sim 0.73$  the constraint is:

$$\Omega_m h_{\text{out}}^2 = 0.185 \pm 0.011, \quad (6.9)$$

- It agrees with our value  $(0.205 \pm 0.01)$  within  $2\sigma$ .

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- On the other hand:

$$\begin{array}{ll} D_A(0.35) = 1375 \text{ Mpc} & \text{for } \Lambda\text{CDM} \\ D_A(0.35) = 1850 \text{ Mpc} & \text{for EdS with } h_{\text{out}} \sim 0.45, \end{array}$$

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$$D_A(0.35) = 1850 \text{ Mpc} \quad \text{for EdS with } h_{\text{out}} \sim 0.45,$$

- **Not** consistent with Eisenstein et al., 2005:

$$D_V(0.35) = 1370 \pm 64 \text{ Mpc},$$

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- The problem is the low value of  $h_{\text{out}}$  from CMB!  
( $h_{\text{out}} \sim 0.56$  would work...)

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- Possible ways out:

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  - *Gpc* Void (Alnes et al., Garcia-Bellido & Haugboelle).  
Fits well, but analysis with full CMB not done yet. It can  
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Fits well, but analysis with full CMB not done yet. It can  
also fit  $D(0.35)/D(0.2)$  (Percival et al.)
  - Fit CMB with higher  $h$   
(**Non-compensated** Void?)



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- It is possible to look (Gaztanaga et al.'08) for the BAO scale only for the radial direction as  $\Delta z$  (model-independent)

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Conclusions

- It is possible to look (Gaztanaga et al.'08) for the BAO scale only for the radial direction as  $\Delta z$  (model-independent)
- Zibin, Moss & Scott '08: it does not fit (*Gpc Void*) together with full CMB (which they fit with very low  $h$  and non-compensated Void)
- Garcia-Bellido & Haugboelle '08: it fits as well as  $\Lambda$ CDM(*Gpc Void*), but only first peak location and SN Union (no full CMB).

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- How much Observer can be off-center?

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- How much Observer can be off-center?
- Observer at Distance  $d_O$
- $\frac{\delta T}{T} \sim v_O \sim \dot{d}_O$
- CMB dipole  $\leq 10^{-3}$  if  $d_O \sim 15 - 20 \text{ Mpc}$  (Tomita et al., Alnes et al.'06)

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- Higher multipoles go as higher powers of  $v_O$ : negligible<sup>10</sup>.
- Bulk dipole of the same size of our dipole 600km/s (Kashlinsky et al. '08: 600 – 1000km/s)

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<sup>10</sup>Alnes et al. '06

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- All objects inside the Void have some peculiar velocity

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- All objects inside the Void have some peculiar velocity
- This gives rise to  $\frac{\delta T}{T} \sim \frac{v}{c}$  and spectrum distortions (kinetic SZ effect)
- Goodman '95:  $v/c \lesssim 0.01$  (at  $z \sim 0.2$ )
- Caldwell-Stebbins '07-'08: rule out Voids with  $z_b > 0.9$

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- Garcia-Bellido & Haugboelle: using 9 clusters ( $0.2 \leq z \leq 0.6$ ) with detection of spectral distortion one finds:  
 $\bar{v} = 320 \text{ km/sec}$  and  $\sigma = 1600 \text{ km/sec}$  ( $\sigma$  expected is only about  $400 \text{ km/sec!}$ )
- Exclude  $L > 1.5 \text{ Gpc}$ , with  $\Omega_{IN} = 0.23$ .

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- Exclude  $L > 1.5 \text{ Gpc}$ , with  $\Omega_{IN} = 0.23$ .
- But Kashlinsky et al. measure high  $\frac{v}{c} \sim 1000 \text{ km/sec}$  on  $300 \text{ Mpc}/h$  (they *assume* kSZ, but do not see spectral distortions).

# Other ways to test the Copernican principle

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- Clarkson, Bassett and Lu '08: a consistency relation:  
 $\mathcal{C}(z) = 0$  for FLRW, at *all*  $z$ ,

$$\mathcal{C}(z) \equiv 1 + H^2(DD'' - D'^2) + HH'DD'$$

- Uzan, Clarkson and Ellis '08: Time drift of the redshift  
(over 10 years)

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- Similarly the expansion is anisotropic if  $d_O$  nonzero<sup>11</sup>.

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- Two papers claim significant anisotropy in  $H$ :
  - D.Schwarz & Weinhorst '07: in the SNIa dataset ( $> 95\% C.L.$ )
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- In addition this should be correlated with CMB dipole
- Also to be explored: non-sphericity of Void

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- It is well-known that linear  $\Phi$  constant in Matter Dominated Universe

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- If they evolve instead  $\Rightarrow$  photon feels  $\Delta\Phi$  inside structures  $\Rightarrow$  additional secondary CMB anisotropy
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- Can we get this in our scenario?
  - Inside the Void
  - If there is curvature
  - If there are other big Voids in the sky  $\Rightarrow$  nonlinear evolution of  $\Phi$
- Effect of order  $(L/r_{\text{hor}})^3 \sim \mathcal{O}(10^{-5})$

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- A Void of at least  $L \sim 200 - 250 \text{ Mpc}/h$  scale consistent with WMAP and SNIa (Riess data), and local  $h$
- $\delta$  quite large ( $\sim 0.4$ )  
Incompatible with the expected value ( $\delta \sim 0.04$ ).

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- But some observations seem to indicate such structures (need for more observations)
- Need for **larger Void** to fit Union data ( $L \gtrsim 500 \text{ Mpc}/h$ )
- More data will discriminate (especially **SDSS-II** for Supernovae)

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- Checking with curved models (*work in progress*)
- Non-compensated Voids? (monopole  $T_0 = 2.73K$  gets large correction) Zibin-Moss-Scott '08

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- Observer has to sit **near the center** (10 – 20Mpc in radial position)

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- ISW effect to be included
- Check if the higher  $\Omega_m/\Omega_b$  is compatible with other data

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## Void:

- Many observations to reconcile together
- Before considering it as a valid alternative to  $\Lambda$ CDM
- More work to be done (and more data will soon discriminate)

# Can an Inhomogeneous Universe mimic Dark Energy?

Alessio Notari <sup>12</sup>

CERN

Dec. 2008 / Workshop @ KEK

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<sup>12</sup>In collaboration with:

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