

Probing the Central Engine of Long Gamma-Ray Bursts and Hypernovae with Gravitational Waves

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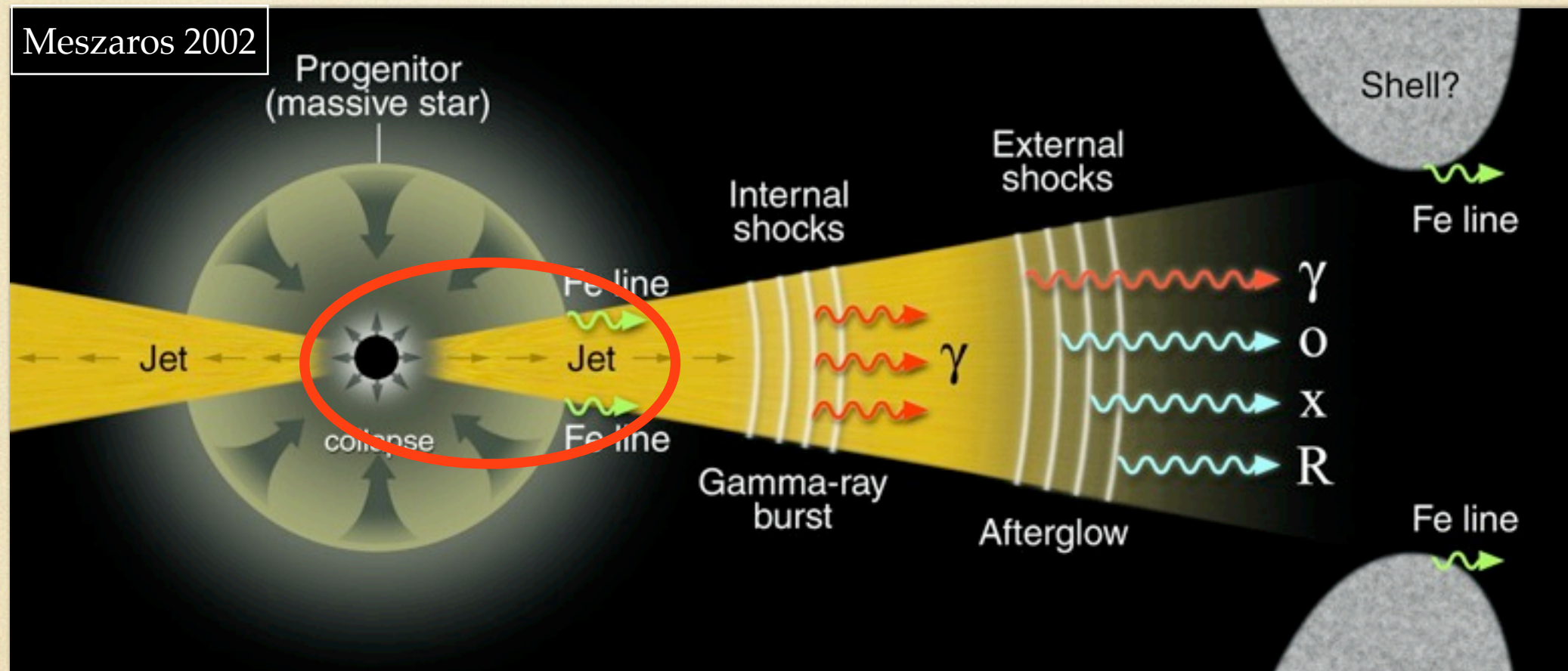
Collaboration with K. Murase (Kyoto U.)

arXiv:0906.3833

Contents

- Introduction
- Collapsar model
- Gravitational wave from collapsar

Gamma-ray bursts



- The most luminous explosions in the Universe
- Small amount of matter accelerated to ultrarelativistic speeds and collimated in a jet
- In many of longer lasting events the total energy in γ rays $\sim 10^{51}$ ergs.
- The required energy for a jet: **$E \sim 10^{52}$ ergs**

GRBs and SNe

GRB \Leftrightarrow SN association

GRB 980425 / SN 1998bw: ($z=0.0085$)

GRB 030329 / SN 2003dh: ($z=0.1687$)

GRB 031203 / SN 2003lw: ($z=0.1055$)

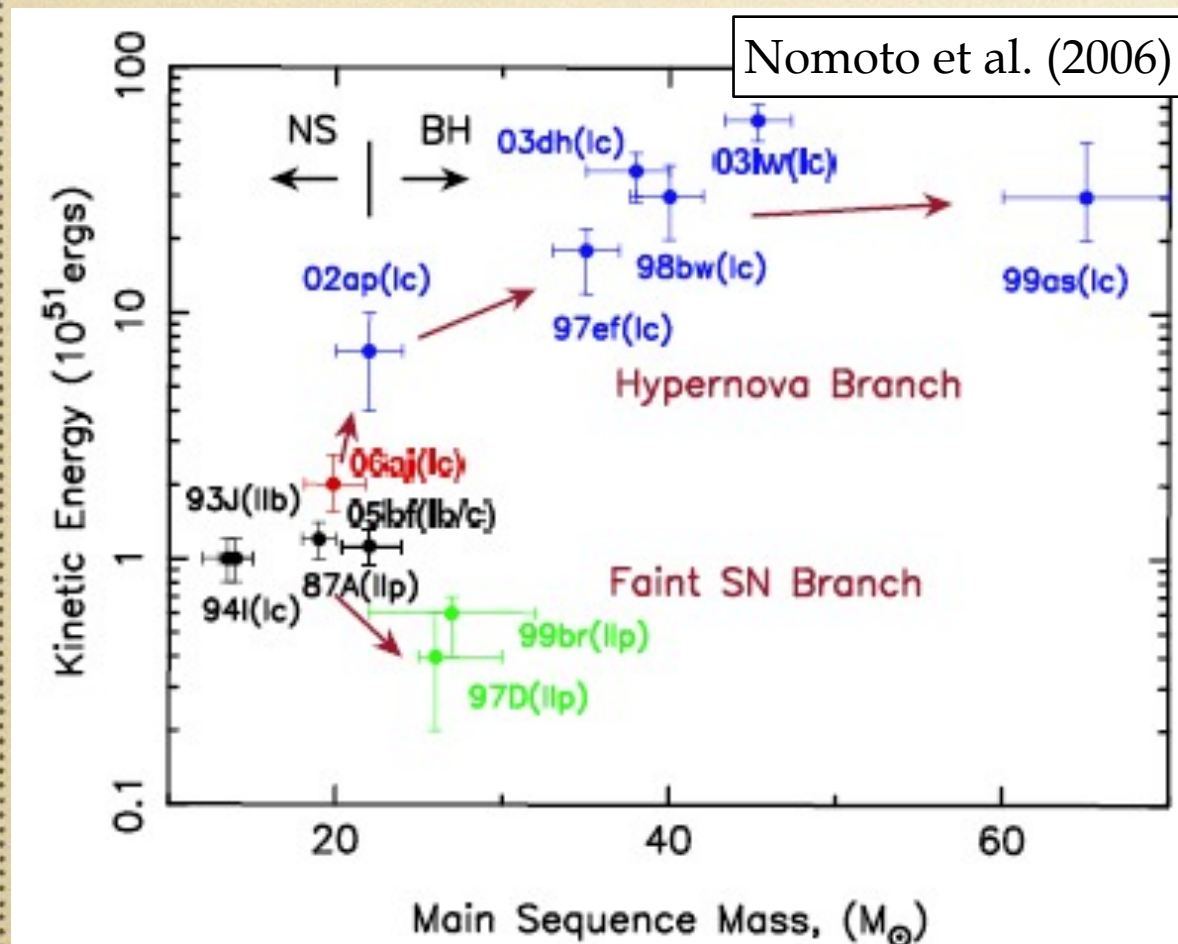
XRF 060218 / SN 2006aj: ($z=0.0335$)

GRB 081007 / SN 2008hw: ($z=0.53$)

Observations of GRB suggest that some GRBs are connected with some kind of SNe.

SNe which associate with GRB are “Hypernovae” (HNe) with explosion energy, $E_{\text{exp}} \sim 10^{52}$ ergs.

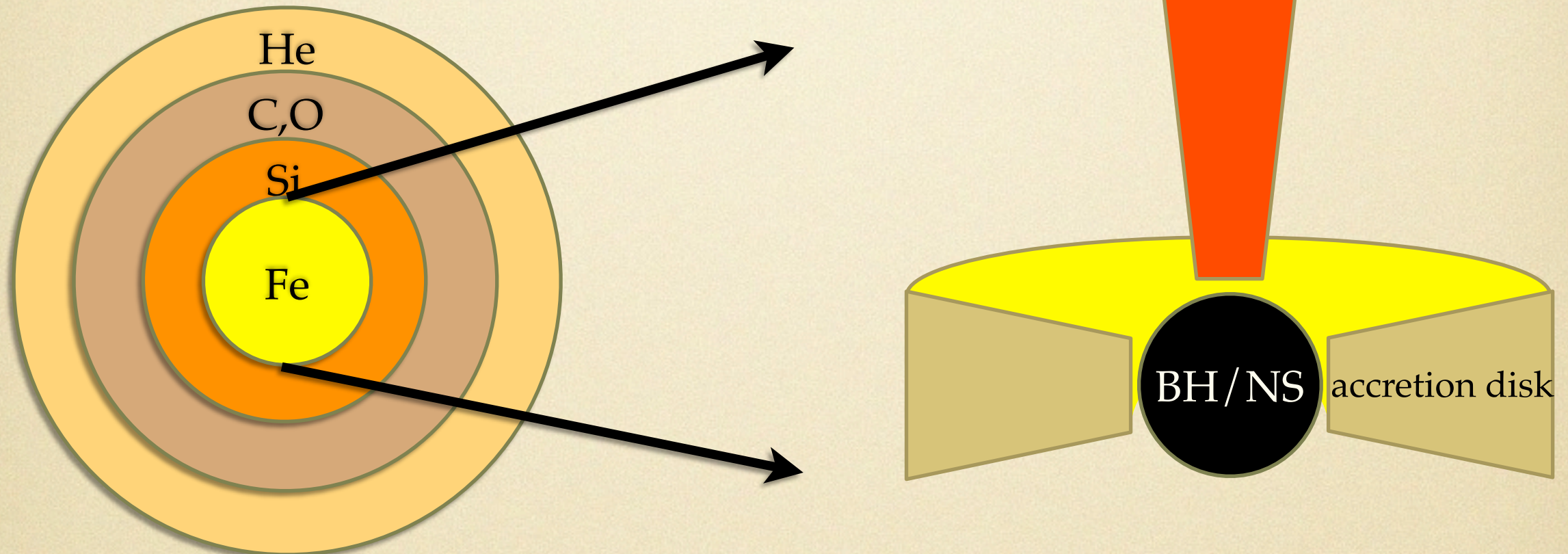
The central engine of GRBs is required to supply such an enormous explosion energy of GRBs/HNe.



Collapsar model

The most promising model of long GRBs

Woosley 93, Paczynski 98, MacFadyen+ 99

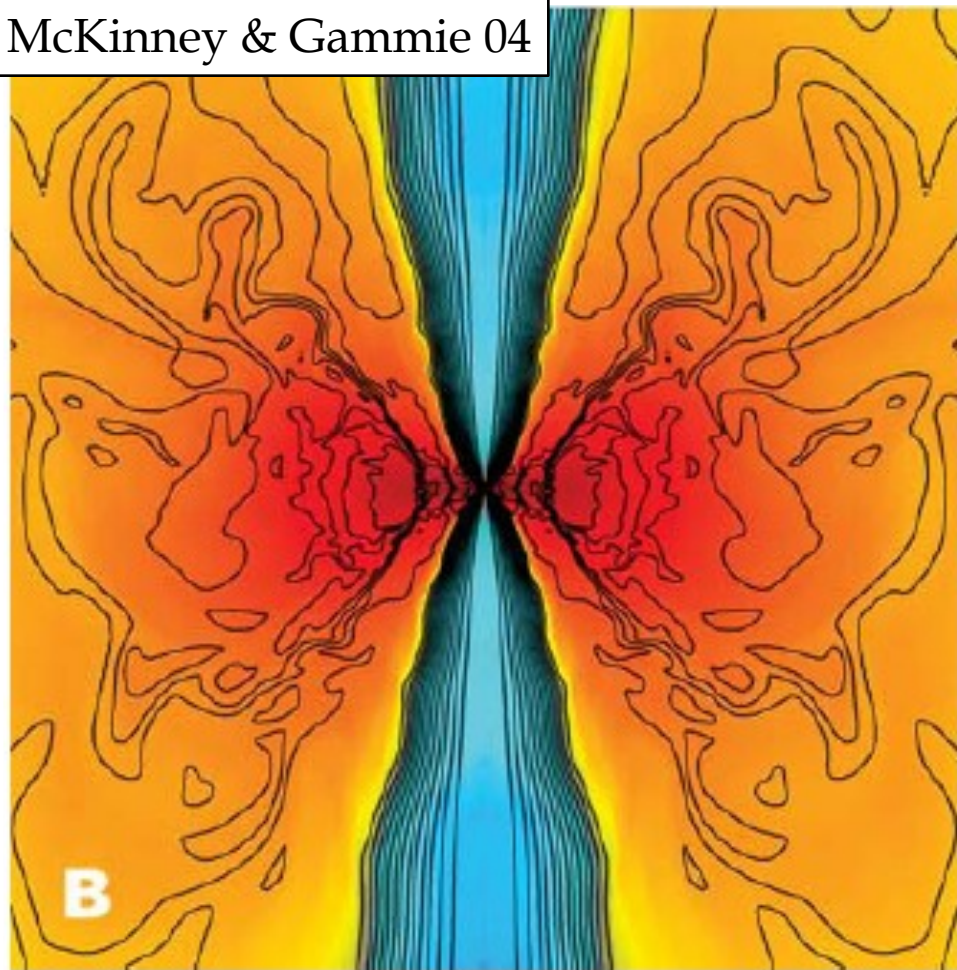


- Iron core of Massive star BH/NS + disk system
- Relativistic jet formation by some kind of “mechanism” ➡ GRB
- Candidates of “mechanism”:
magnetic force and **neutrino annihilation**

Magneto-driven jet

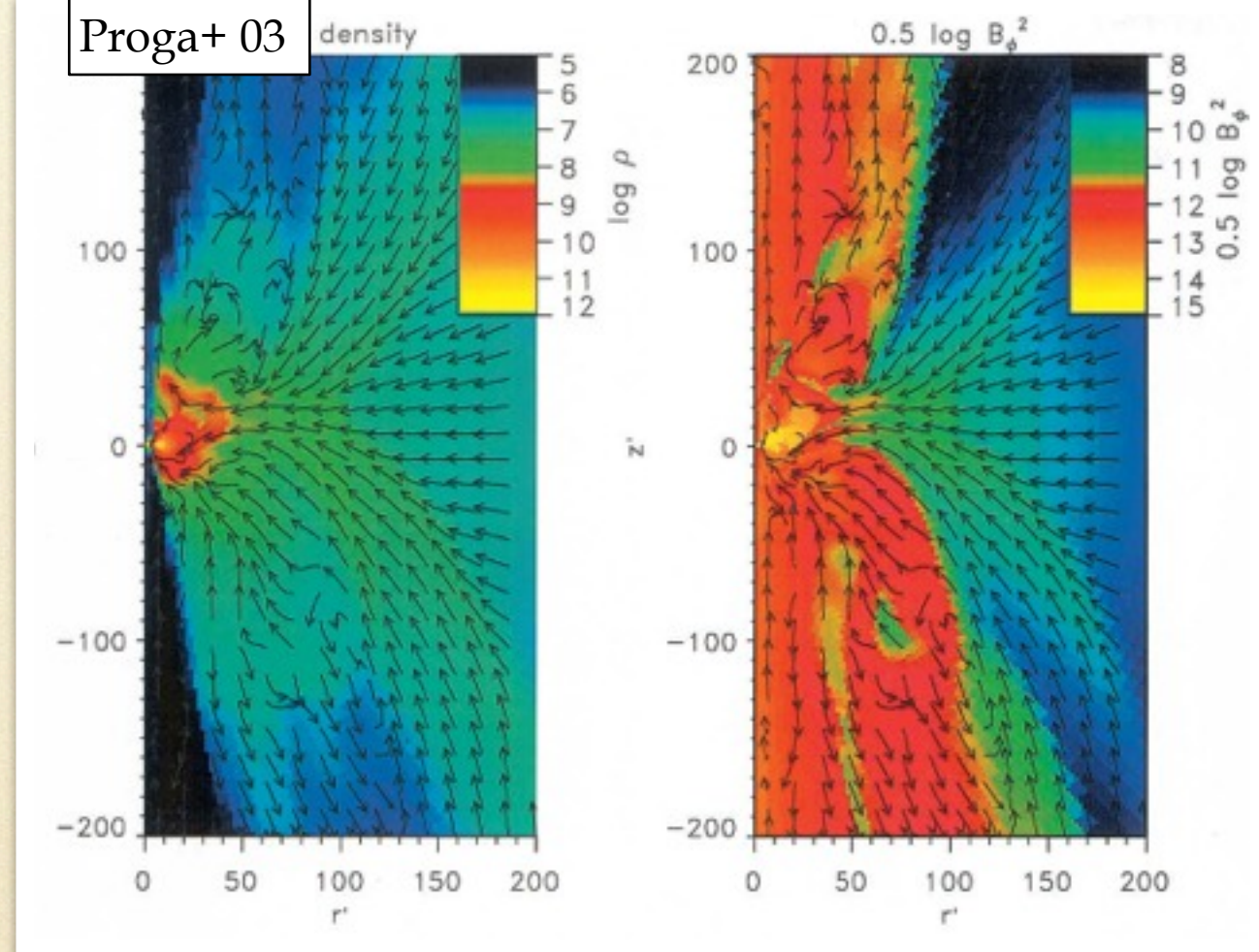
Blandford-Znajek process

McKinney & Gammie 04



MHD process

Proga+ 03



rotation energy of a BH

☞ Poynting flux ☞ **jet production**

$$P_{\text{BZ}} \sim 10^{51} \tilde{a}^2 \left(\frac{M_{\text{BH}}}{3M_{\odot}} \right)^2 \left(\frac{B}{10^{15} \text{G}} \right)^2 \text{ ergs s}^{-1} \quad \text{Lee+ 00}$$

$$E_{\text{rot}} = 5.4 \times 10^{54} f(\tilde{a}) \left(\frac{M}{3M_{\odot}} \right) \text{ ergs} \quad f(\tilde{a}) = 1 - \sqrt{\frac{1}{2}[1 + \sqrt{1 - \tilde{a}^2}]}$$

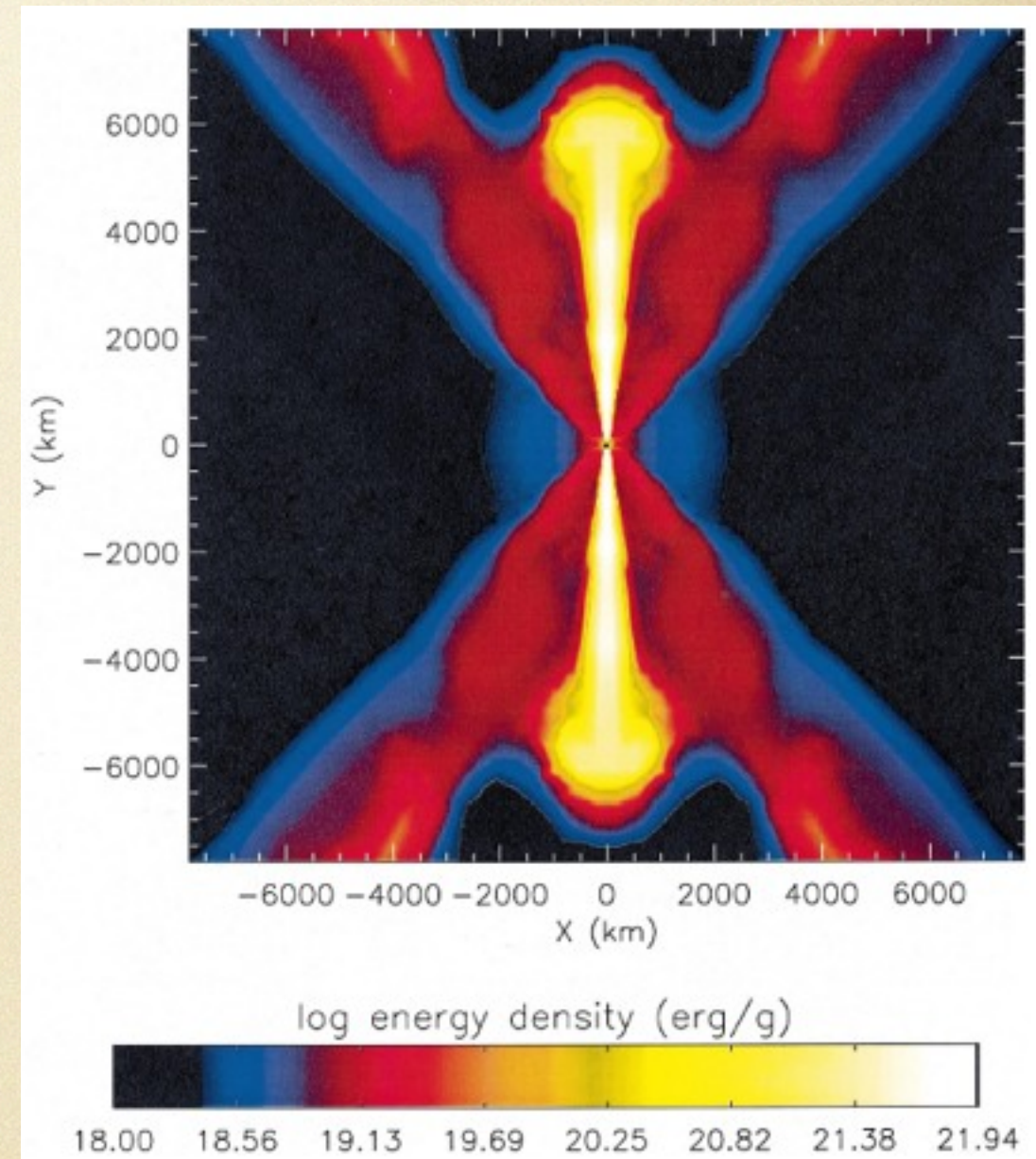
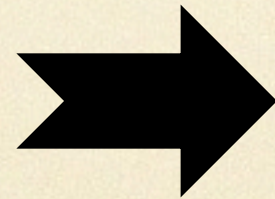
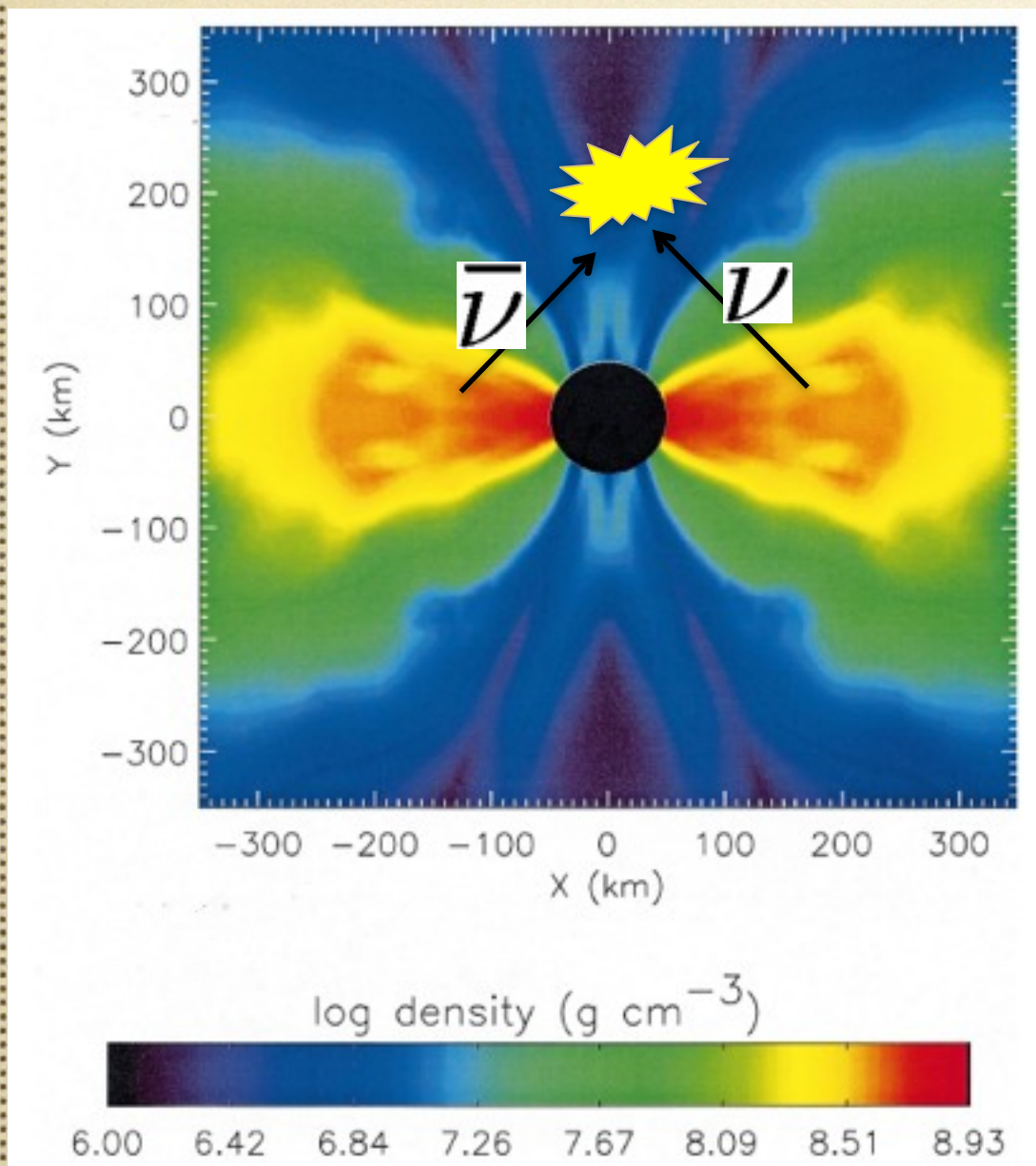
binding energy of an accretion disk

☞ Poynting flux ☞ **jet production**

$$P_{\text{MHD}} \sim P_{\text{BZ}}$$

Neutrino-driven jet

McFadyen & Woosley 99

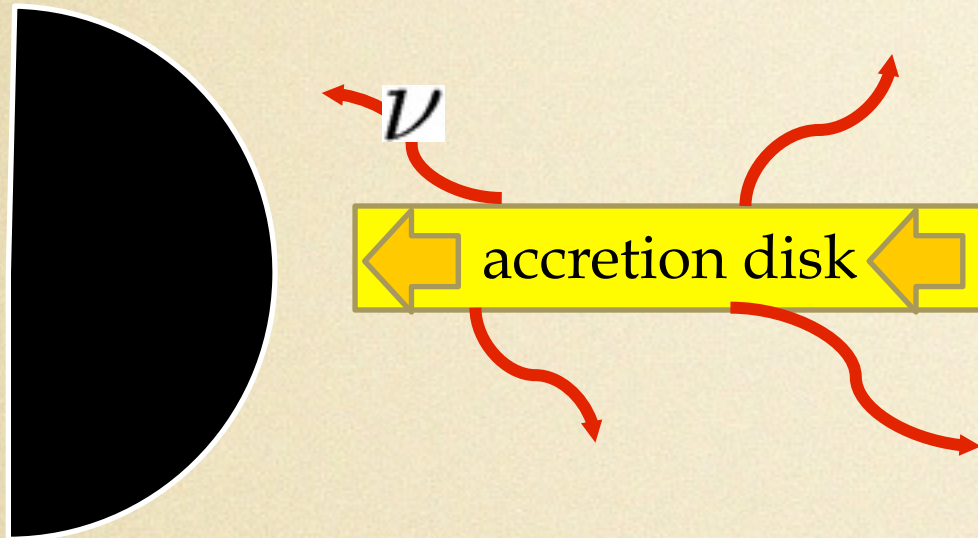


Neutrino pairs are generated in the hot disk

☞ Impact each other ☞ Energy deposition at rotational axis ☞ **jet production**

NDAF

Neutrino Dominated Accretion Flow

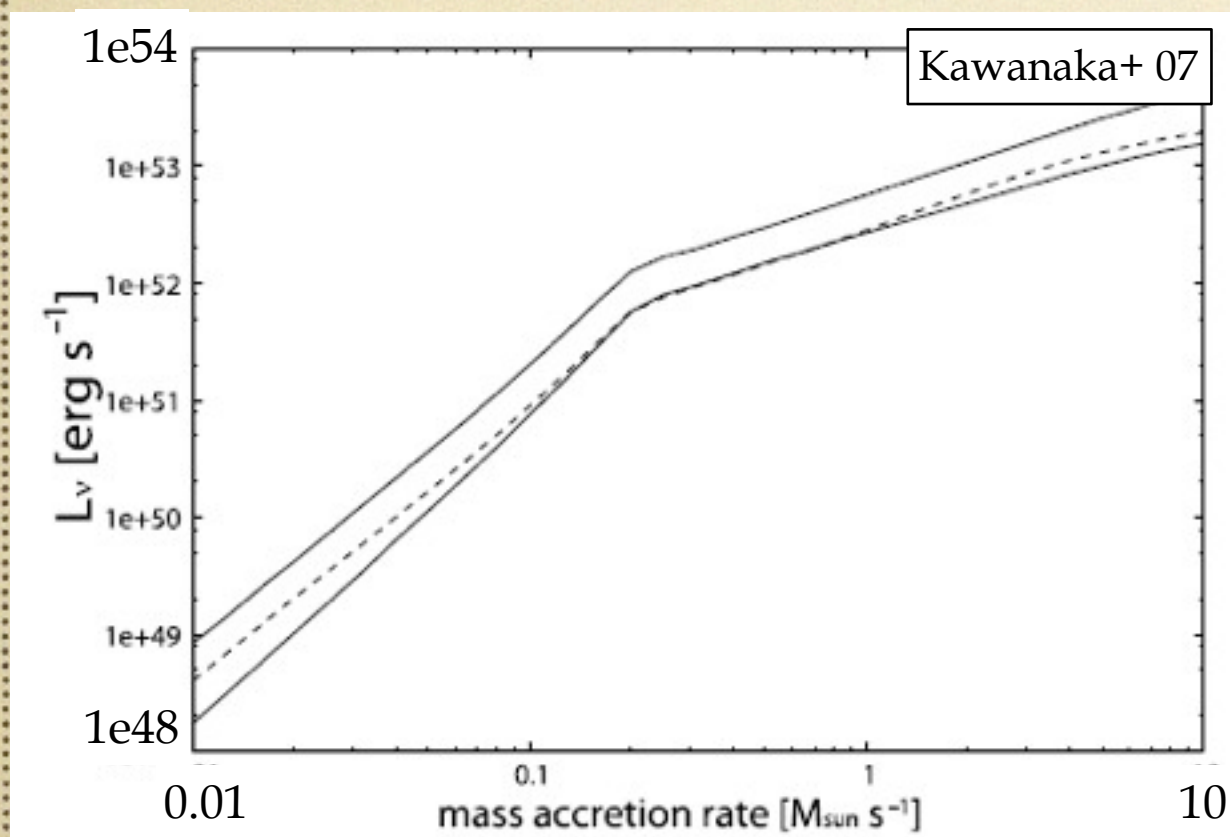


- Model for hyper accretion disk

- The neutrino emission is the dominant source of cooling.

(These neutrinos are suggested for the source of a jet)

- Depending on the accretion rate, the neutrino luminosity, L , becomes as large as 10^{53} erg/sec!



NDAF

Popham+ 99

TABLE 3
NEUTRINO ANNIHILATION EFFICIENCY

\dot{M} ($M_{\odot} \text{ s}^{-1}$)	α	a	M (M_{\odot})	L_{ν} ($10^{51} \text{ ergs s}^{-1}$)	$L_{\nu\bar{\nu}}$ ($10^{51} \text{ ergs s}^{-1}$)	Efficiency (%)
0.01	0.1	0	3	0.015	3.9×10^{-8}	0.0003
0.01	0.03	0	3	0.089	2.9×10^{-7}	0.0003
0.01	0.01	0	3	0.650	9.0×10^{-6}	0.001
0.01	0.1	0.5	3	0.036	5.9×10^{-7}	0.002
0.01	0.01	0	10	0.049	6.4×10^{-9}	10^{-5}
0.05	0.1	0.5	3	1.65	1.8×10^{-3}	0.11
0.1	0.1	0	3	3.35	3.0×10^{-3}	0.09
0.1	0.03	0	3	6.96	1.7×10^{-3}	0.02
0.1	0.01	0	3	6.15	8.0×10^{-4}	0.01
0.1	0.1	0.5	3	8.03	0.039	0.5
0.1	0.1	0.95	3	46.4	2.0	4.2
0.1	0.1	0.95	6	26.2	0.79	3.0
1.0	0.1	0	3	86.3	0.56	0.6
1.0	0.1	0.5	3	142	3.5	2.5
10.0 ^a	0.1	0	3	(781)	(200)	(26)
10.0 ^a	0.1	0.5	3	(1280)	(820)	(64)

^a The assumption that the neutrinos are optically thin breaks down for accretion rates of $10 M_{\odot} \text{ s}^{-1}$ and above. The neutrino annihilation luminosities and energies listed for these high-accretion simulations are upper limits.

The required energy for producing GRB jet is $\sim 10^{51} \text{ erg/sec}$.

In order to achieve this, the total neutrino luminosity must be as large as $\sim 10^{53} \text{ erg/sec}$!

(The efficiency of energy conversion is quite low.)

Summary of energy budget

Energy of GRB jet

$$E_{\text{GRB-jet}} = 10^{52} \text{ ergs}$$

Energy available by models of central engine

BZ process

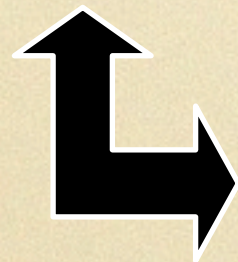
$$P_{\text{BZ}} \sim 10^{51} \tilde{a}^2 \left(\frac{M_{\text{BH}}}{3M_{\odot}} \right)^2 \left(\frac{B}{10^{15} \text{ G}} \right)^2 \text{ ergs s}^{-1}$$

MHD process

$$P_{\text{MHD}} \sim P_{\text{BZ}}$$

ν process

$$P_{\nu\bar{\nu}} \sim 10^{51} \text{ ergs s}^{-1}$$



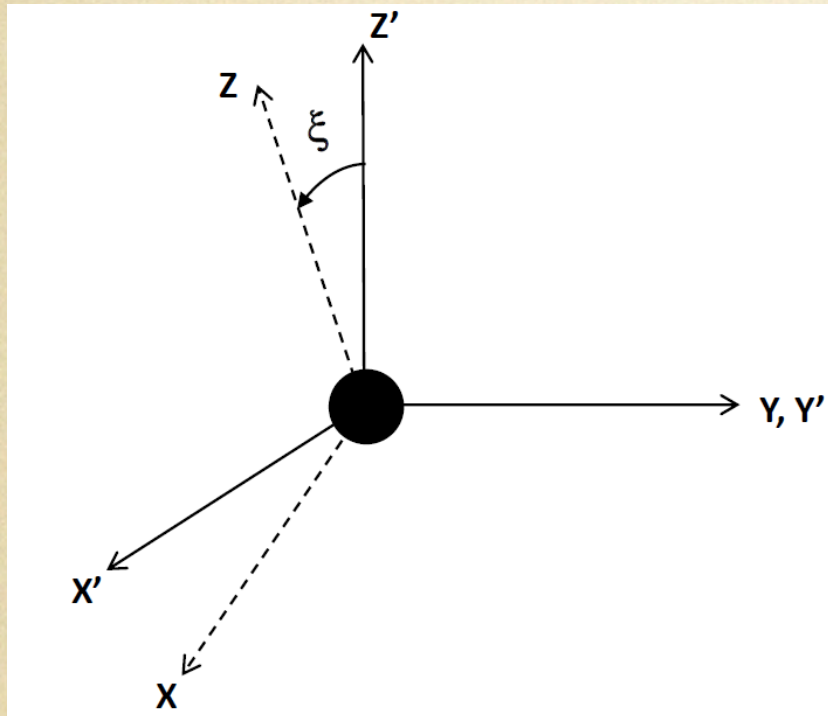
$$L_{\nu} \sim 10^{53} \text{ ergs s}^{-1}$$



Can we constrain this component using GW??

GW from anisotropic neutrino radiation

Epstein 78, Turner 78, Mueller & Janka 97



GW amplitude

$$h_+(t) = \frac{2G}{c^4 R} \int_{-\infty}^{t-R/c} dt' \int_{4\pi} d\Omega' \Psi(\Omega') \frac{dL_\nu}{d\Omega'}(\Omega', t')$$

neutrino luminosity per unit solid angle

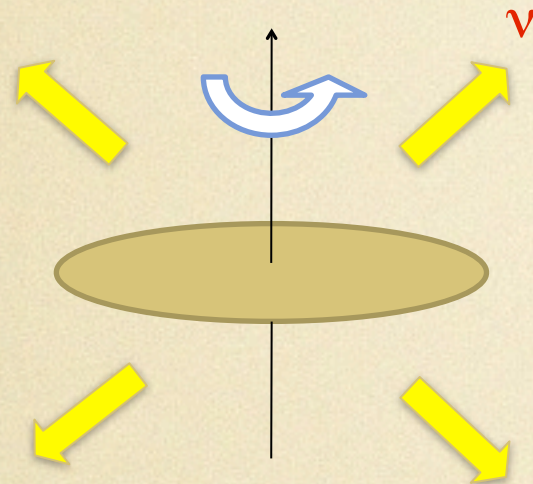
$$\Psi(\Omega') = (1 + \cos \theta' \cos \xi + \sin \theta' \cos \phi' \sin \theta') \frac{(\sin \theta' \cos \phi' \cos \xi - \cos \theta' \sin \xi)^2 - \sin^2 \theta' \sin^2 \phi'}{(\sin \theta' \cos \phi' \cos \xi - \cos \theta' \sin \xi)^2 + \sin^2 \theta' \sin^2 \phi'}$$

Axisymmetric case

$$h_+(t) = \frac{2G}{c^4 R} \int_{-\infty}^{t-R/c} dt' \int_0^\pi \sin \theta' d\theta' \Phi(\theta') \frac{dL_\nu}{d\Omega'}(\theta', t'),$$

$$\Phi(\theta') = \begin{cases} -2\pi [1 + \cos \theta' (2 + \cos \xi)] \tan^2 \left(\frac{\xi}{2} \right) & (\text{for } \theta \geq \xi) \\ -2\pi [1 + \cos \theta' (-2 + \cos \xi)] \cot^2 \left(\frac{\xi}{2} \right) & (\text{for } \theta < \xi) \end{cases}$$

GW from thin disk

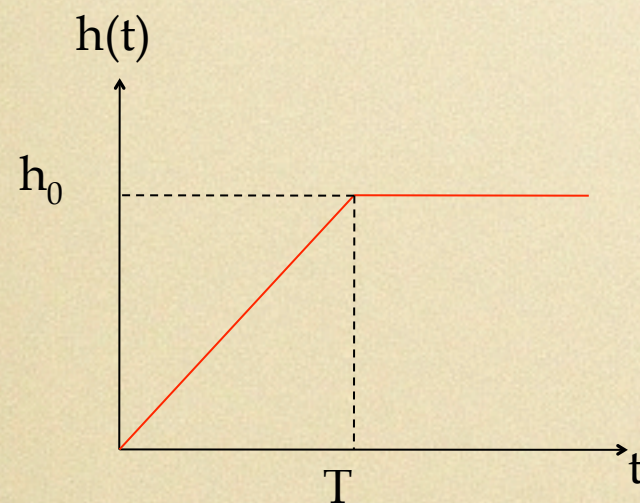


Neutrino emission

$$\frac{dL_\nu}{d\Omega'} = \frac{L_\nu}{2\pi} |\cos \theta|$$

GW amplitude ($\xi=\pi/2$)

$$\begin{aligned} h_+(t) &= \frac{2G}{c^4 R} \int_{-\infty}^{t-R/c} dt' \int_0^\pi d\theta' \Phi(\theta') \frac{dL_\nu}{d\Omega'}(\theta', t') \\ &= \frac{2G}{c^4 R} \int_{-\infty}^{t-R/c} dt' L_\nu(t') \times \int_0^\pi d\theta' (-1 + 2 \cos \theta') \sin \theta' |\cos \theta'| \\ &= \frac{2G}{c^4 R} \left(\frac{1}{3} \right) \int_{-\infty}^{t-R/c} L_\nu(t') dt'. \end{aligned}$$



Final converged value

$$h_0 \sim 1.8 \times 10^{-21} \left(\frac{10 \text{ Mpc}}{R} \right) \left(\frac{E_\nu}{10^{54} \text{ ergs}} \right)$$

R: distance
E_ν: total energy emitted by neutrino

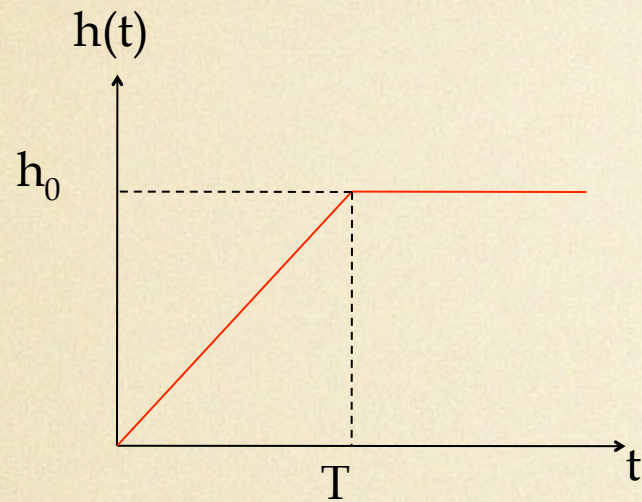
$$\Leftrightarrow E_{\text{tot}} = 3 \times 10^{53} \text{ ergs for ordinary SNe}$$

cf.) GW from jet itself

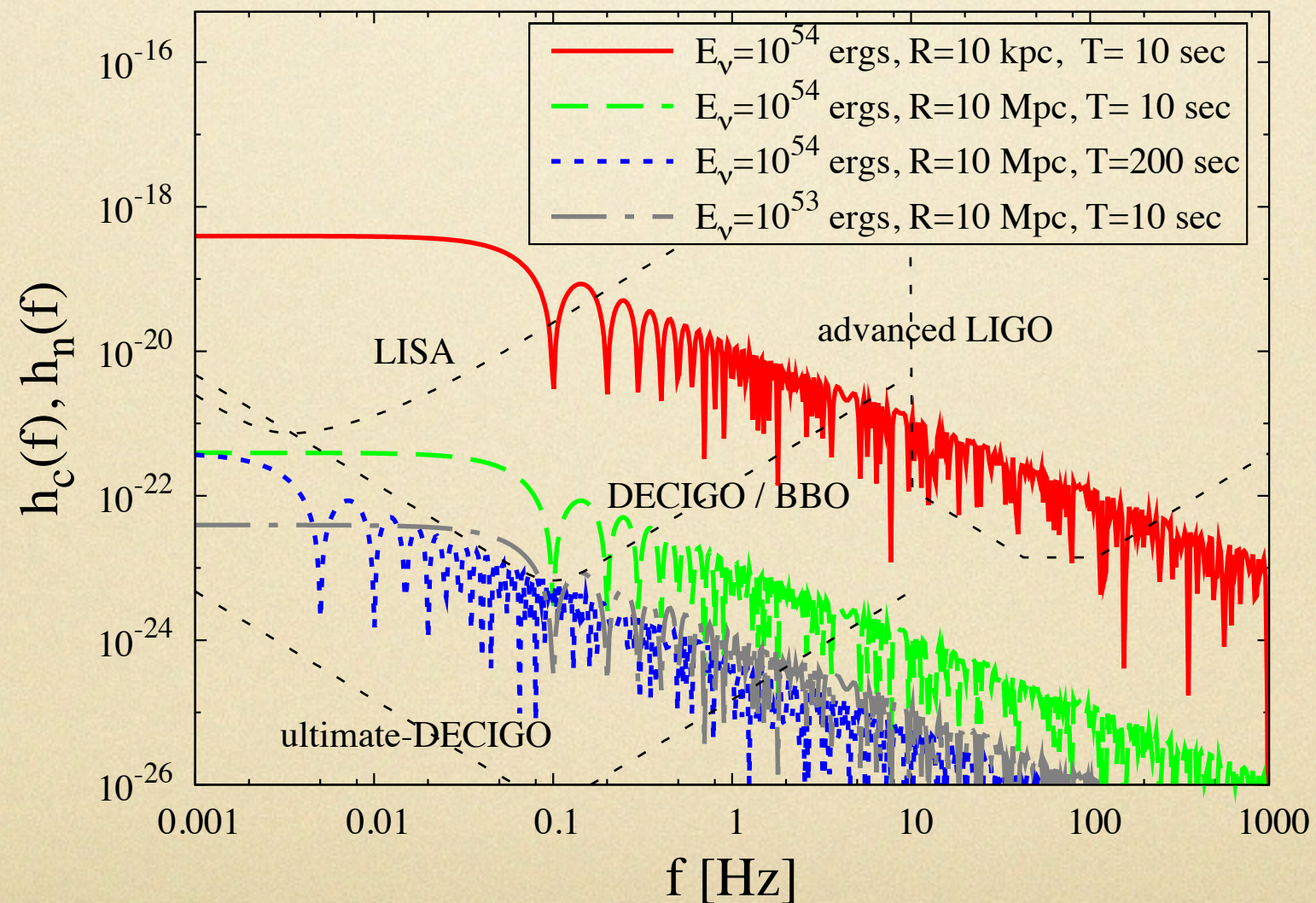
$$h_0 \sim 8.5 \times 10^{-24} \left(\frac{10 \text{ Mpc}}{R} \right) \left(\frac{E_{\text{jet}}}{10^{52} \text{ ergs}} \right)$$

Sago+ 04
Hiramatsu+ 05

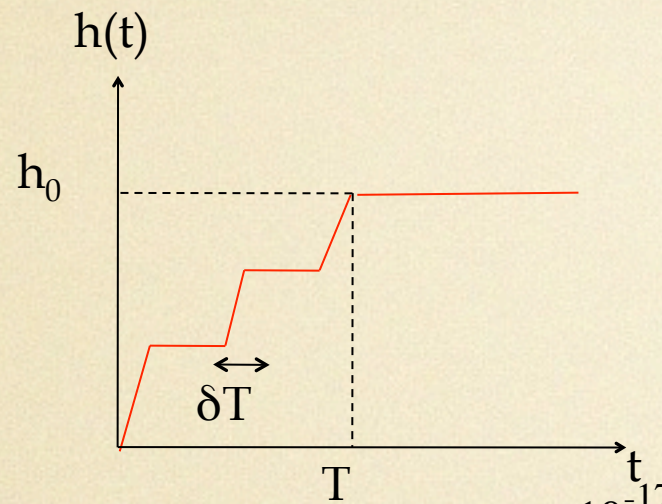
GW spectrum



$$h_c(f) \equiv f |\tilde{h}(f)| = \frac{h_0}{2\pi^2 T f} |\sin(\pi T f)|$$

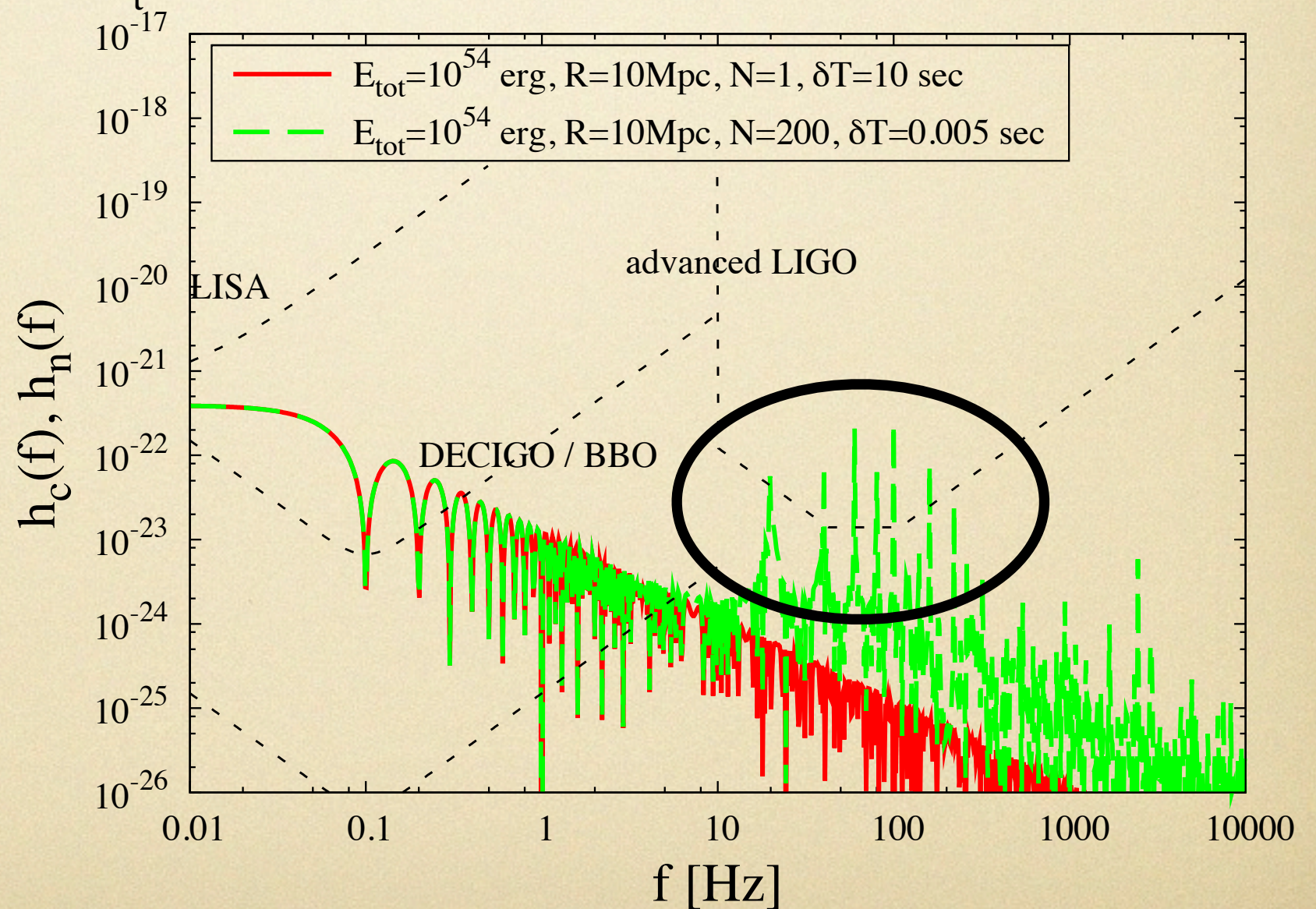


Effect of time variability

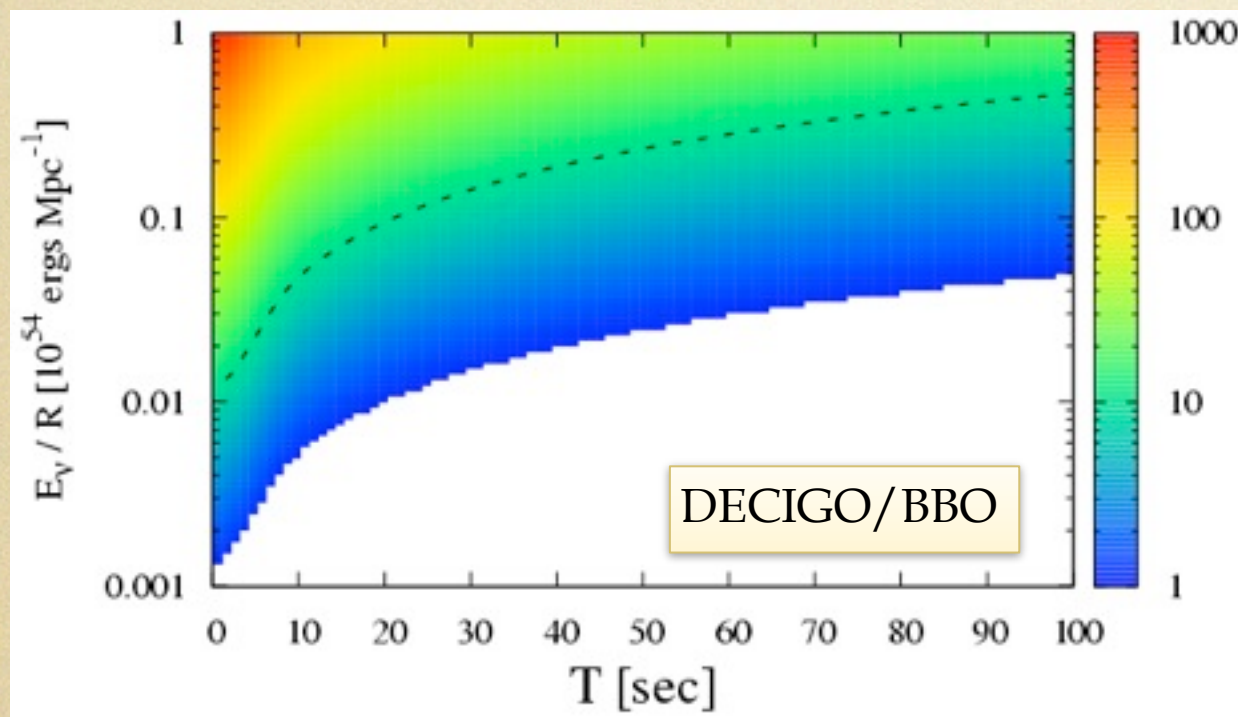
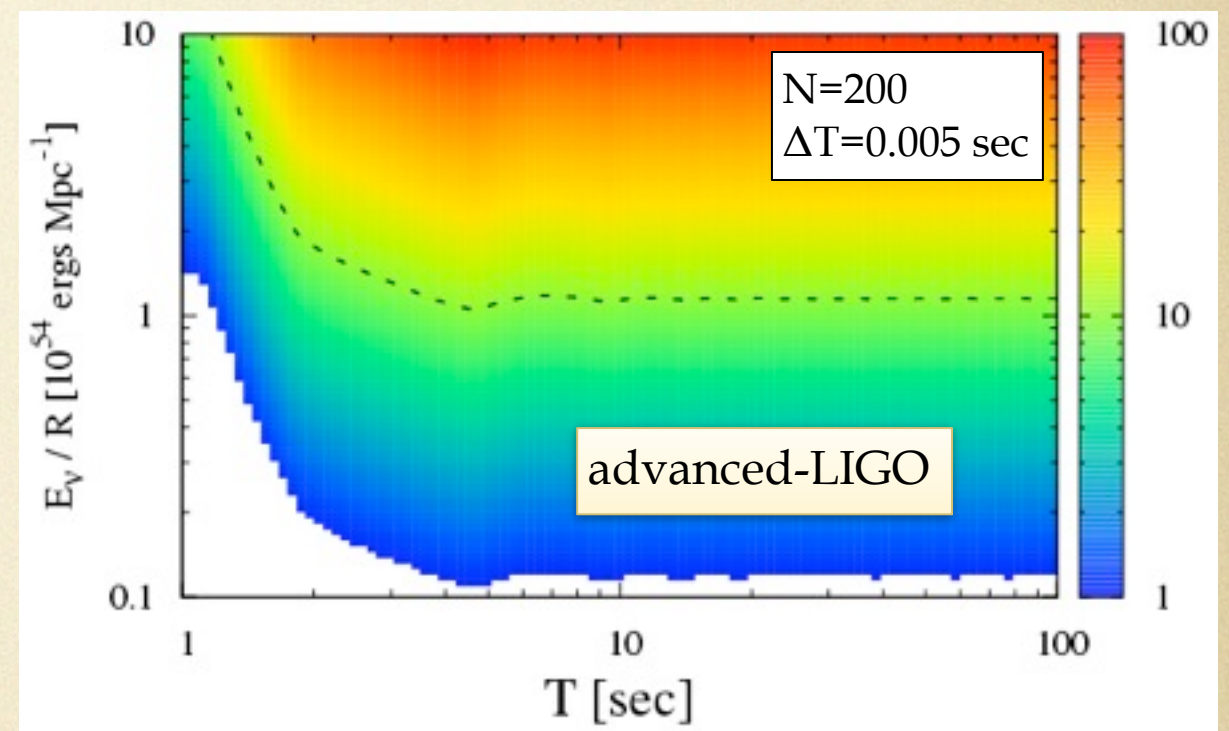
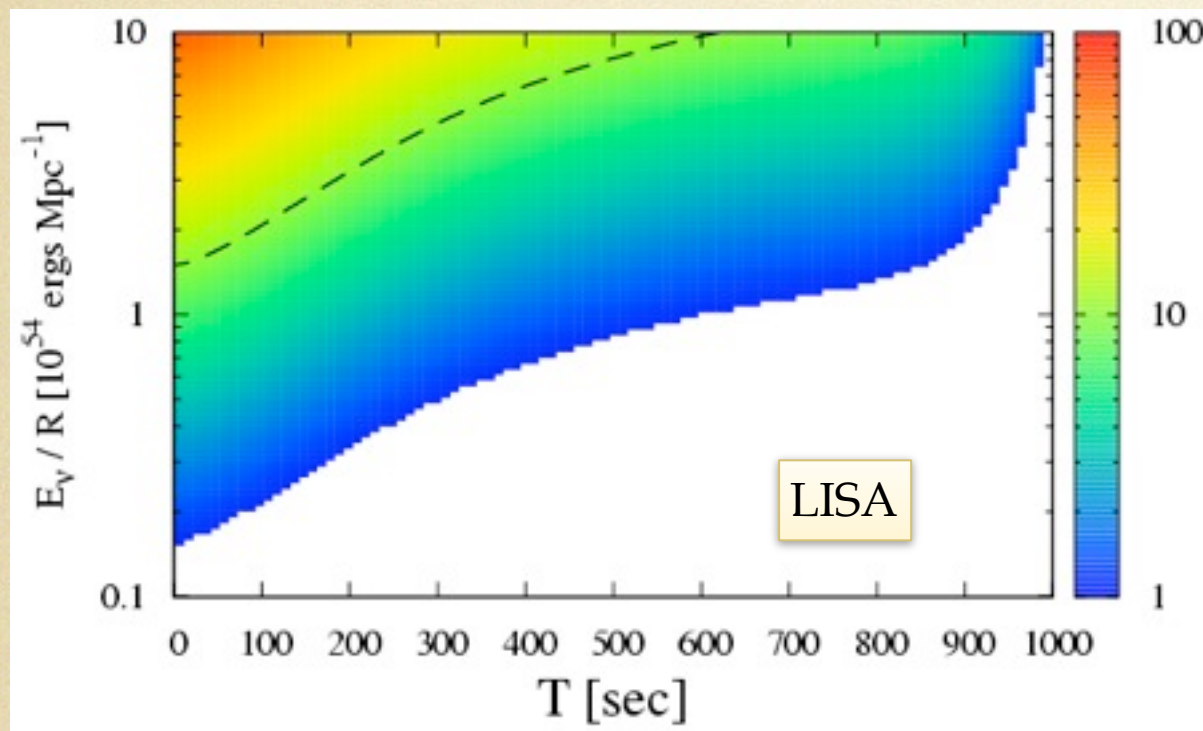


$$h_c(f) = \frac{\sqrt{\beta} h_\infty}{\pi^2 N \delta t f} \left| \frac{\sin(\pi \delta t f) \sin(\pi T f)}{\sin(\pi T f / N)} \right|$$

N: number of bursts



Signal-to-noise ratio



$$\text{SNR} = \sqrt{\int_0^\infty d(\ln f) \frac{h_c(f)^2}{h_n(f)^2}}$$

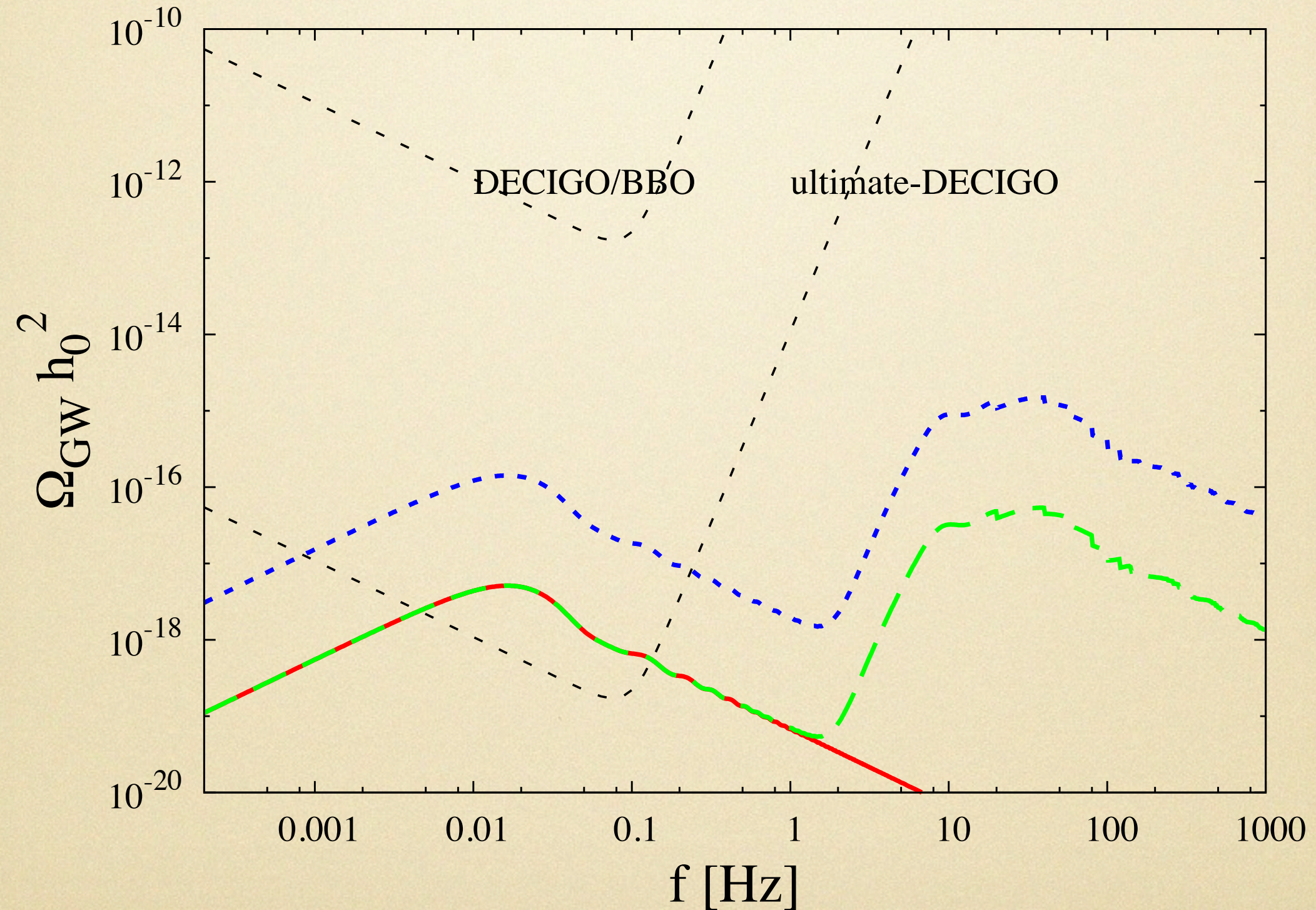
$\text{SNR} > 10 \Leftrightarrow$

- < 1 Mpc for ad.-LIGO (2015)
- < 1 Mpc for LISA (2019)
- < 100 Mpc for DECIGO/BBO (2024)

cf.) GRB rate \Leftrightarrow

$\sim 10 - \sim 1000 \text{ Gpc}^{-3} \text{ yr}^{-1}$ (Guetta+ 07)

Gravitational Wave Background



Summary

Energy of GRB jet

$$E_{\text{GRB-jet}} = 10^{52} \text{ ergs}$$

Energy available by models of central engine

BZ process

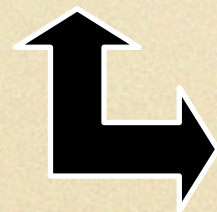
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MHD process

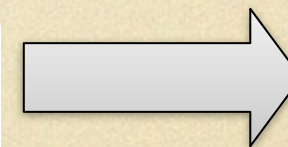
$$P_{\text{MHD}} \sim P_{\text{BZ}}$$

ν process

$$P_{\nu\bar{\nu}} \sim 10^{51} \text{ ergs s}^{-1}$$



$$L_{\nu} \sim 10^{53} \text{ ergs s}^{-1}$$



GW