

Jamming transition of non-spherical particles

[EPJE, 44, 120 \(2021\) or arXiv:2012.07294](#)

Harukuni Ikeda

Department of Physics, Gakushuin University

- **Introduction**
- **Jamming of spherical particles**
- **Jamming of non-spherical particles**

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- **Jamming of non-spherical particles**

Introduction

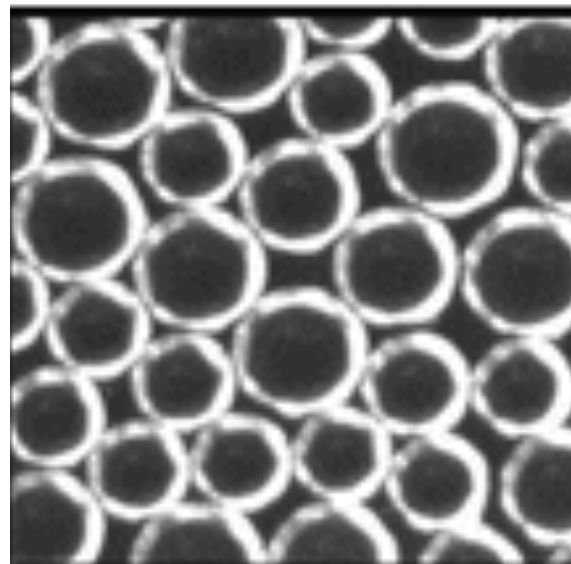
Granular matters are ubiquitous in our daily life!

Assembly of macroscopic particles (0.1mm~)

Thermal fluctuations are completely negligible = Non equilibrium



M&M Candies



Forms



Sand & rock



Grains

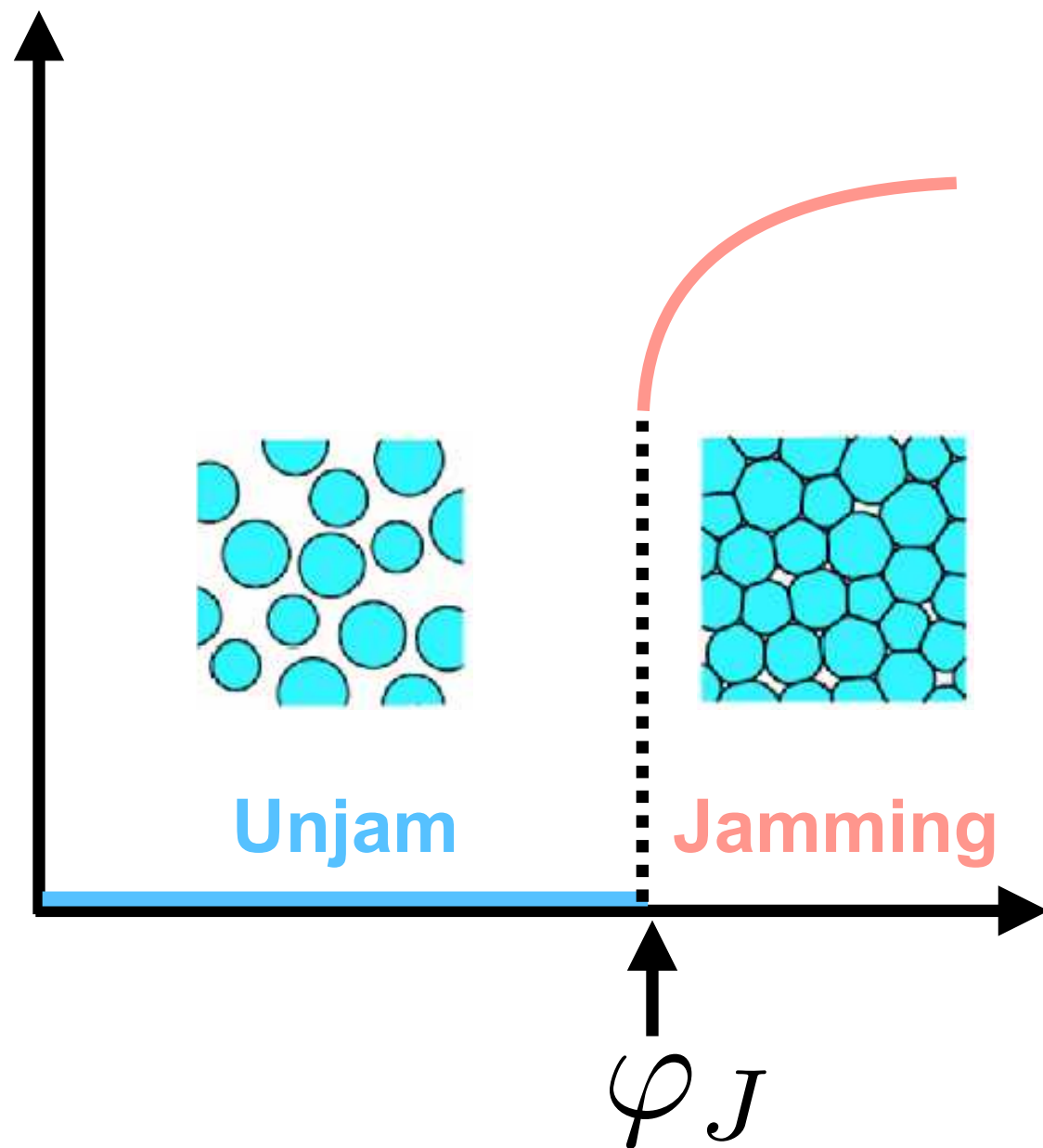


Snow powders

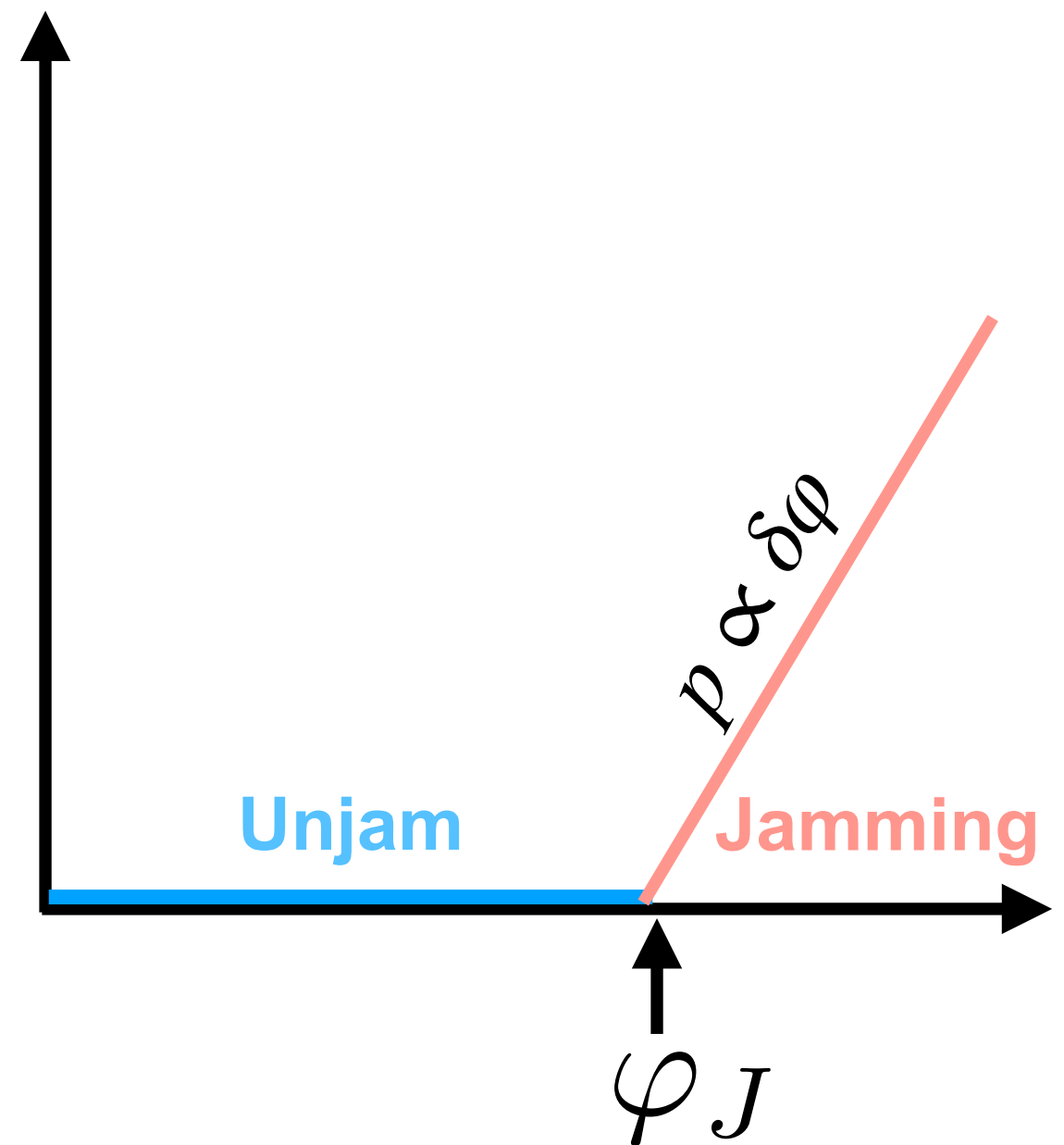
Introduction

What is the jamming transition?

Contact number



Pressure

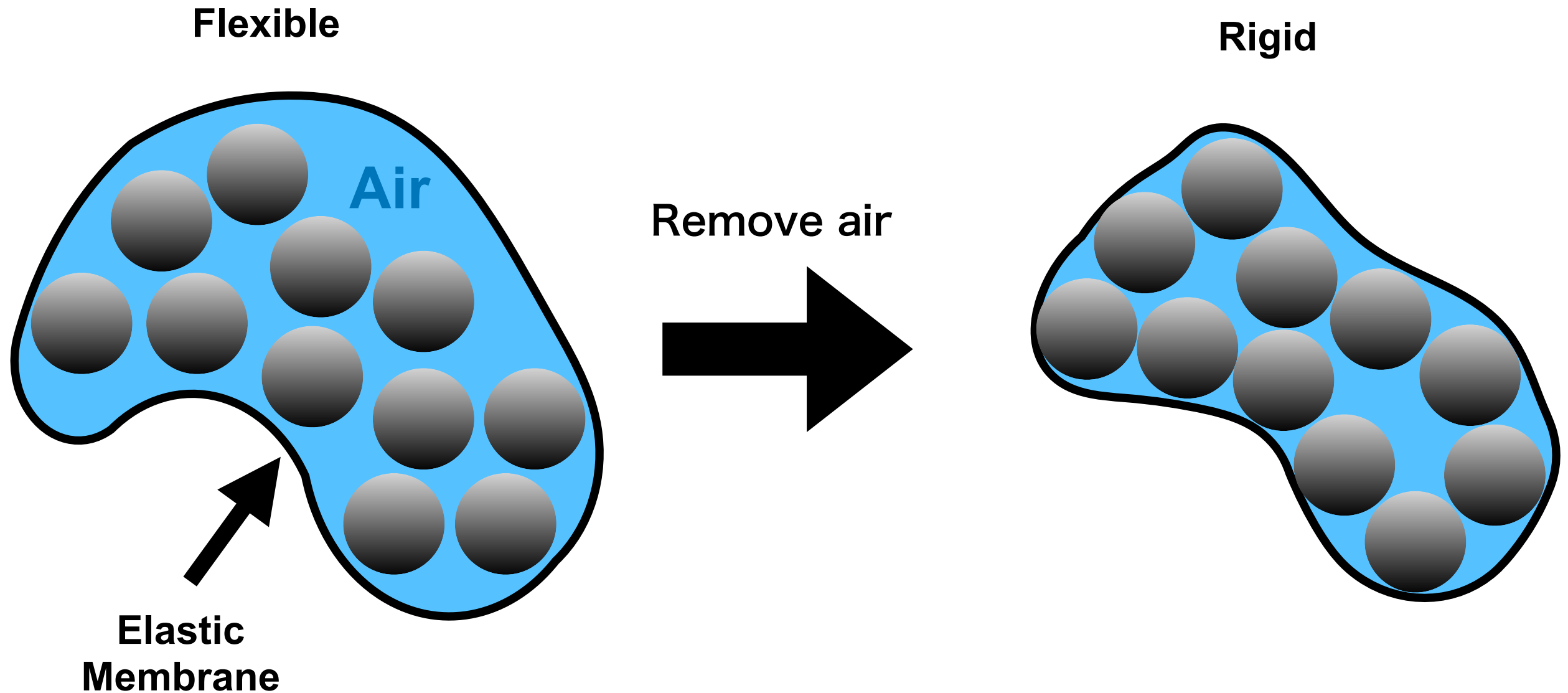


The jamming transition is a phase transition from solid to liquid at zero temperature

Introduction

Related phenomena

“Soft” robot



Introduction

Related phenomena

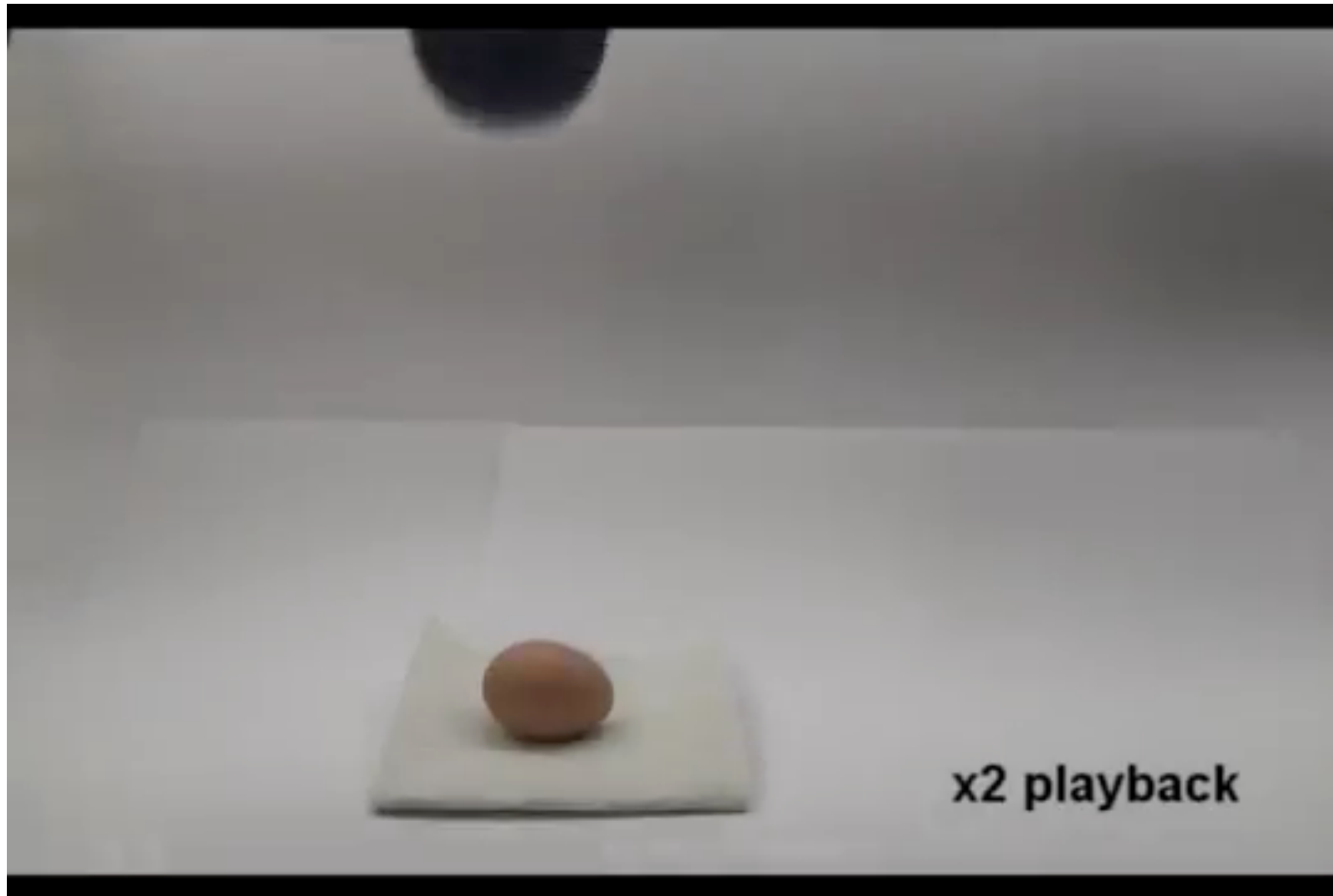
“Soft” robot



Introduction

Related phenomena

“Soft” robot



Cornell Creative Machines Lab: <http://creativemachines.cornell.edu/>

Introduction

Related phenomena

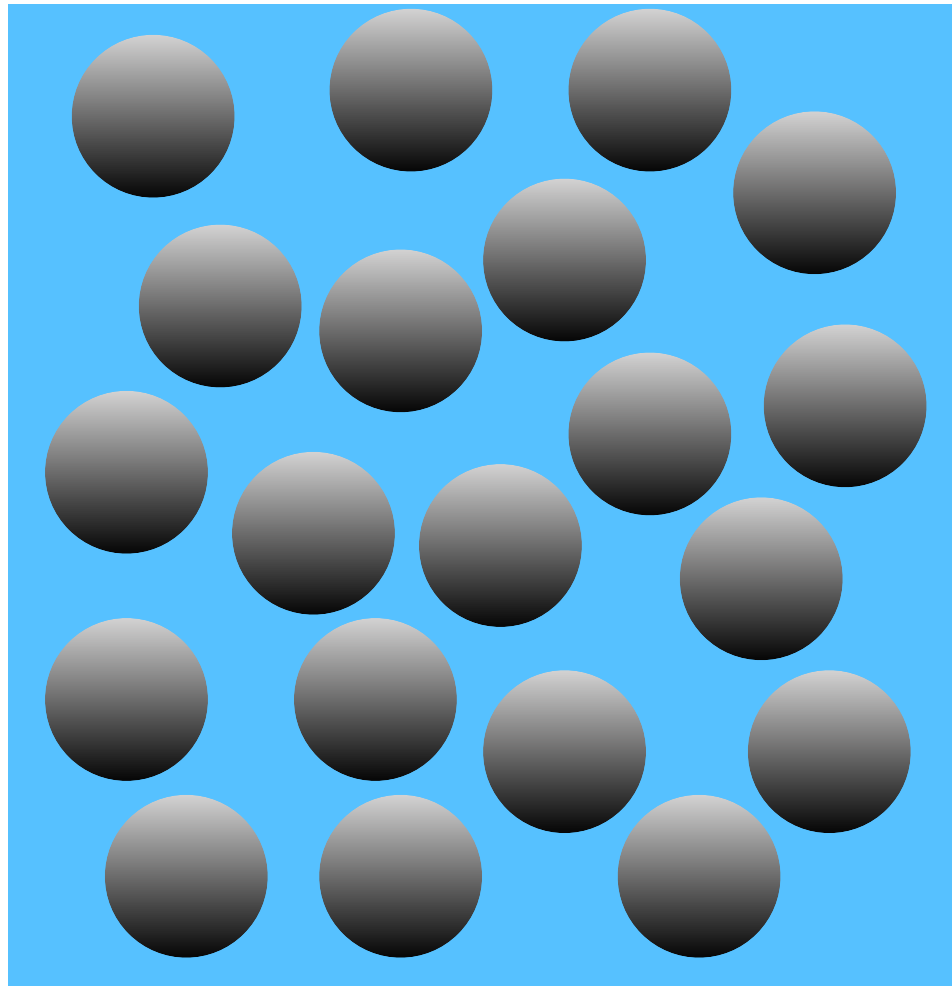
Dilatancy (shear thickening)



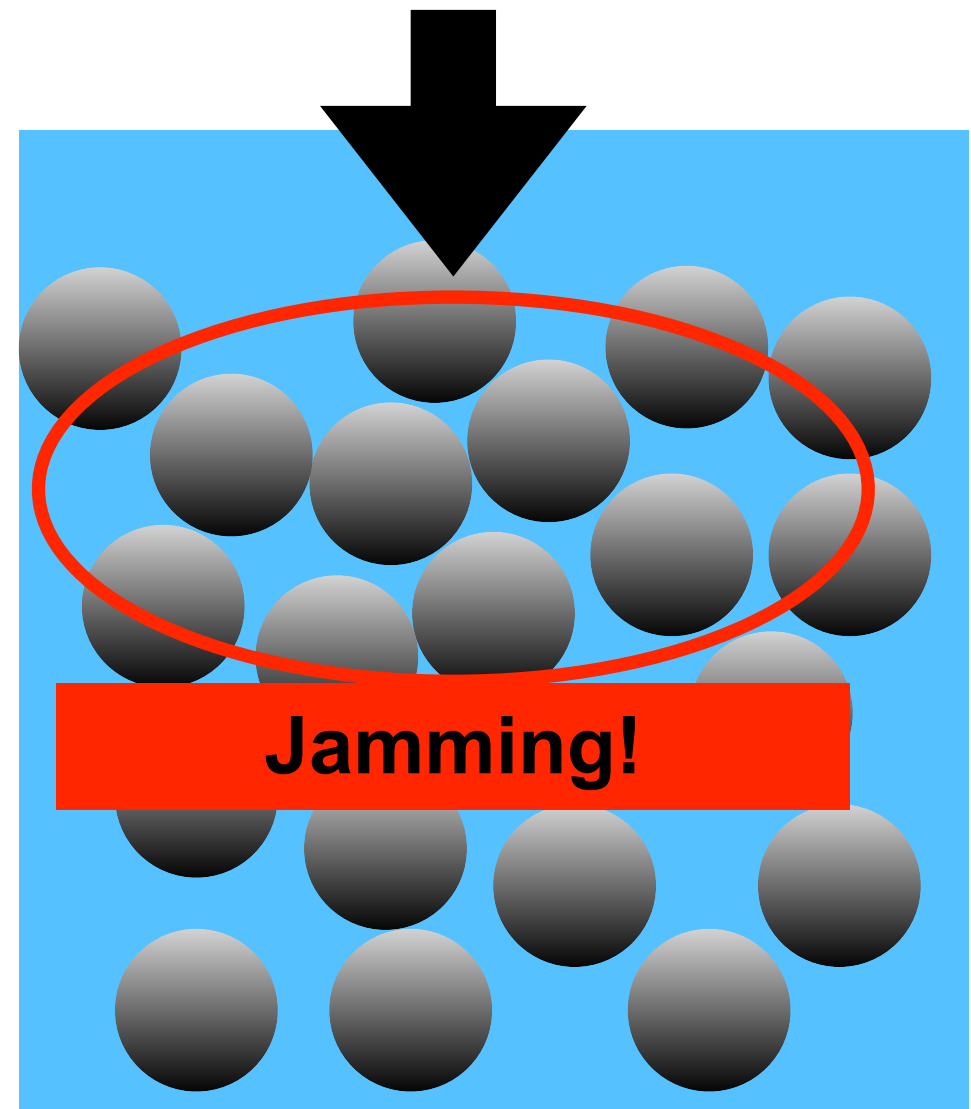
Introduction

Related phenomena

Dilatancy (shear thickening)



Without external force



External force induces
the jamming transition

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Spherical particles

Motivation

Granular particles have complex shapes.



M&M Candies



Sand & rock



Grains



Snow powders

Spherical particles

Motivation

To simplify the problem, we consider spherical particles



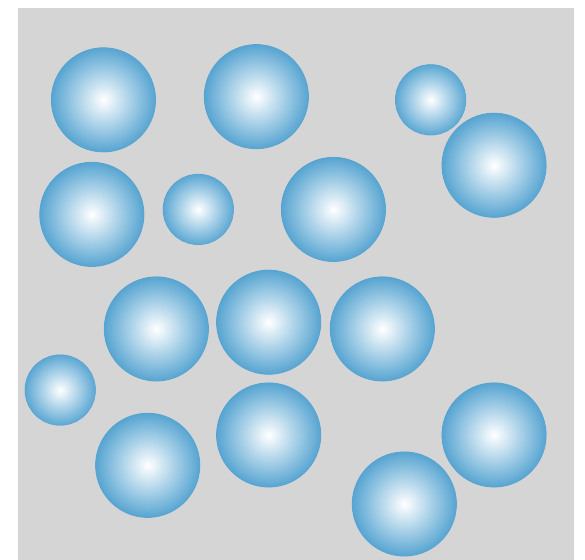
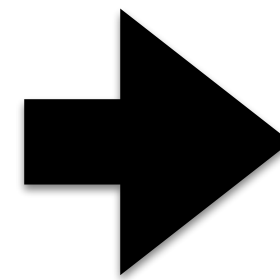
M&M Candies



Sand & rock



Grains



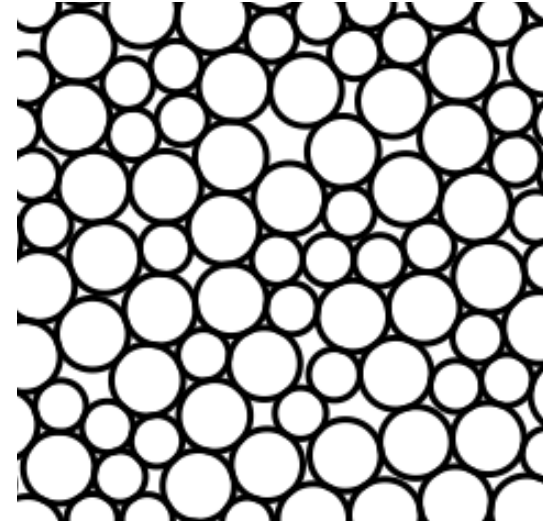
spherical particles

Spherical particles

Motivation



metallic balls



**hard spheres
(simulation)**



Mables

For frictionless spherical particle
The jamming occurs at $\varphi_J \sim 0.64$

Spherical particles

Contact number



J. C. Maxwell
1831-1879

Simple stability argument by Maxwell

of constraints > # of degree of freedom

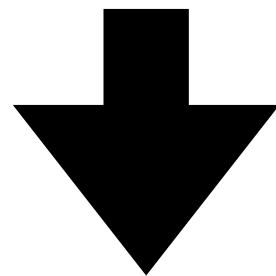
of constraints = # of contacts = $Nz/2$

of degree of freedom = Nd

N: # of particles

z: # of contacts per particle

d: spatial dimensions



$$z > 2d$$

Transition may occur at $z_J = 2d$

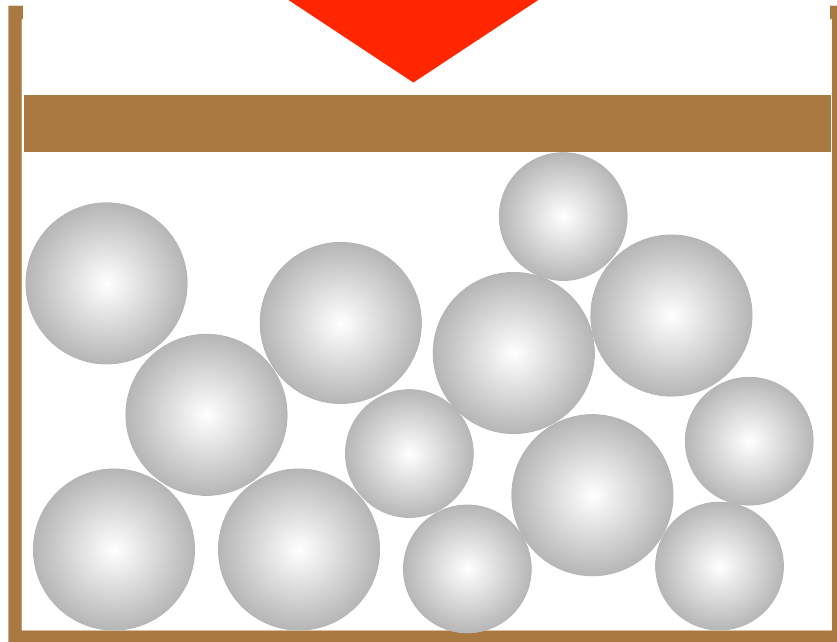
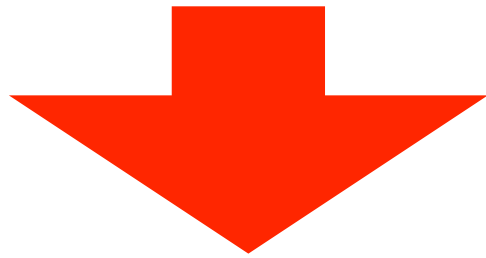
Spherical particles

Contact number

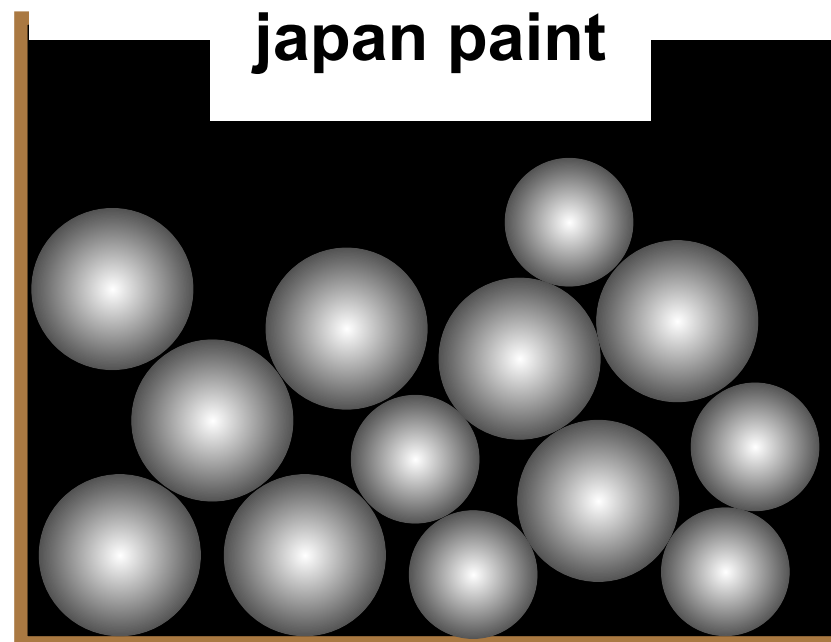
Experiment

J. Bernal and J. Mason (1960)

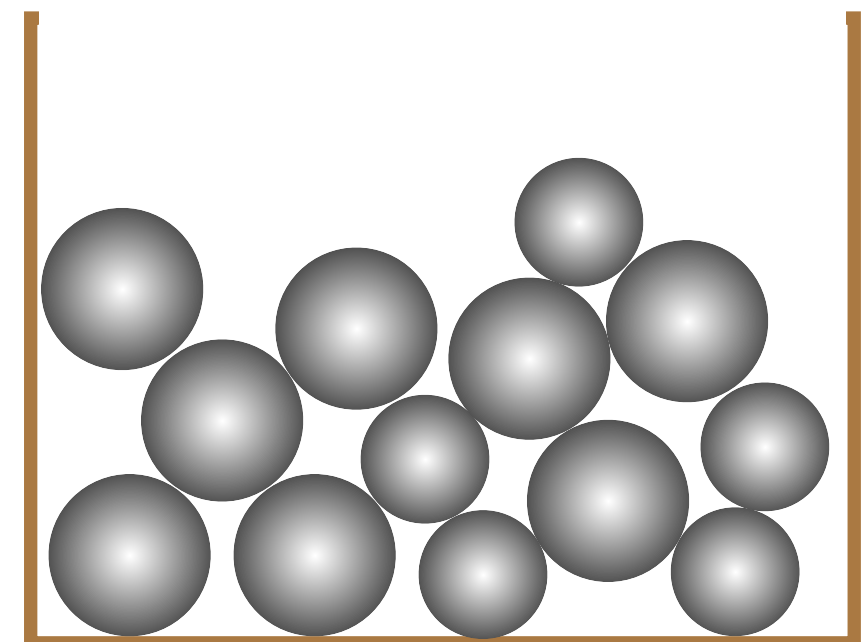
1. Compress



2. Put Japan paint
(Urushi)



3. Remove Japan paint

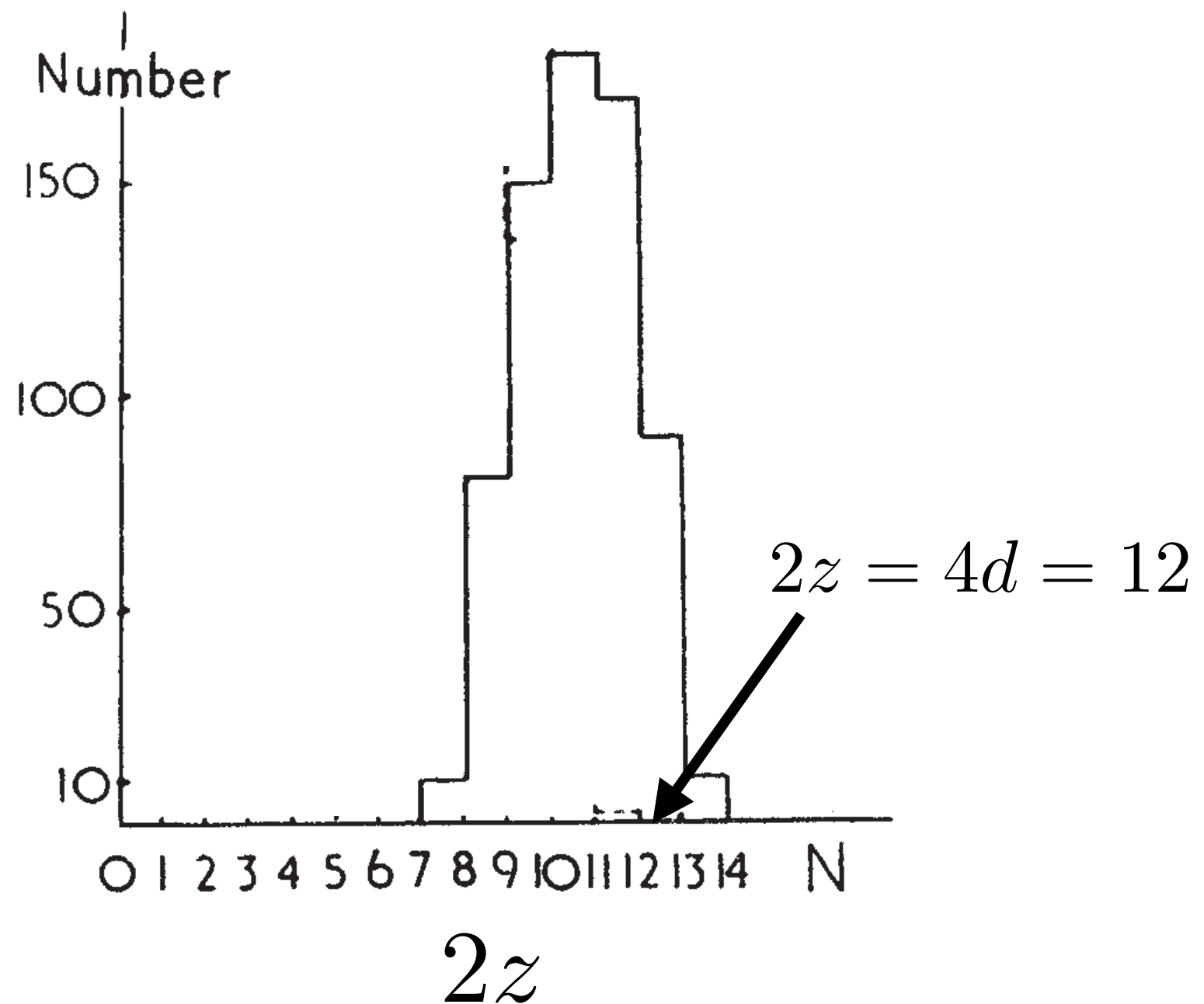
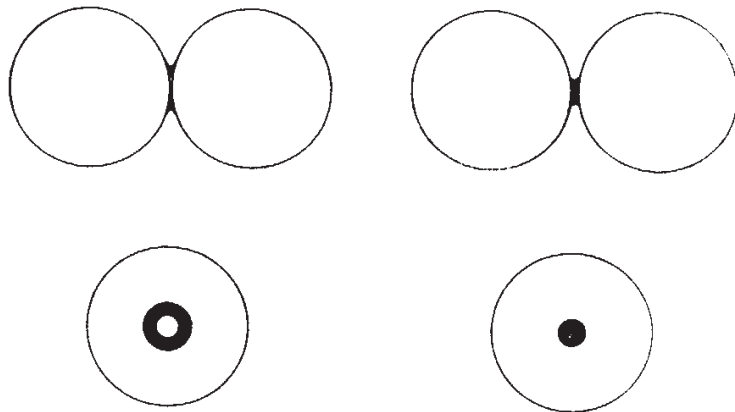


Spherical particles

Contact number

J. Bernal and J. Mason (1960)

N= 500
metallic balls painted
with black japan paint



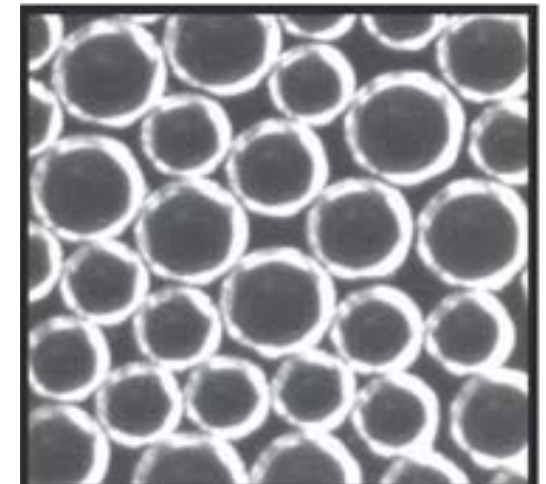
Spherical particles

Contact number

Simulation

Harmonic potential

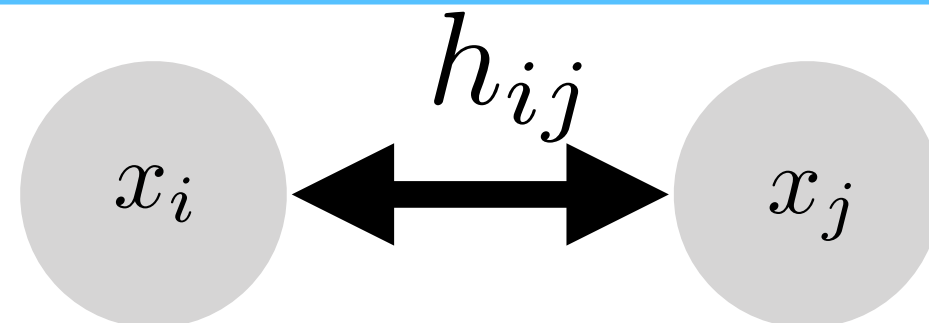
$$E = \sum_{i < j} v(h_{ij}) = \frac{\varepsilon}{2} \sum_{i < j} h_{ij}^2 \theta(-h_{ij})$$



Wet foams

Gap function

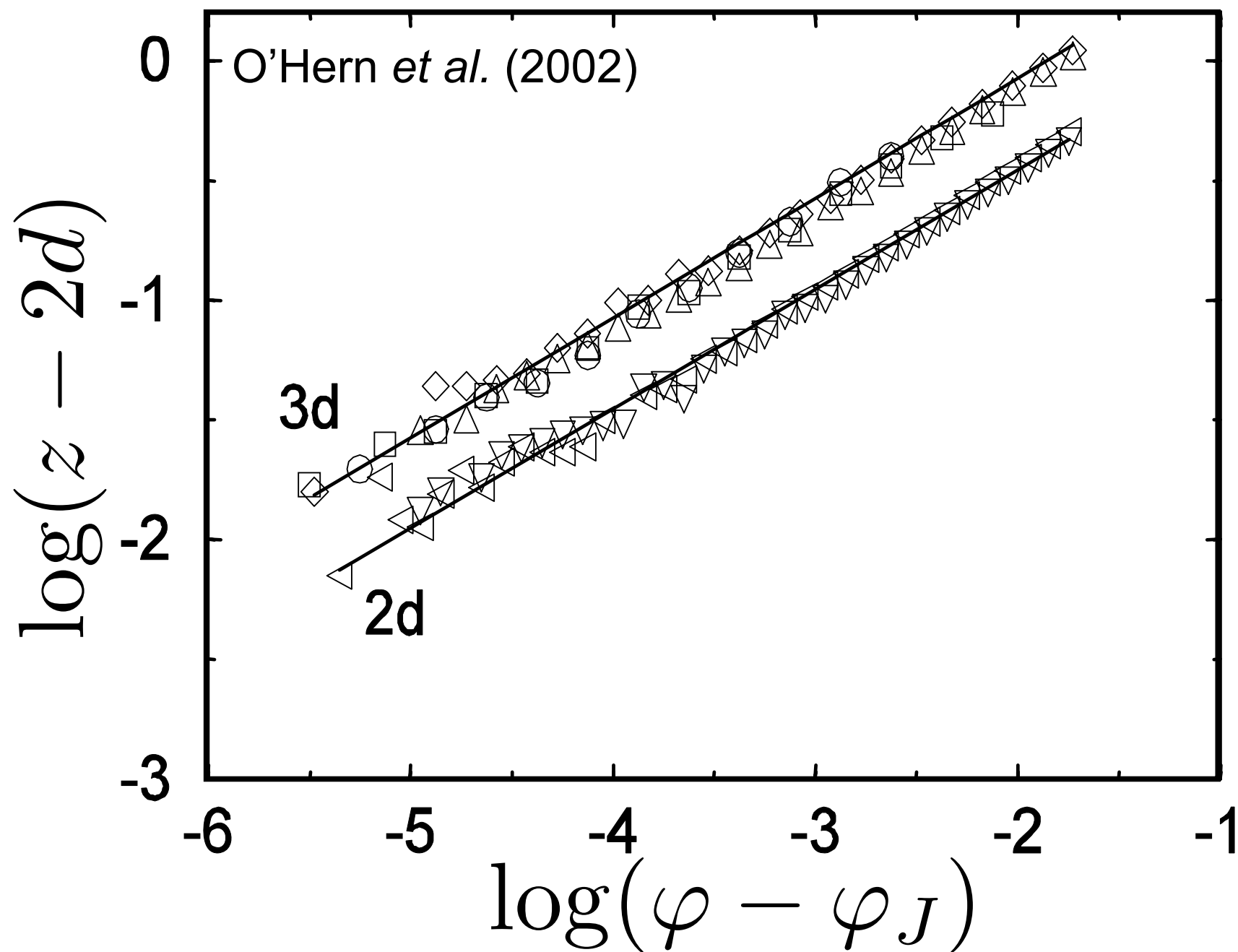
$$h_{ij} = |x_i - x_j| - \sigma_{ij}$$



Spherical particles

Power-law

Contact number of hard spheres



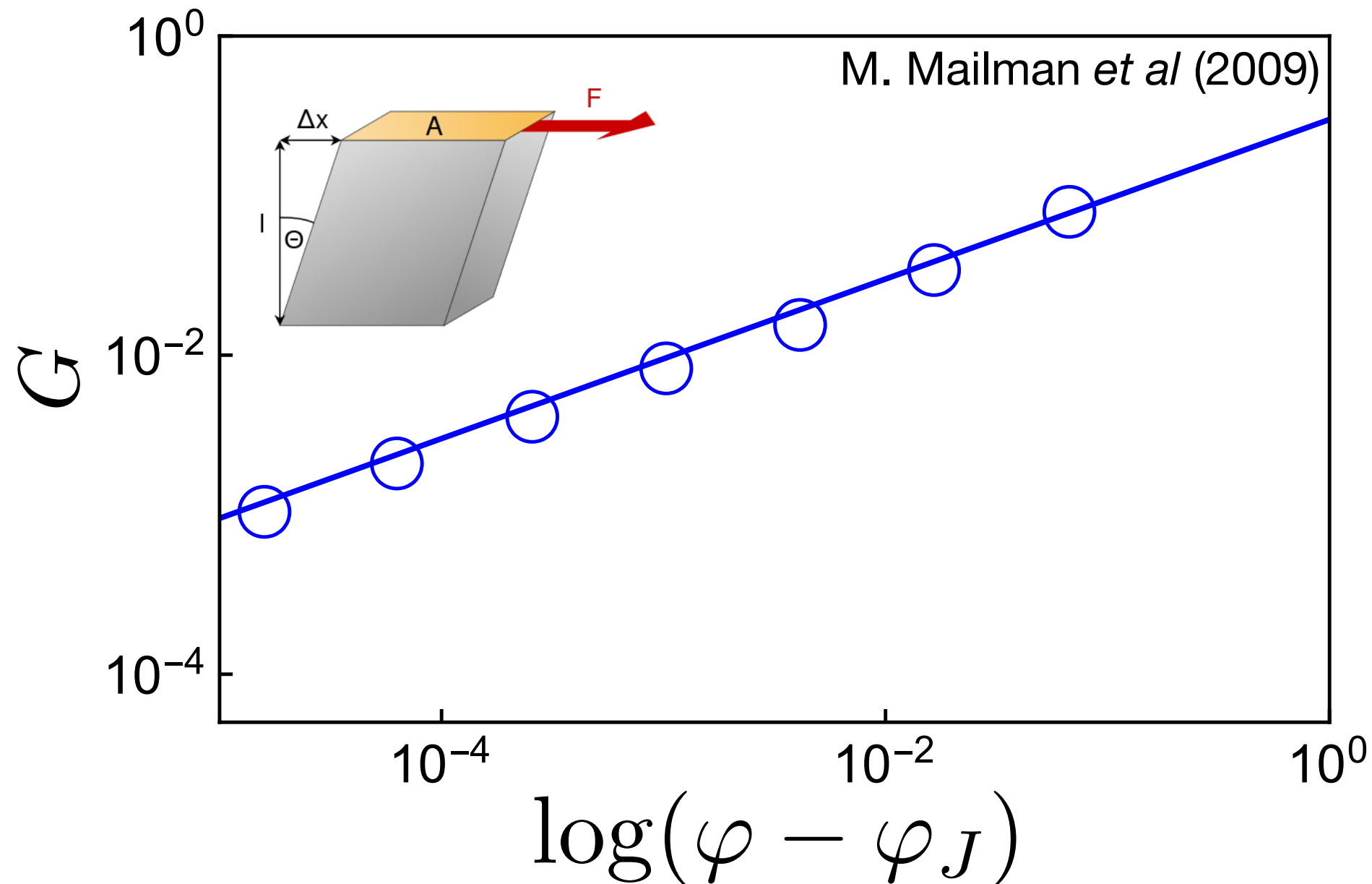
Scaling

$$z - 2d \sim (\varphi - \varphi_J)^{1/2}$$

Spherical particles

Power-law

Shear modulus of harmonic spheres



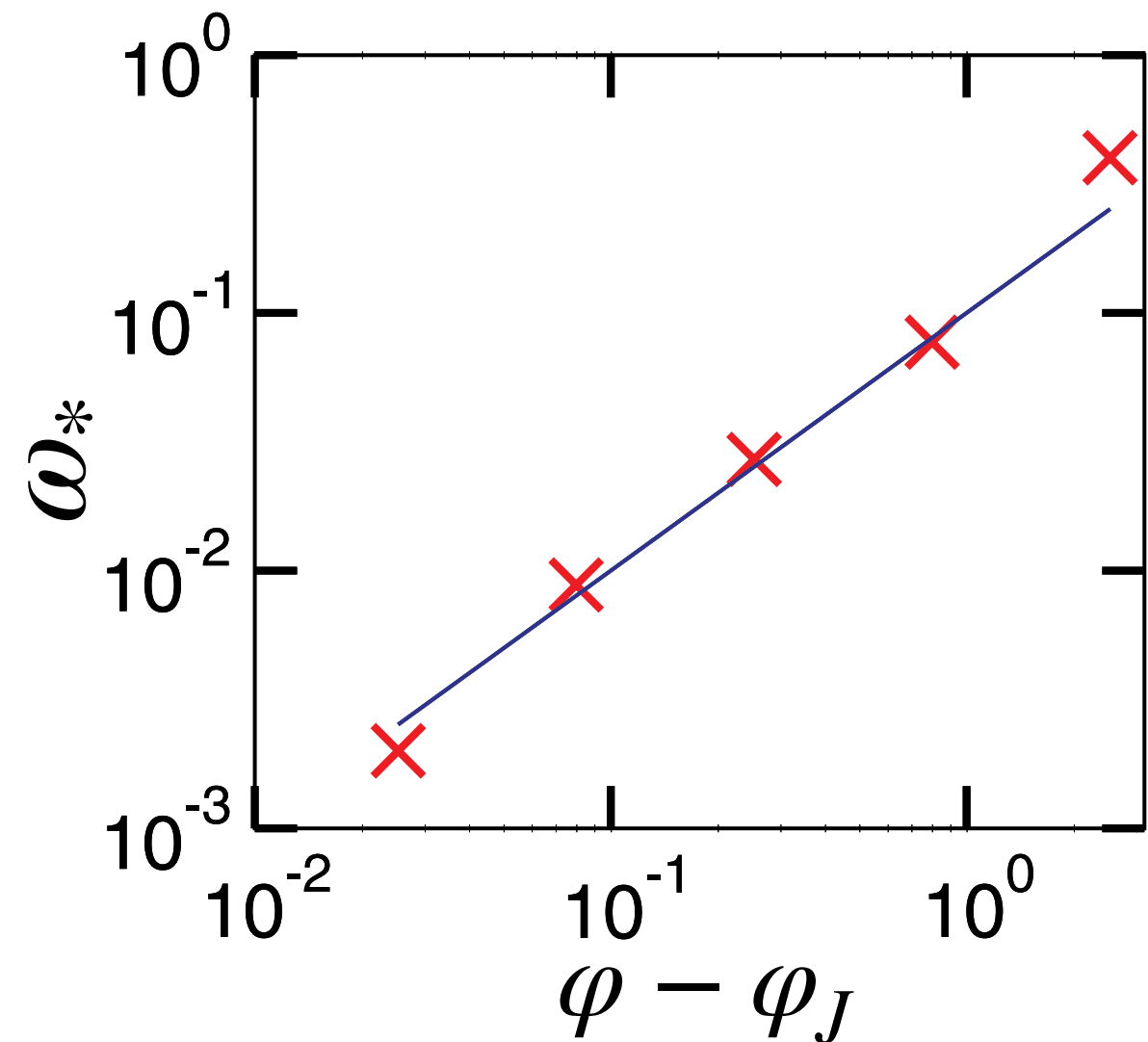
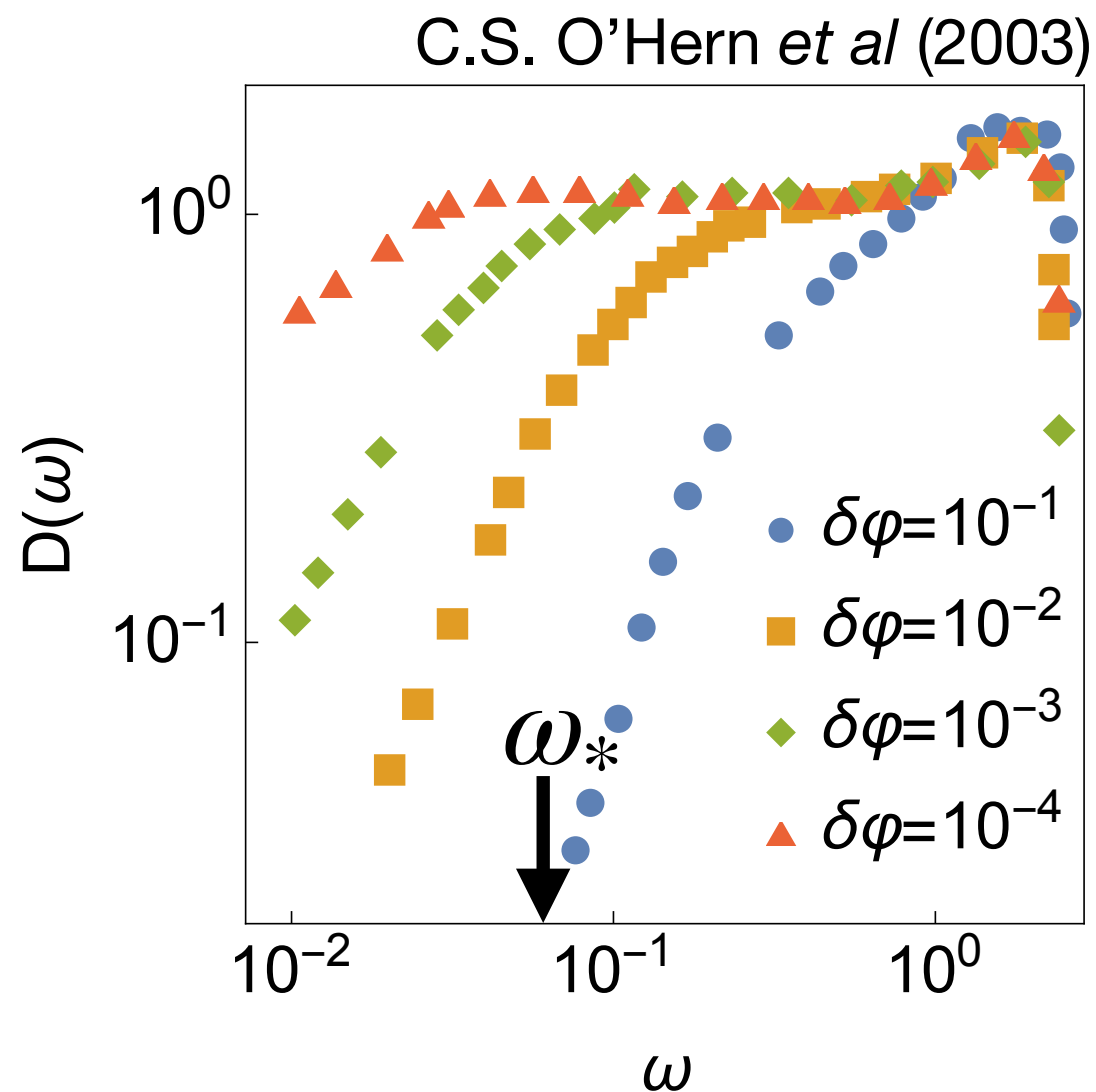
Scaling

$$G \sim (\varphi - \varphi_J)^{1/2}$$

Spherical particles

Power-law

Scaling of the vibrational density of states



Scaling

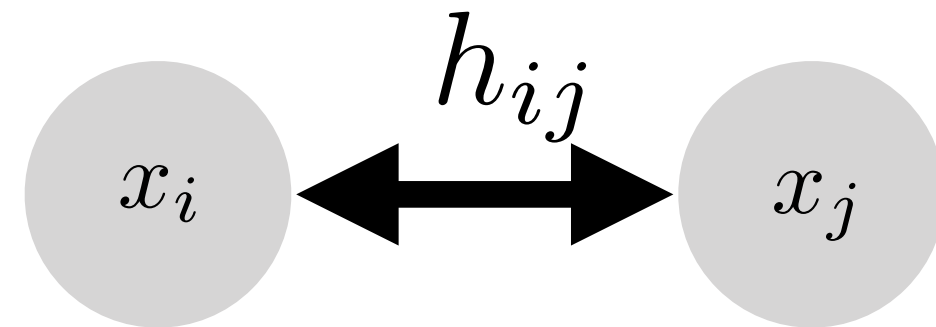
$$\omega_* \sim (\phi - \phi_J)^{1/2}$$

Spherical particles

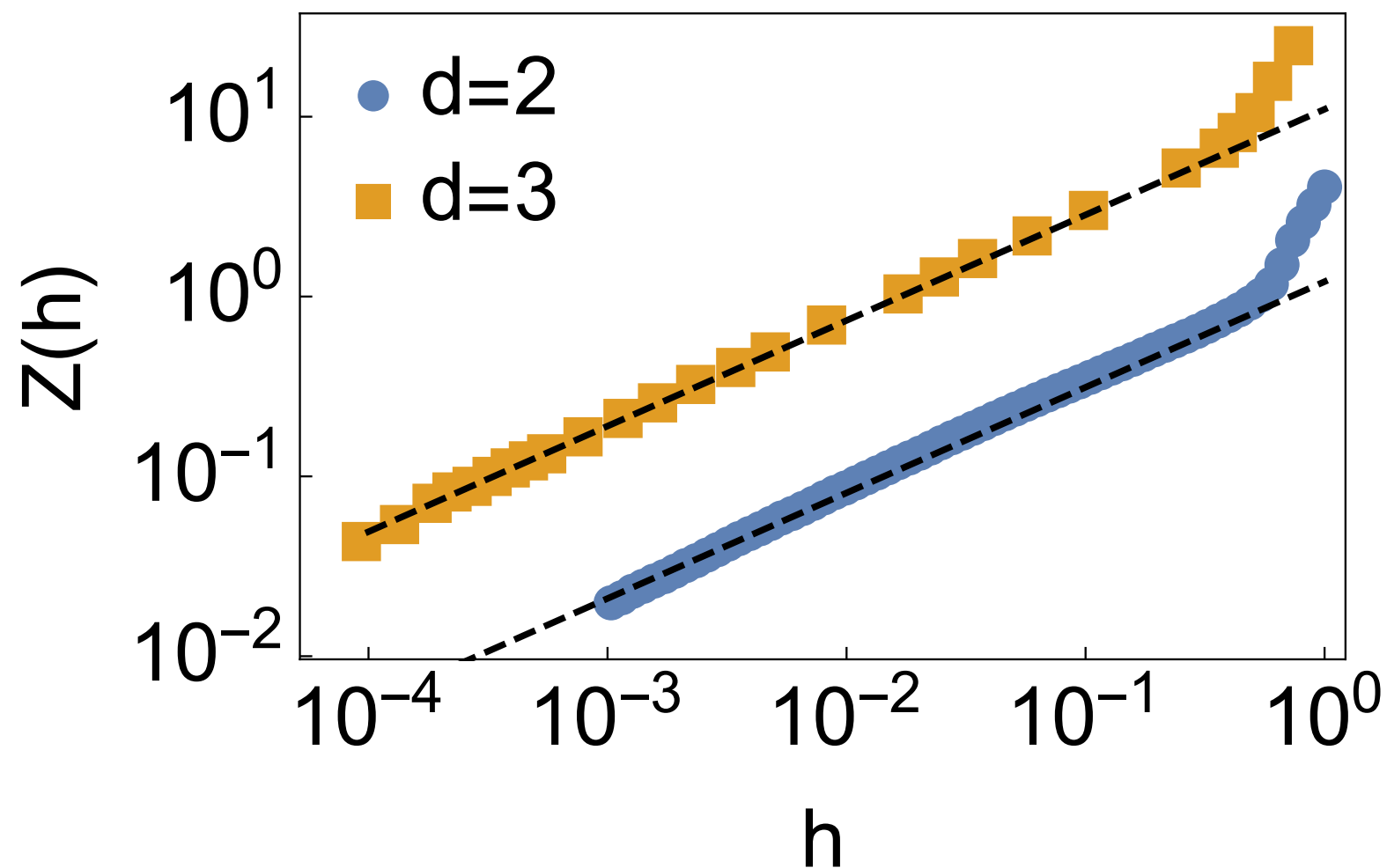
Power-law

Gap distribution function

$$g(h) = \frac{1}{N_c} \sum_{i < j} \delta(h_{ij} - h), \quad Z(h) = \int_0^h dh' g(h')$$



P. Charbonneau *et al.* (2014)



Scaling

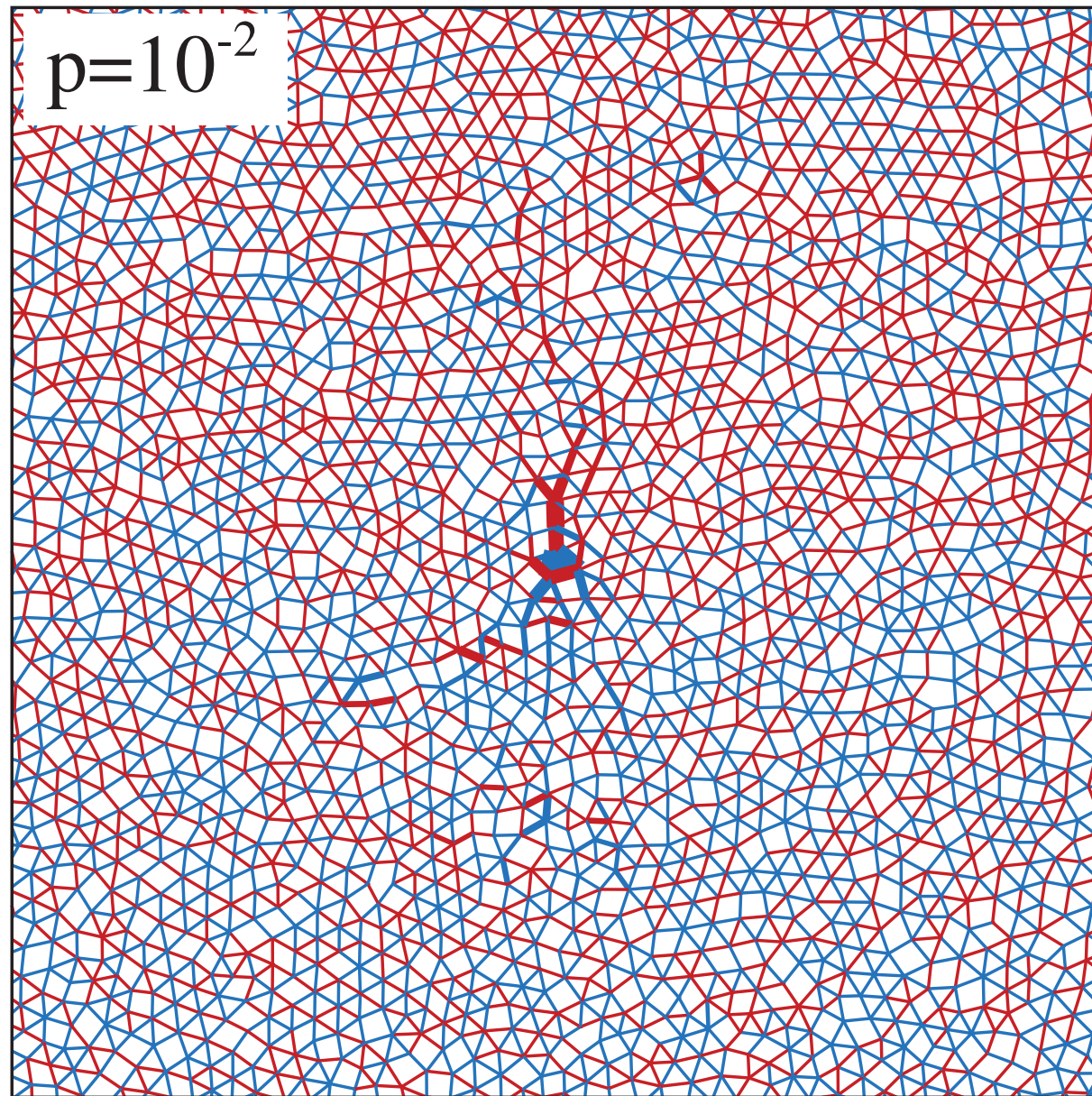
$$g(h) \sim h^{-\gamma}, \quad \gamma = 0.41$$

Spherical particles

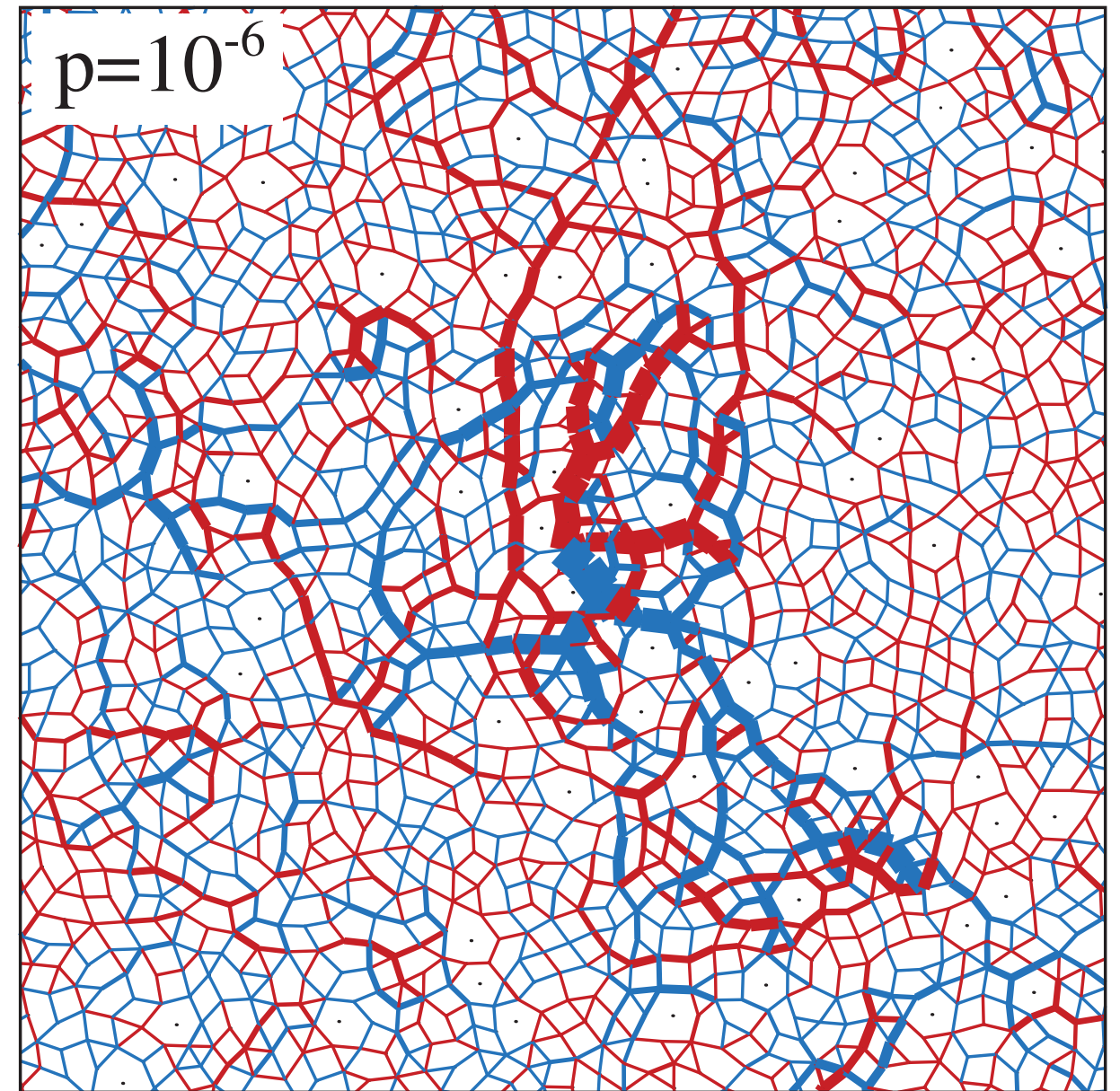
Power-law

W. G. Ellenbroek *et al.* (2006)

Far from the jamming transition point



Near the jamming transition point



Correlation length

$$\xi \sim (\varphi - \varphi_J)^{-1/4}$$

E. Lerner *et al.* (2015), M. Shimada *et al.* (2019)

Spherical particles

Summary

Several physical quantities exhibit the power-law behavior

Contact number

$$z - 2d \sim \delta\varphi^{1/2}$$

Shear modulus

$$G \sim \delta\varphi^{1/2}$$

Correlation length

$$\xi \sim \delta\varphi^{-1/4}$$

Pair correlation

$$g(h) \sim \begin{cases} \delta\varphi^{-\mu\gamma} (h \ll \delta\varphi^\mu) \\ h^{-\gamma} & (h \gg \delta\varphi^\mu) \end{cases}$$

$\gamma = 0.41, \mu = \frac{1}{1-\gamma}$

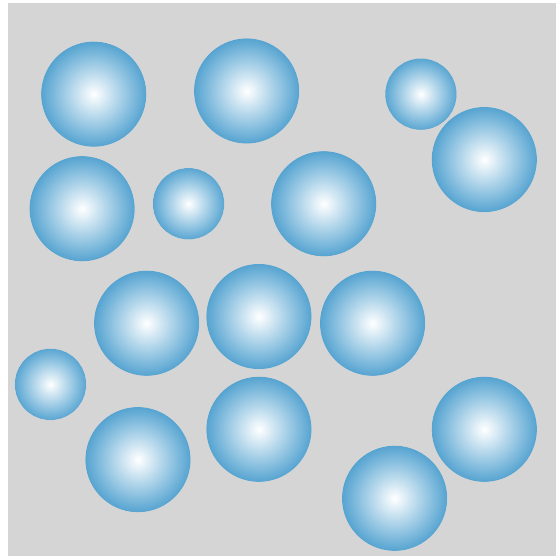
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- **Jamming of non-spherical particles**

Outline

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- **Jamming of non-spherical particles**

Non-spherical particles

Motivation



Spherical particles

Non-spherical particles

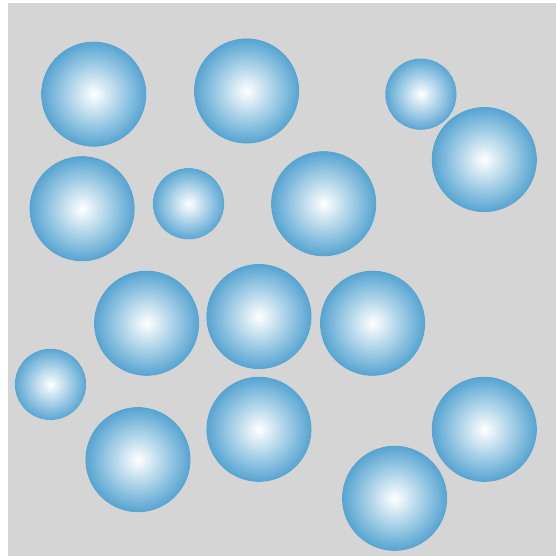
Motivation



Sand & rock



Snow powders



Spherical particles



Grains



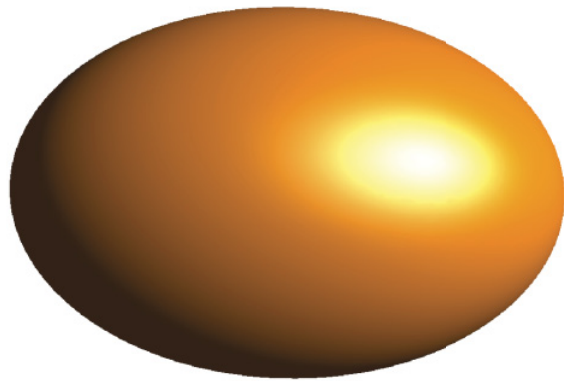
M&M Candies

However, in most cases, the constituent particles are **non-spherical**.

Non-spherical particles

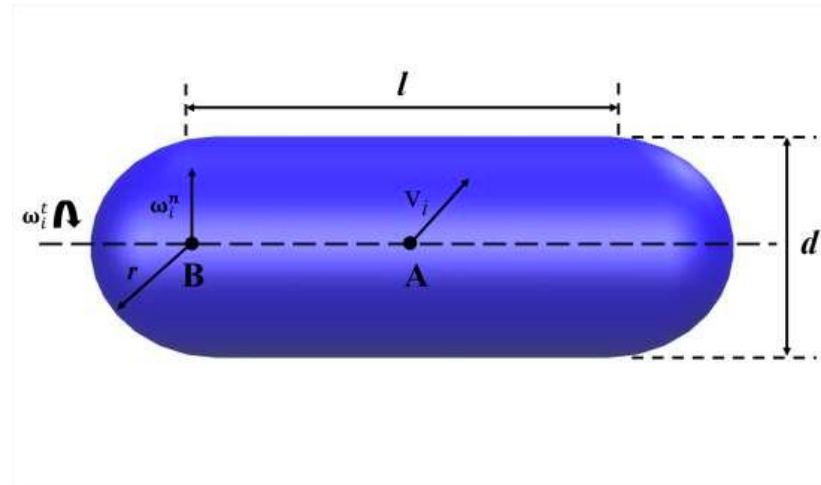
Previous numerical studies

Ellipsoid



A. Donev *et al* (2007)
M. Mailman *et al* (2009)
C. F. Schreck *et al.* (2010, 2012)

Sphero-cylinder



T. Marschall and S. Teitel (2018)

Other shapes

???

Purpose of this study

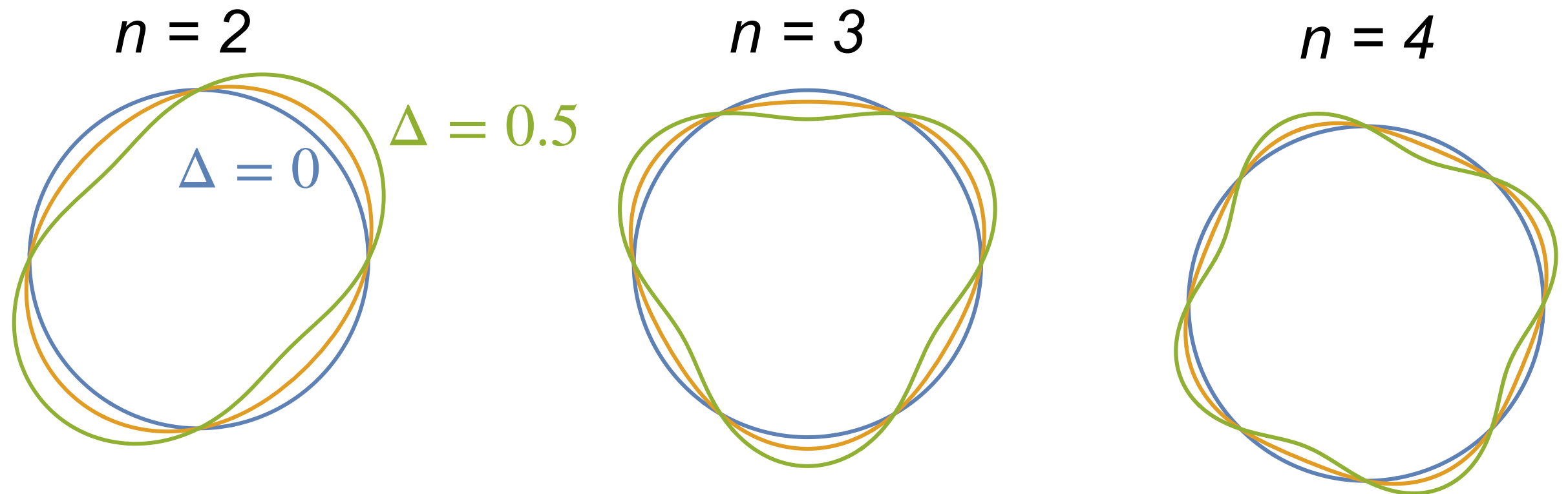
Systematic investigation of various shapers of particles

Non-spherical particles

Model

Radius of a particle

$$R_i(\theta) = R_i^0 \left(1 + F_n(\theta) \right), \quad F_n(\theta) = \frac{\Delta}{n} \sin(n\theta)$$



We investigate the jamming transitions of various shapes by changing n and Δ systematically.

Non-spherical particles

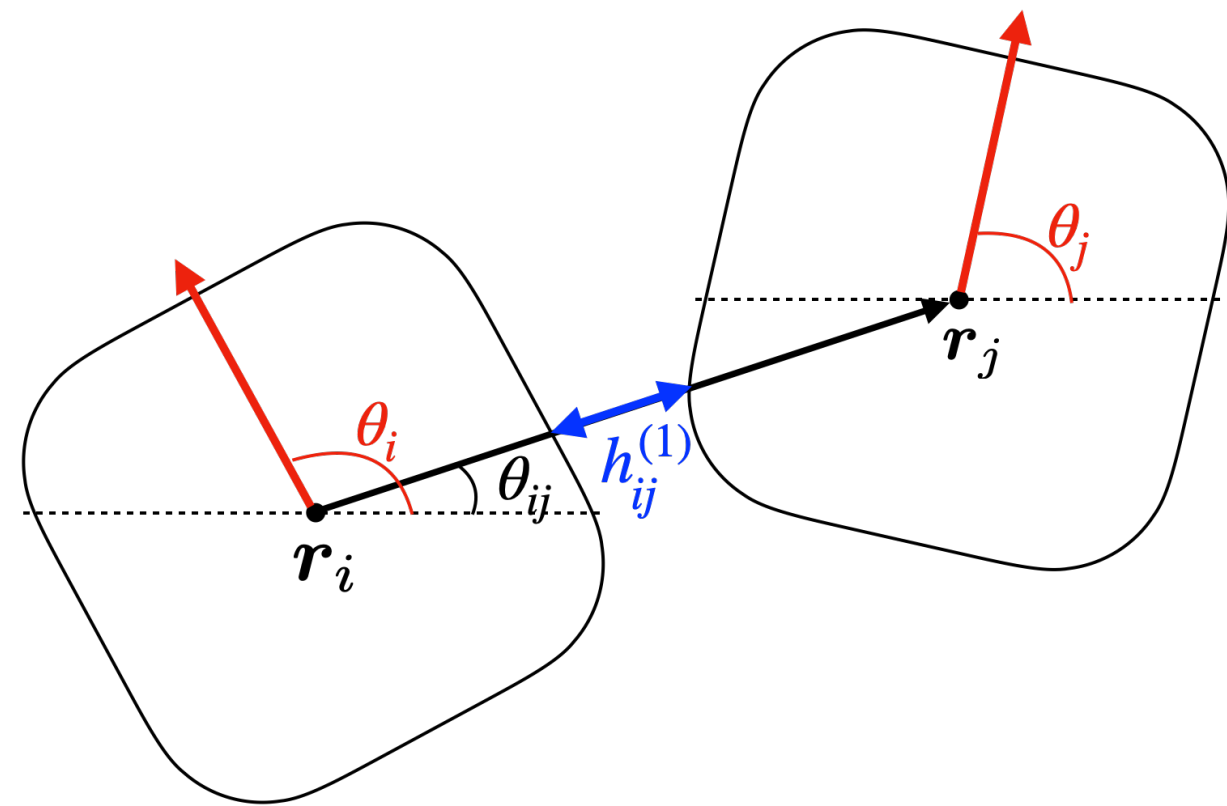
Model

Harmonic potential

$$V_N = \sum_{i < j} \frac{h_{ij}^2}{2} \theta(-h_{ij})$$

Gap function: the minimal distance between two particles

For nearly spherical particles, we can derive an analytic form by using a perturbation w.r.t Δ .



Small Δ expansion

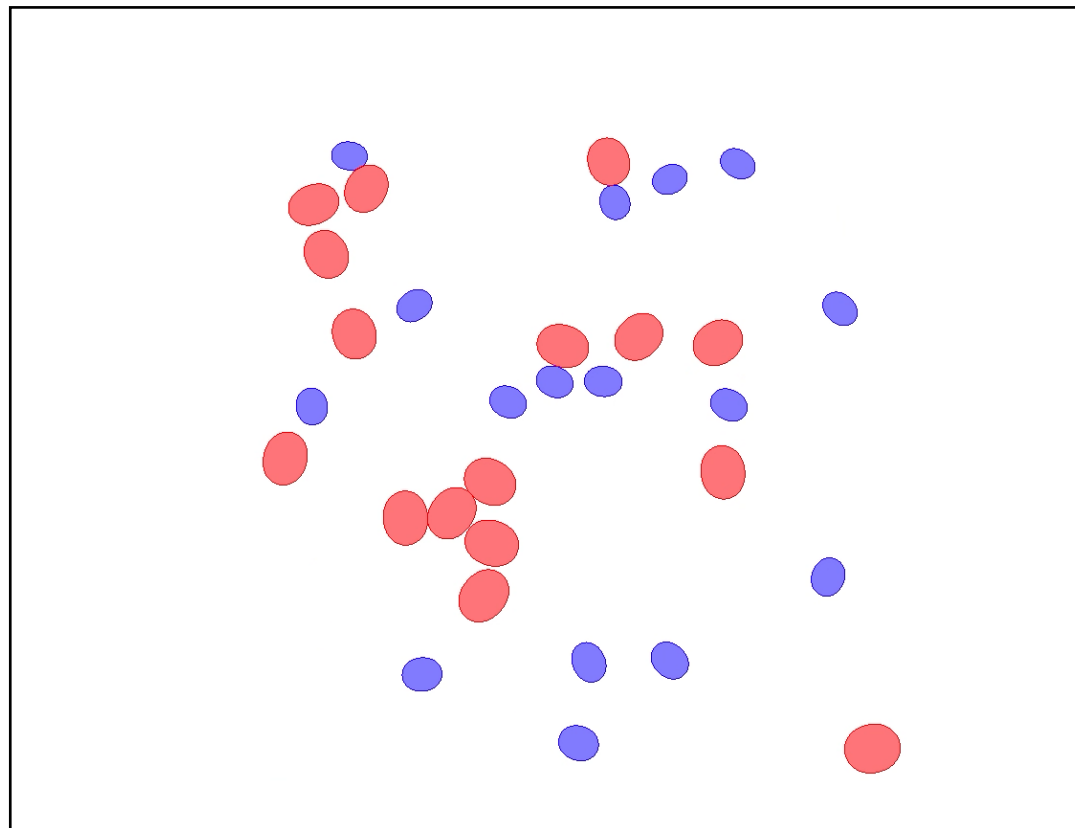
$$\begin{aligned} h_{ij} &= h_{ij}^{(1)} + O(\Delta^2), \\ h_{ij}^{(1)} &= |\mathbf{r}_i - \mathbf{r}_j| - R_i^0 - R_j^0 \\ &\quad - \Delta [R_i^0 F_n(\theta_i - \theta_{ij}) + R_j^0 F_n(\theta_j - \theta_{ji})] \end{aligned}$$

Non-spherical particles

Numerics

Algorithm

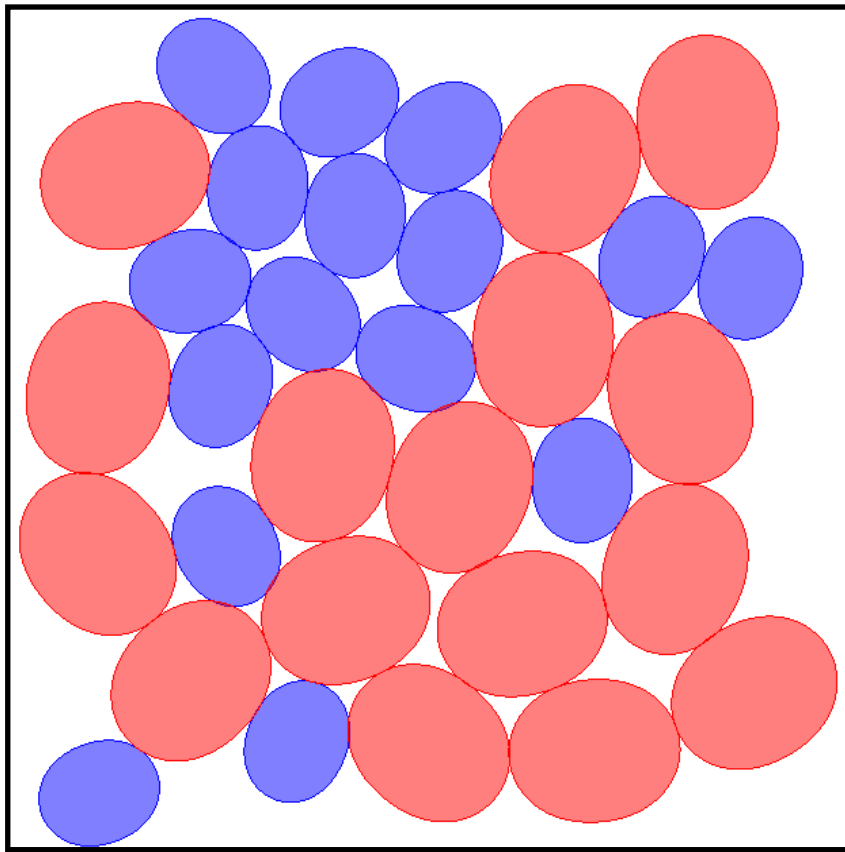
1. Start from a random configuration.
2. Increase the packing fraction $\phi \rightarrow \phi + \delta\phi$.
3. Remove contact by the energy minimization.
4. Repeat 2 and 3 until the jamming transition point.



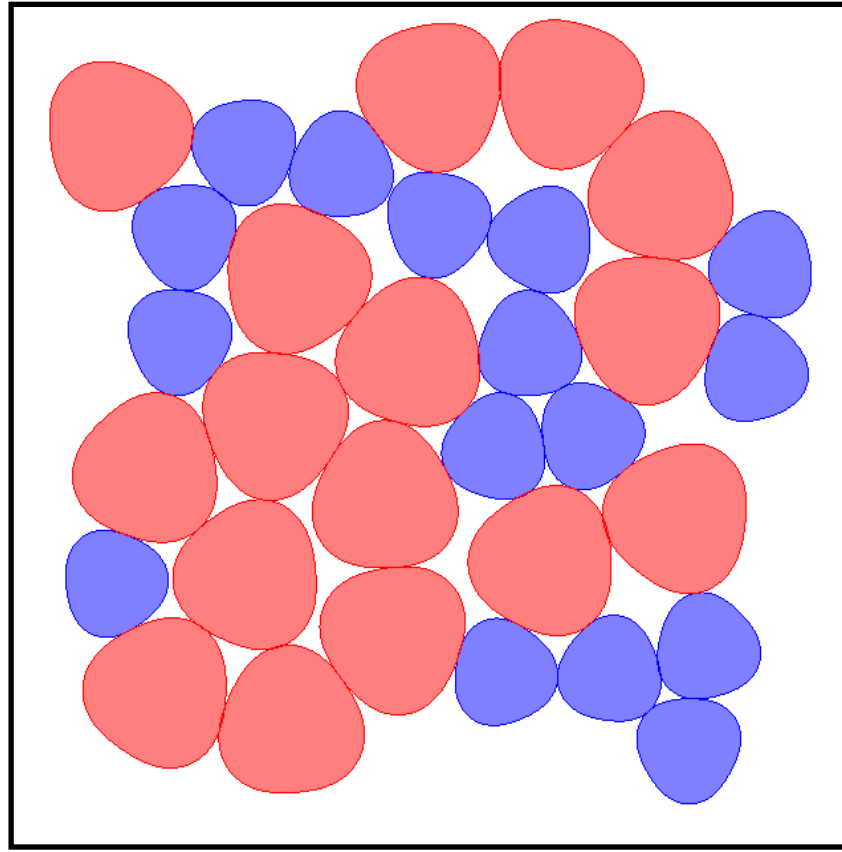
Non-spherical particles

Numerics

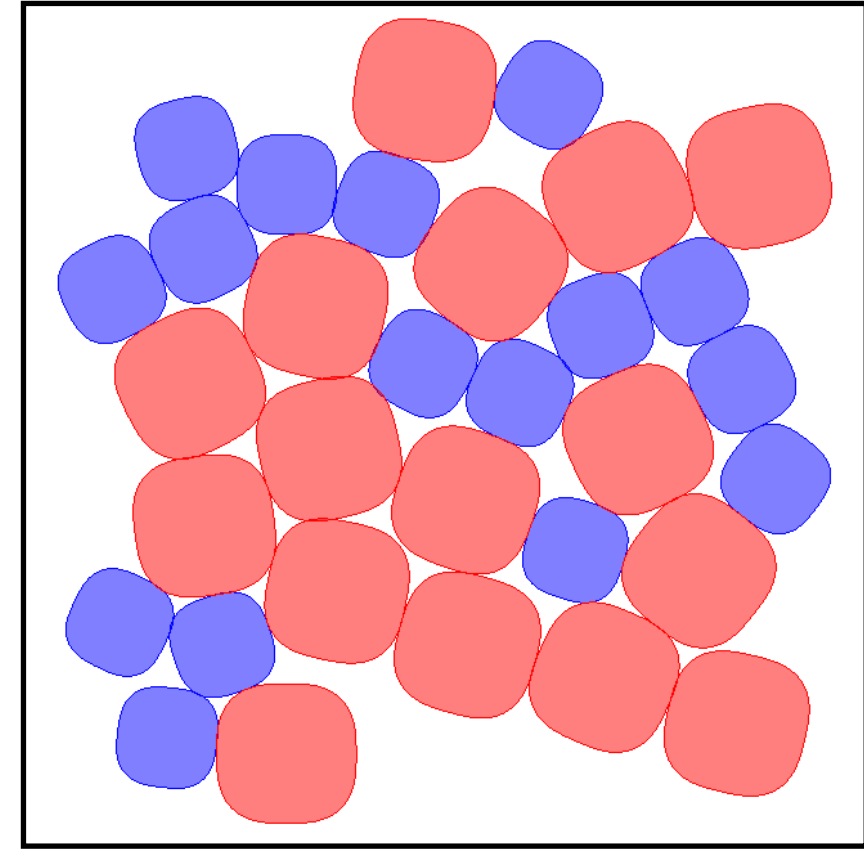
$n=2$



$n=3$



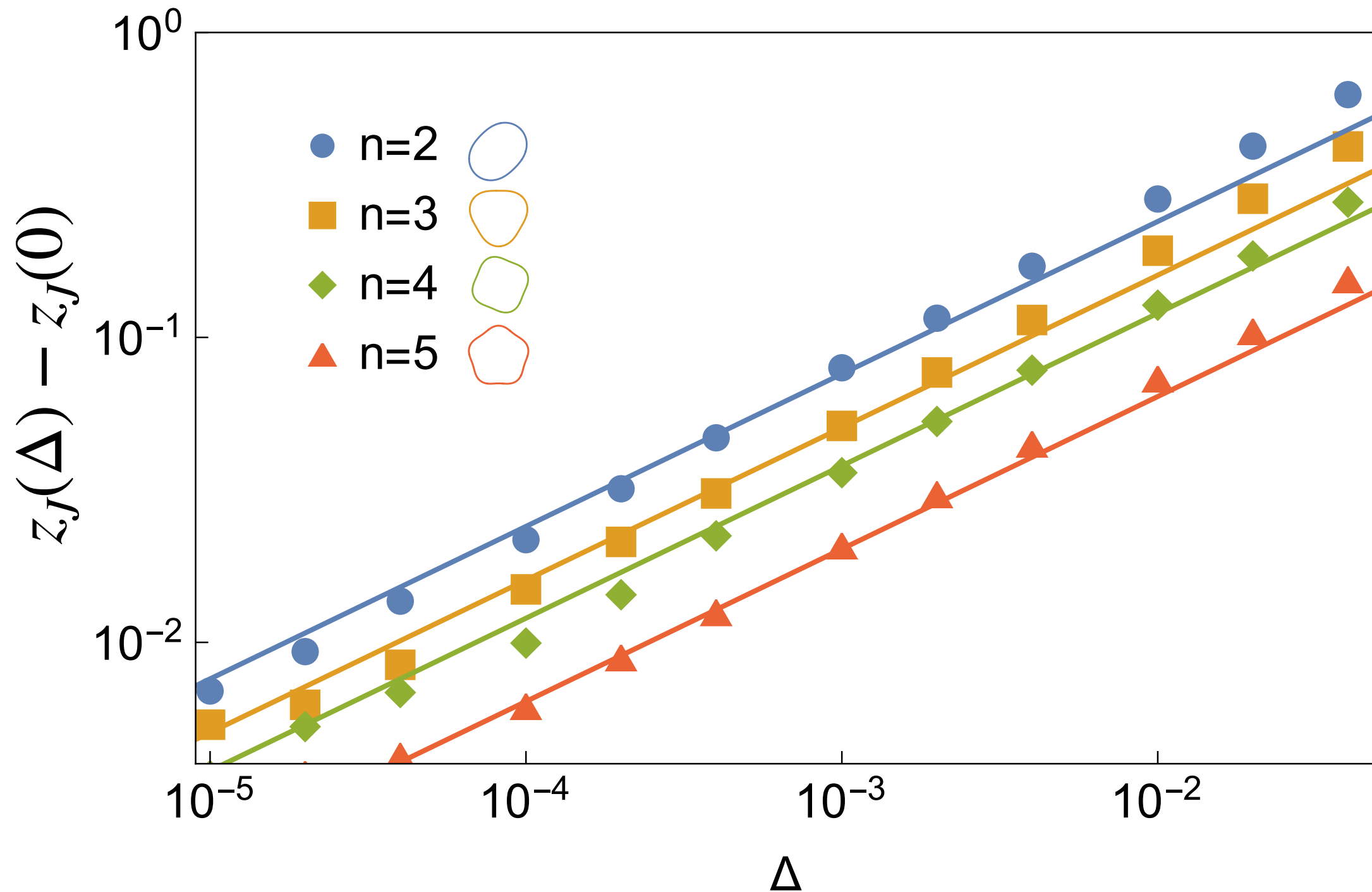
$n=4$



We focus on the configurations at the jamming transition point.

Non-spherical particles

Contact number



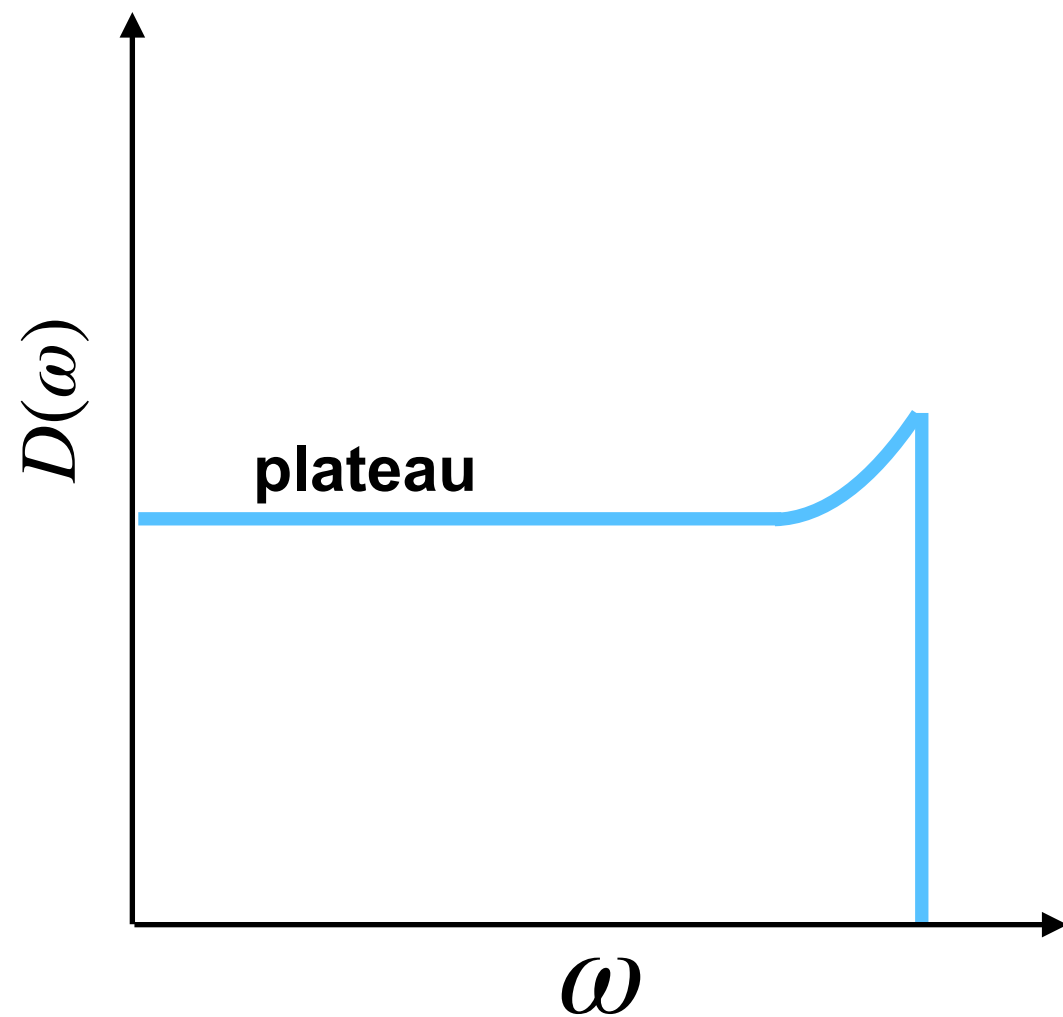
The contact number increases with the aspect ratio as

$$z_J(\Delta) - z_J(0) \sim \Delta^{1/2}$$

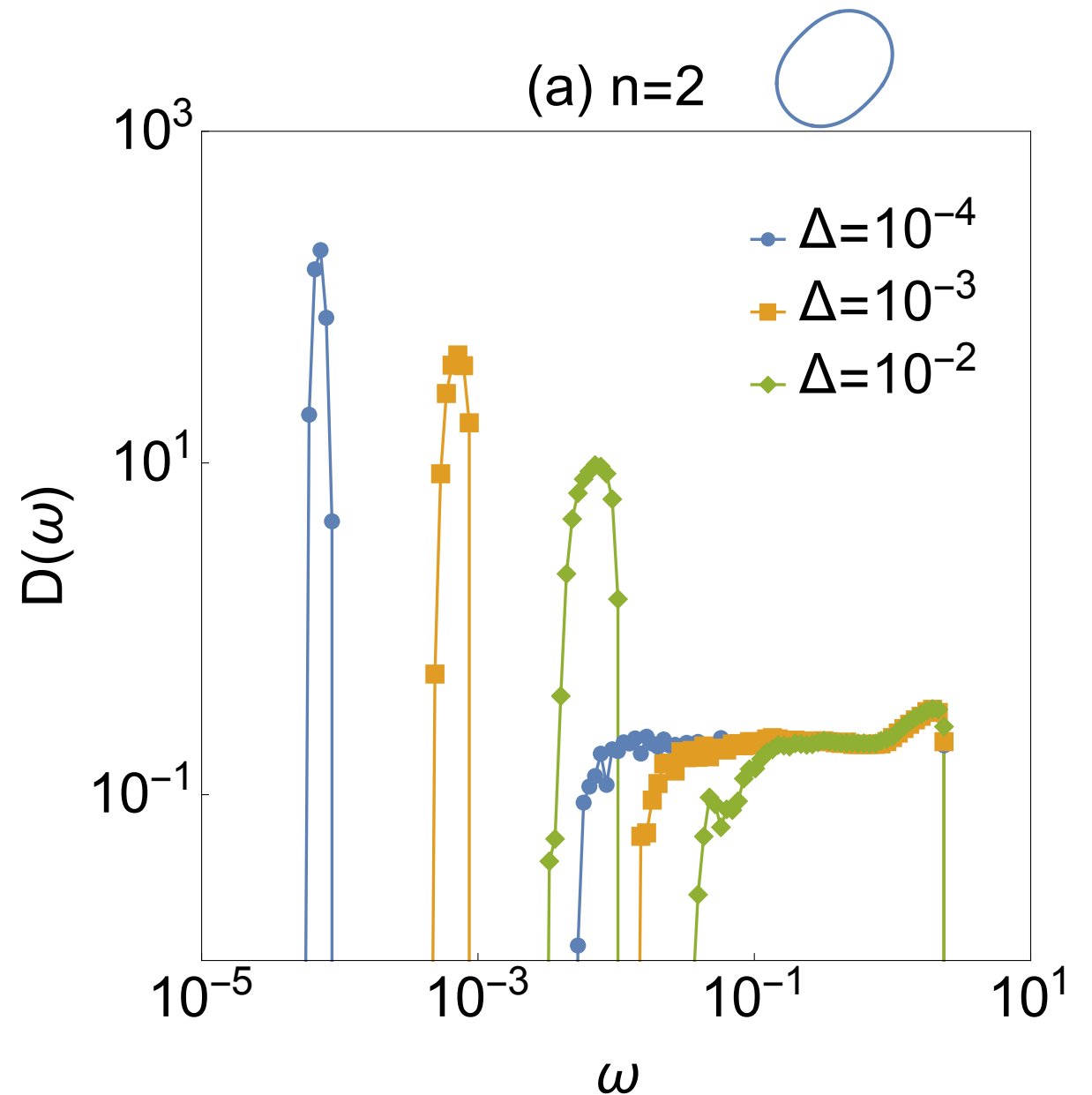
Non-spherical particles

Vibrational density of states

**Spherical
(at jamming)**



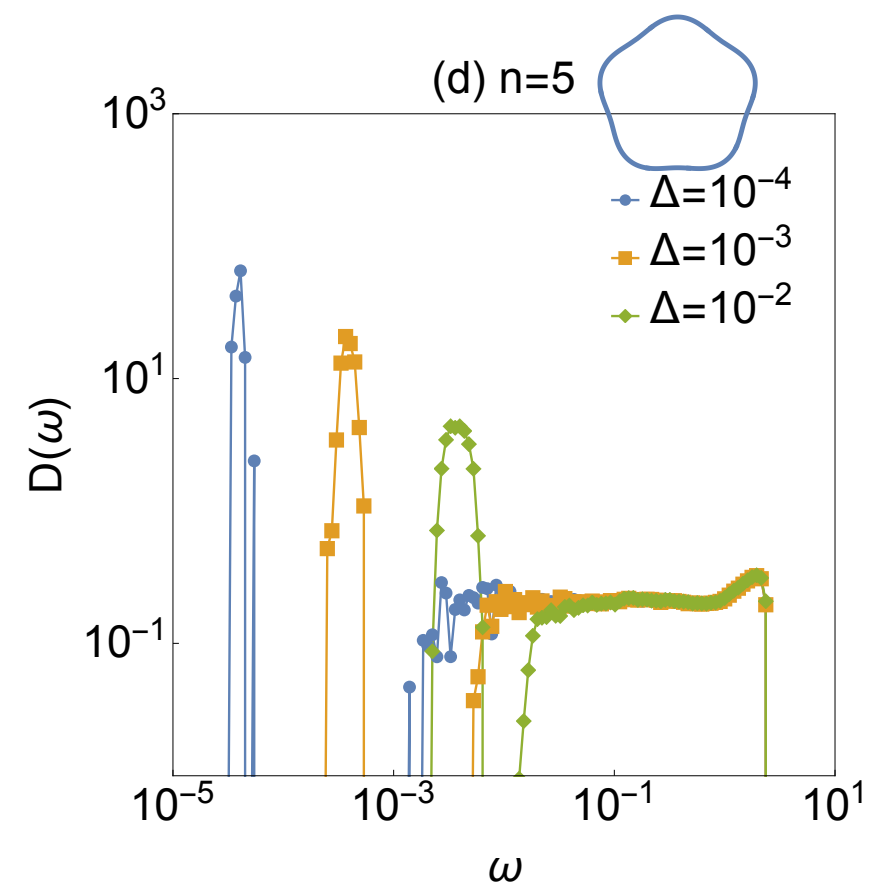
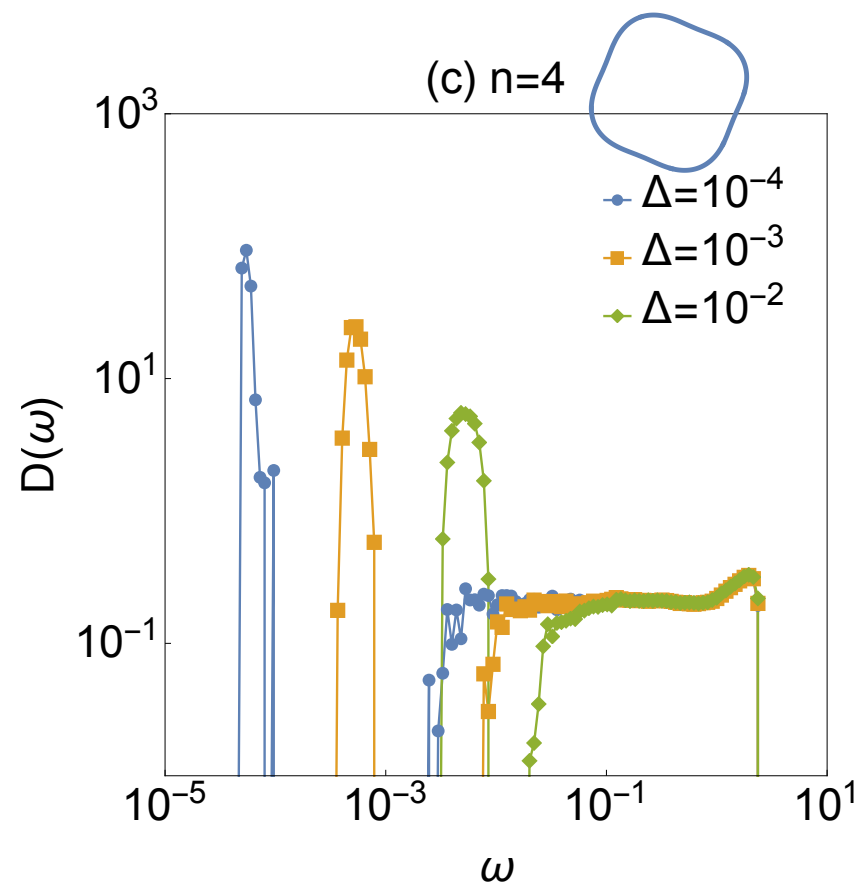
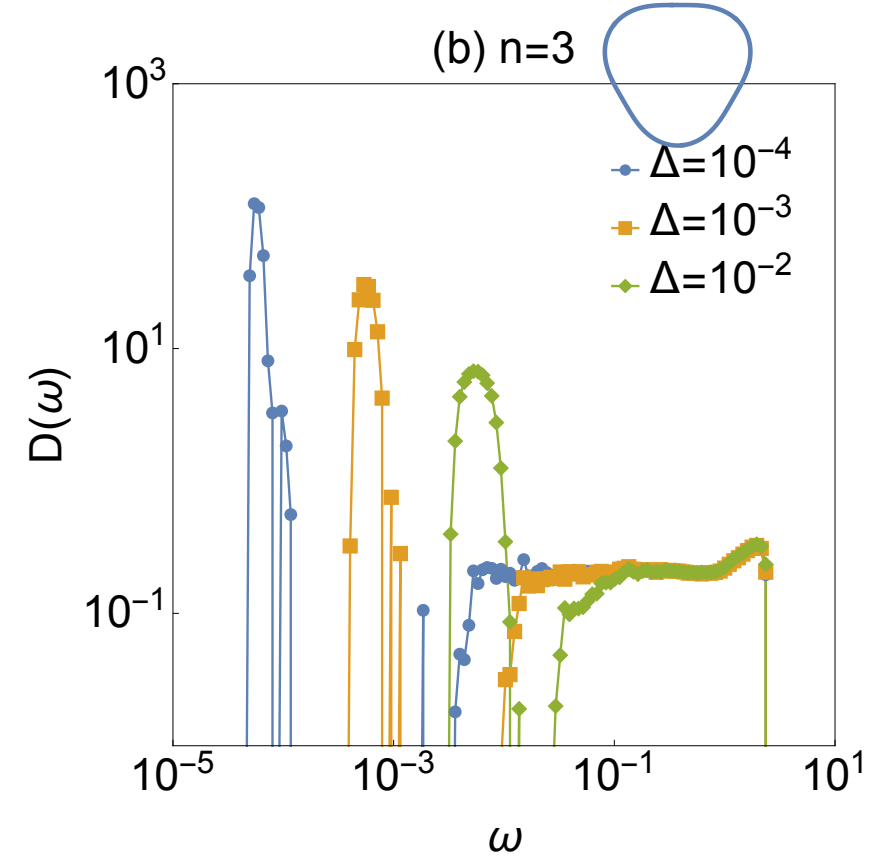
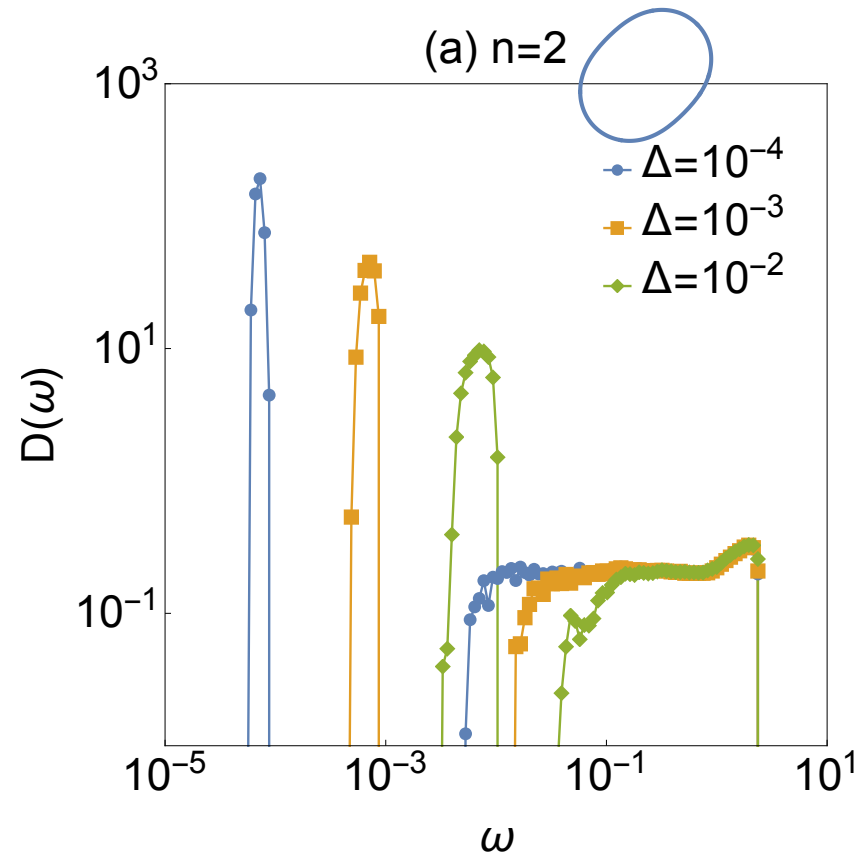
Non-spherical



vDOS of non-spherical particles is gapped!

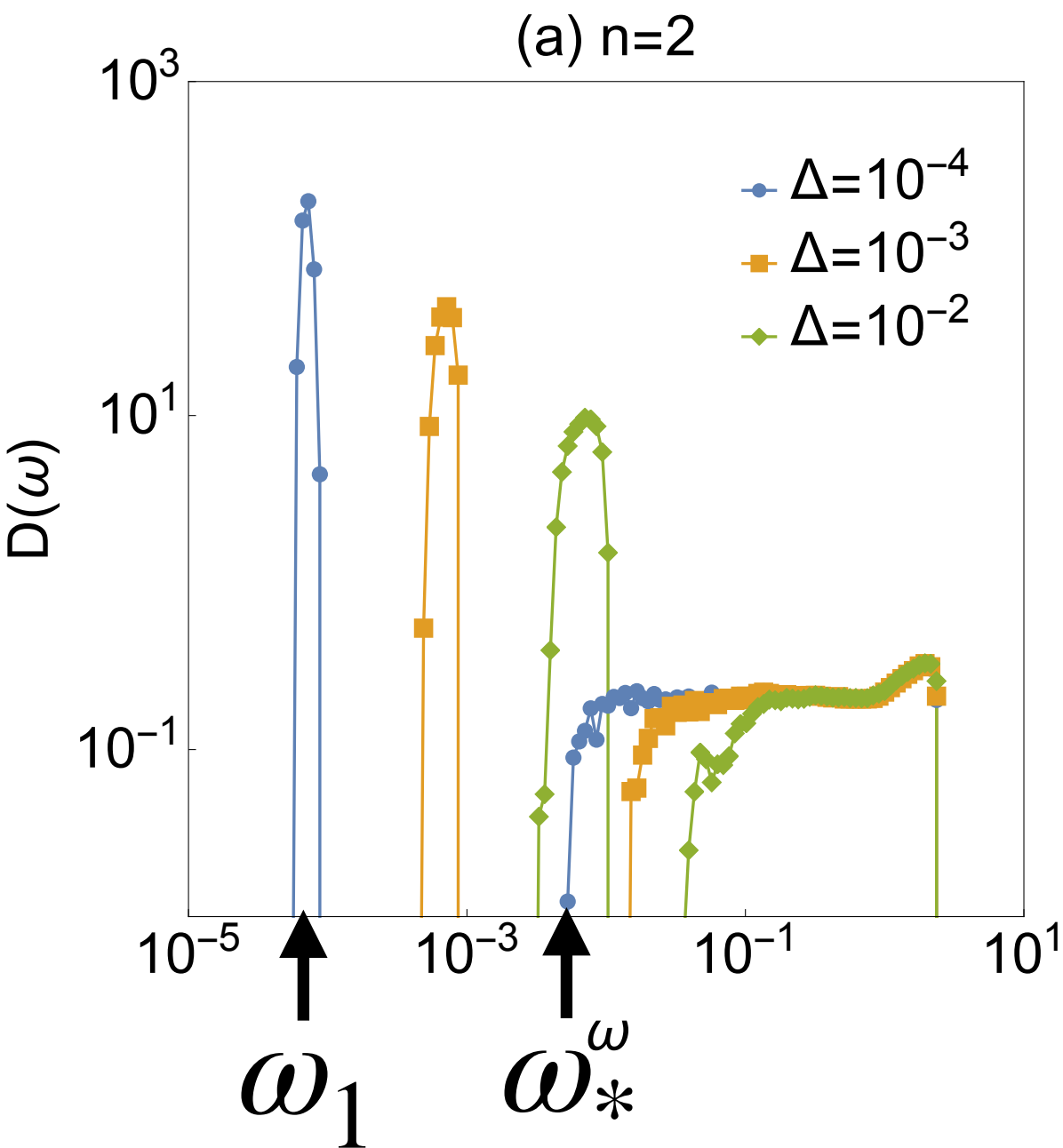
Non-spherical particles

Vibrational density of states

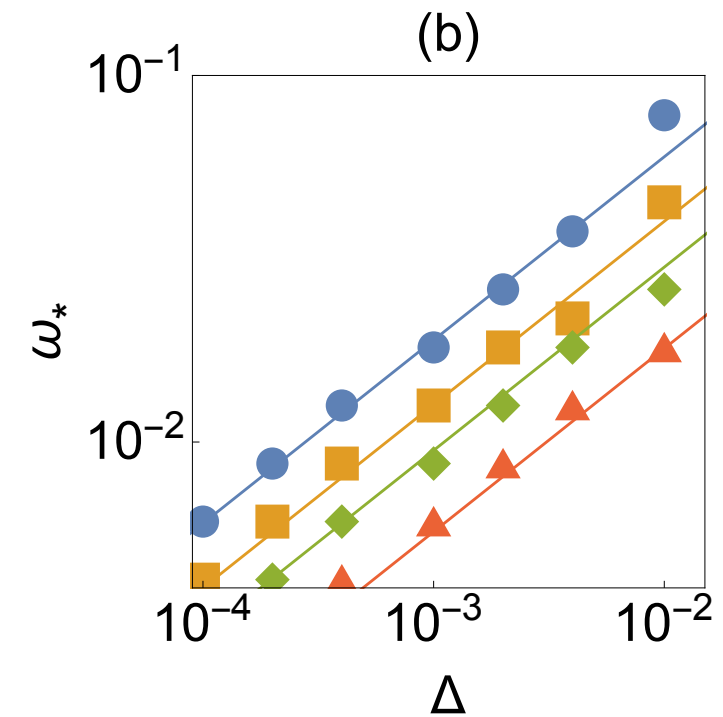
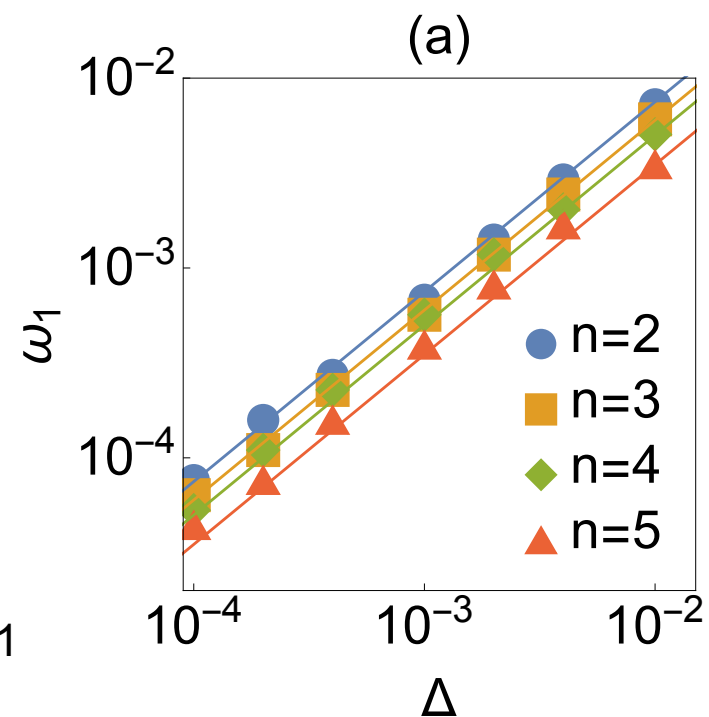


Non-spherical particles

Vibrational density of states



Scalings of characteristic frequencies



Scalings of characteristic frequencies

$$\omega_1 \sim \Delta, \quad \omega_* \sim \Delta^{1/2}$$

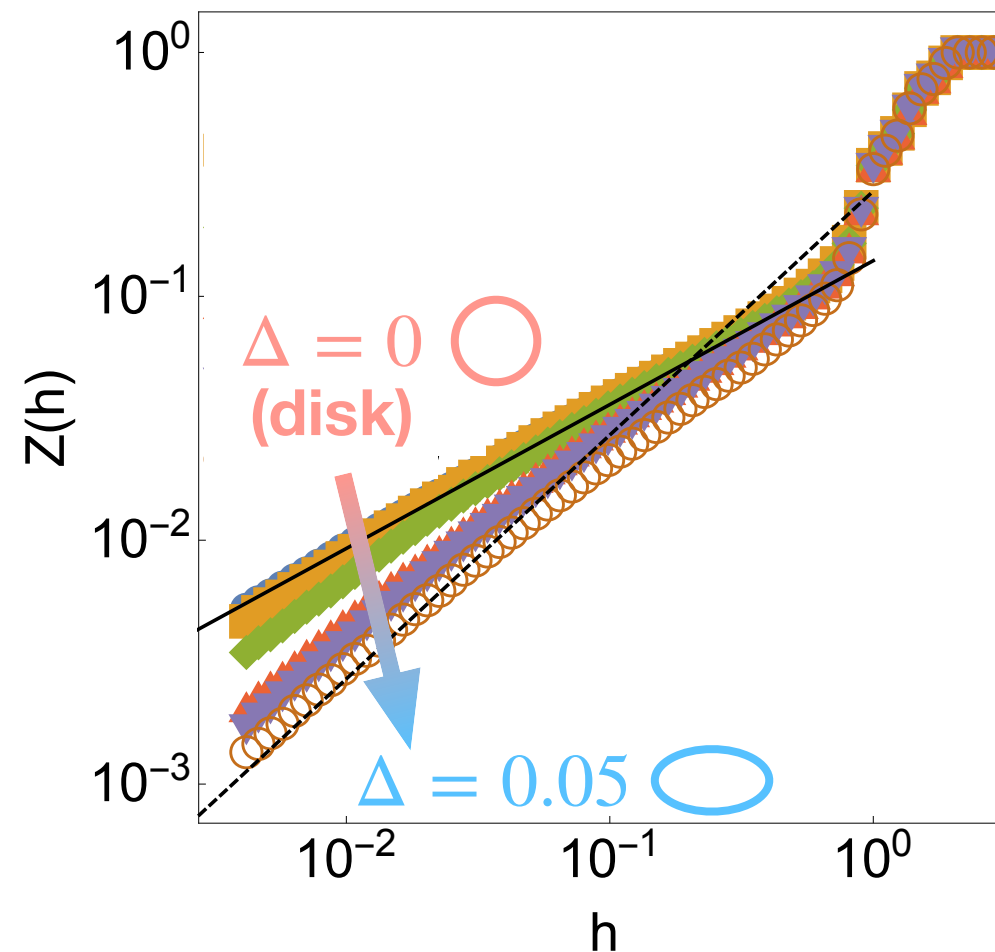
Non-spherical particles

Gap distribution

Cumulative distribution function

$$Z(h) = \int_0^h dh g(h), \quad (g(h) \sim h^{-\gamma} \rightarrow Z(h) \sim h^{1-\gamma})$$

$n=2$

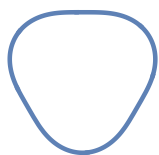


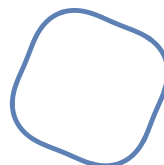
Gap distribution function is finite and regular

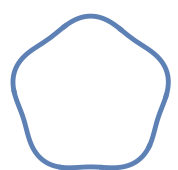
$$Z(h) \sim h \rightarrow g(h) \sim \text{const}.$$

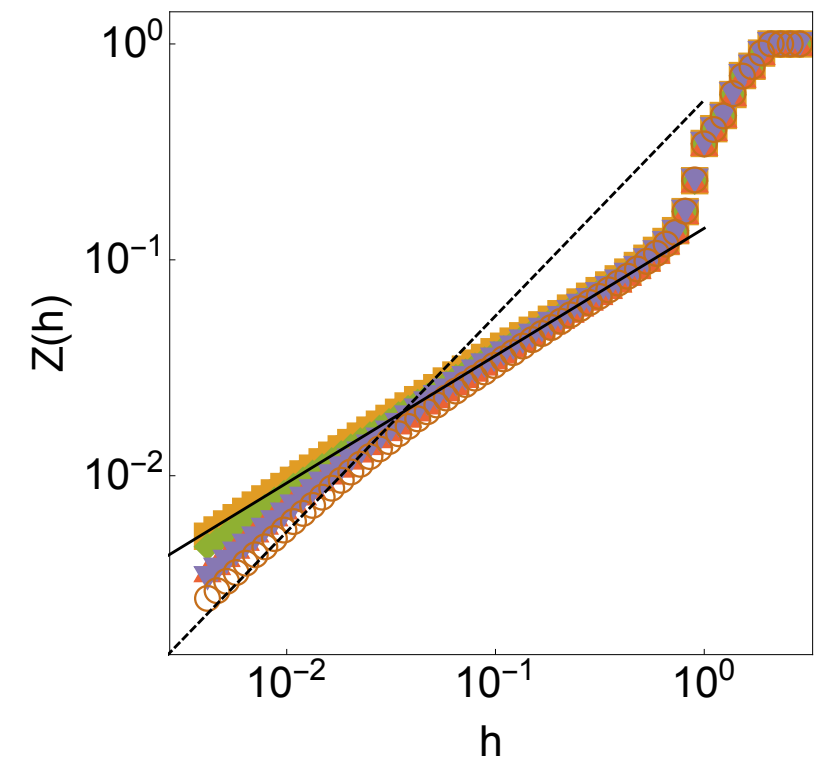
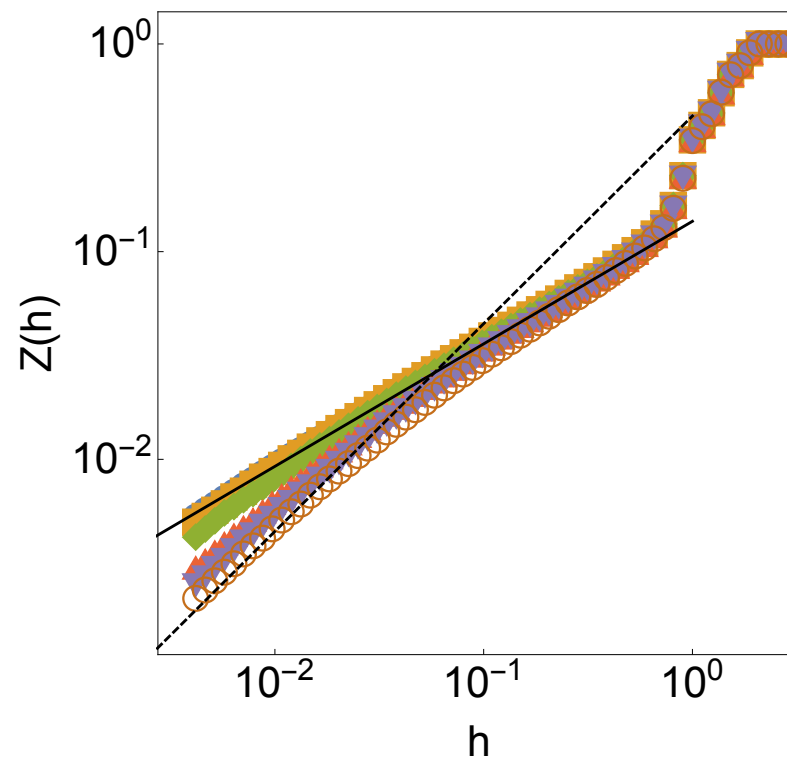
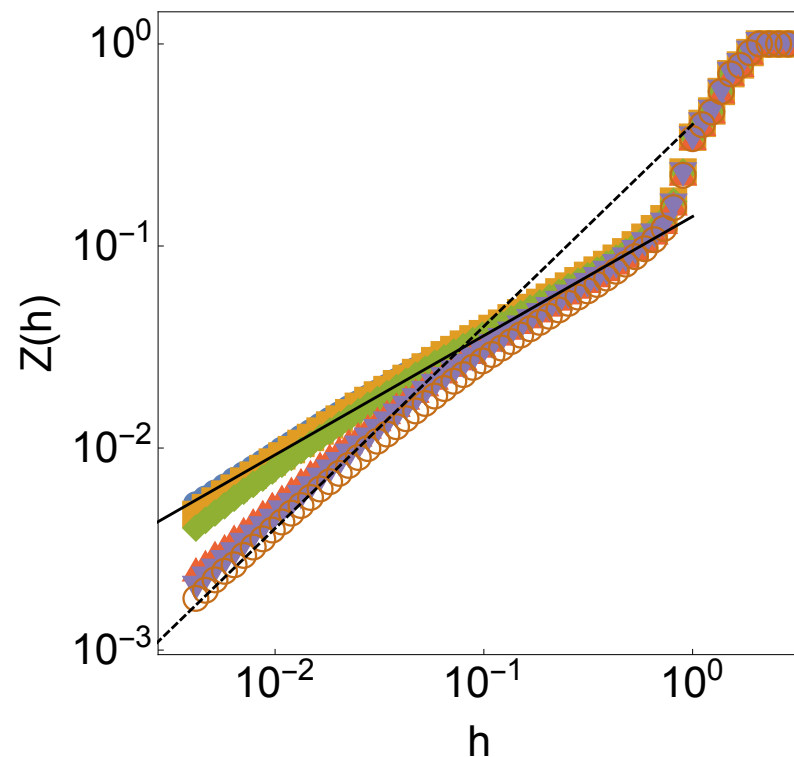
Non-spherical particles

Gap distribution

$n=3$ 

$n=4$ 

$n=5$ 



Gap distribution function is finite and regular

$$Z(h) \sim h \rightarrow g(h) \sim \text{const}.$$

Discussion

Spherical vs non-spherical

Spherical particles
(near jamming)

$$z(\varphi) - 2d \sim \delta\varphi^{1/2}$$

$$\omega_* \sim \delta\varphi^{1/2}$$

$$g(h) \sim \begin{cases} \delta\varphi^{-\mu\gamma} (h \ll \delta\varphi^\mu) \\ h^{-\gamma} & (h \gg \delta\varphi^\mu) \end{cases}$$

Non-spherical
(at jamming)

$$z_J - 2d \sim \Delta^{1/2}$$

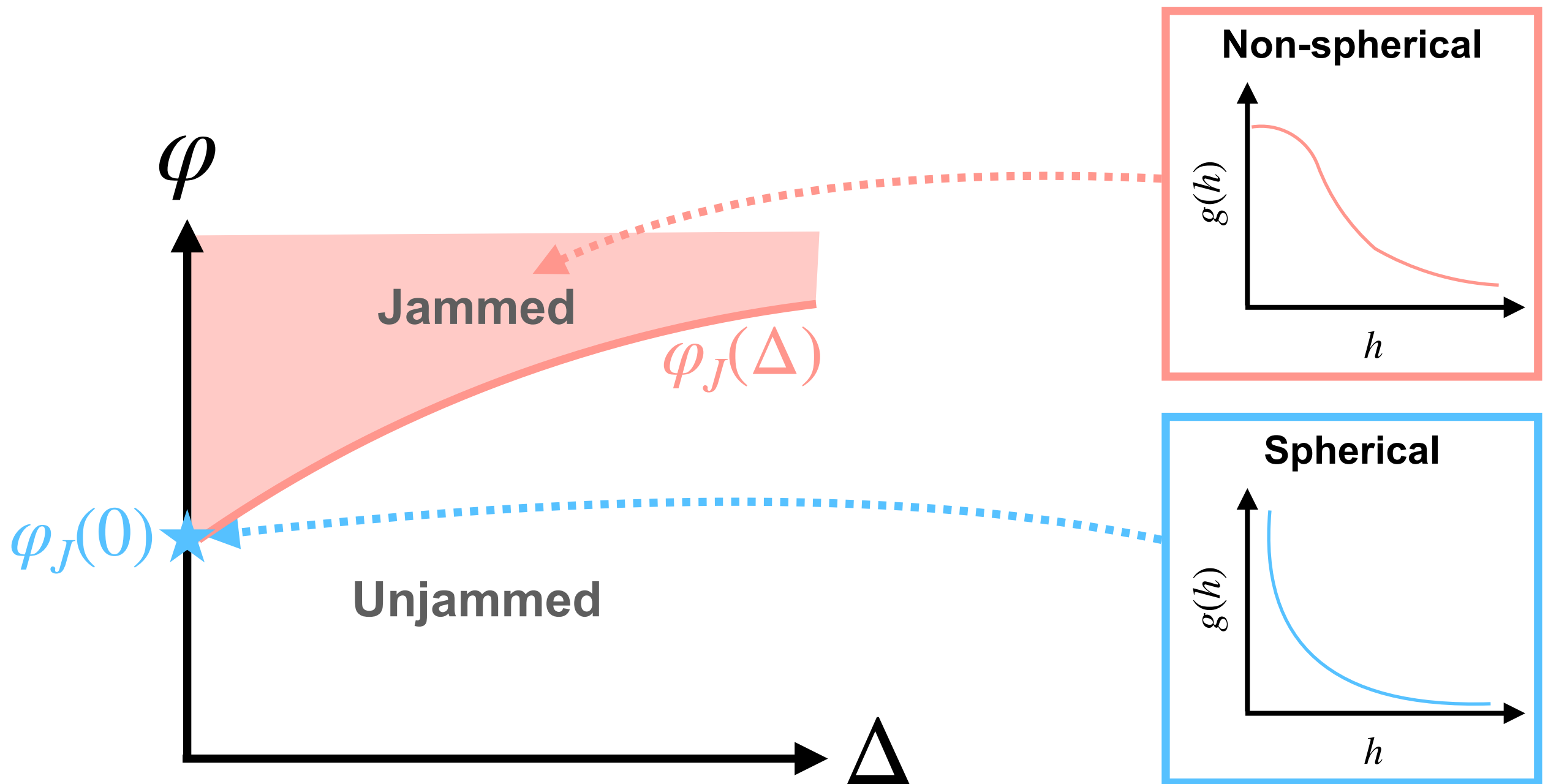
$$\omega_* \sim \Delta^{1/2}$$

$$g(h) \sim \begin{cases} \Delta^{-\mu\gamma} (h \ll \Delta^\mu) \\ h^{-\gamma} & (h \gg \Delta^\mu) \end{cases}$$

Δ plays the same role as $\delta\varphi$!

Discussion

Phase diagram



The jamming transition point of non-spherical particles is not critical in terms of $g(r)$!!!

Discussion

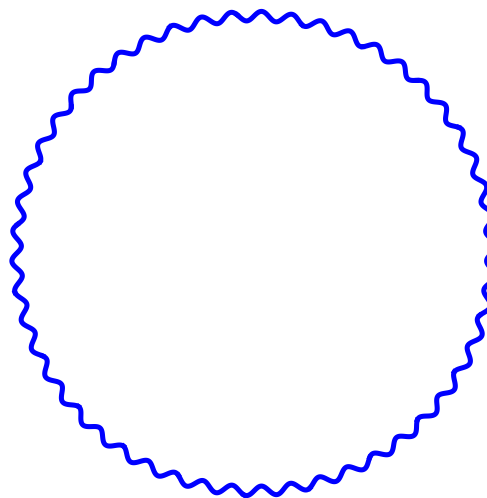
Summary and Future works

Summary

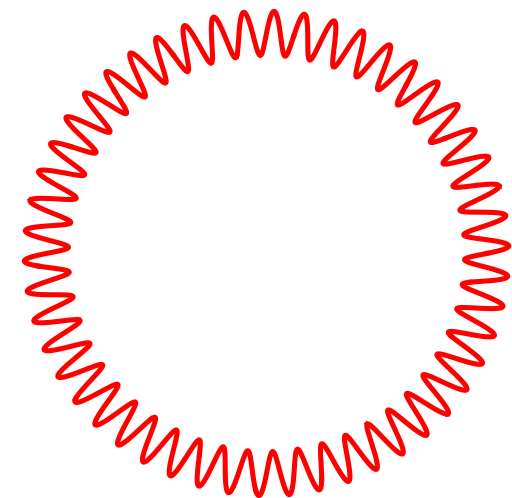
Deformation = compression

Future works

- Does the correlation length diverge at the jamming transition point of non-spherical particles?
- Investigate scalings above jamming
cf: ellipsoids $G \sim (\varphi - \varphi_J)^1$ $z - z_J \sim (\varphi - \varphi_J)^1$ M. Mailman *et al* (2009), C. F. Schreck *et al.* (2012)
- Investigate scalings below jamming (rheology)
- Geometric friction
H. Ikeda *et al.* PRL **124**, 208001 (2020)



Weak friction

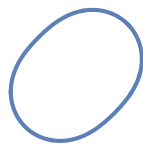


Strong friction

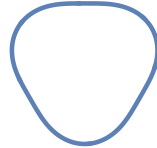
Non-spherical particles

Gap distribution

$n=2$



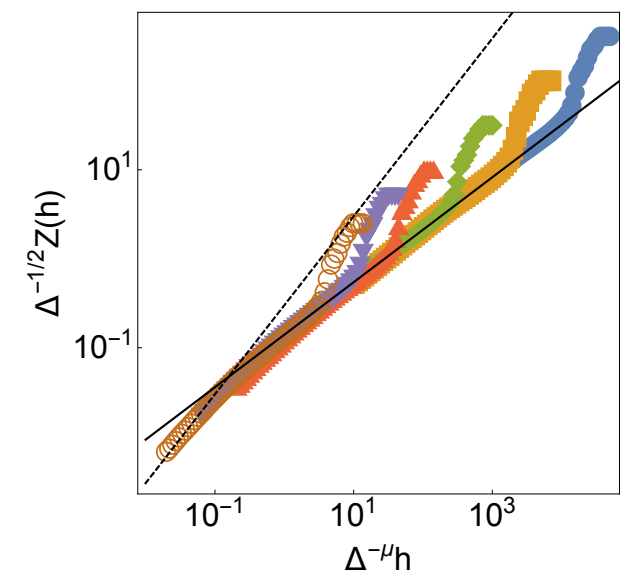
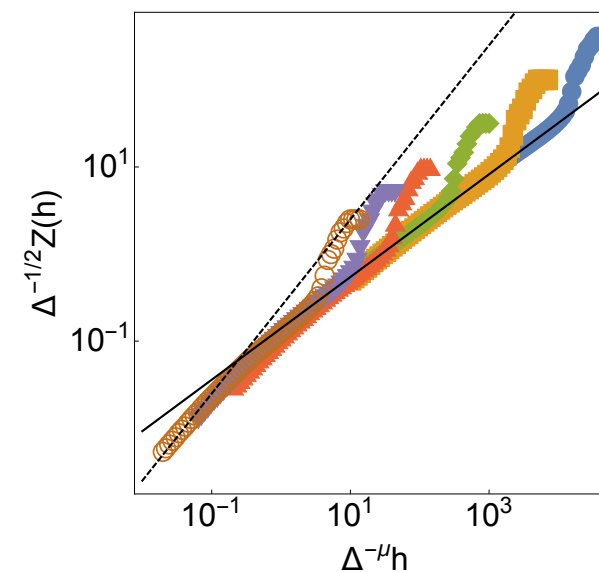
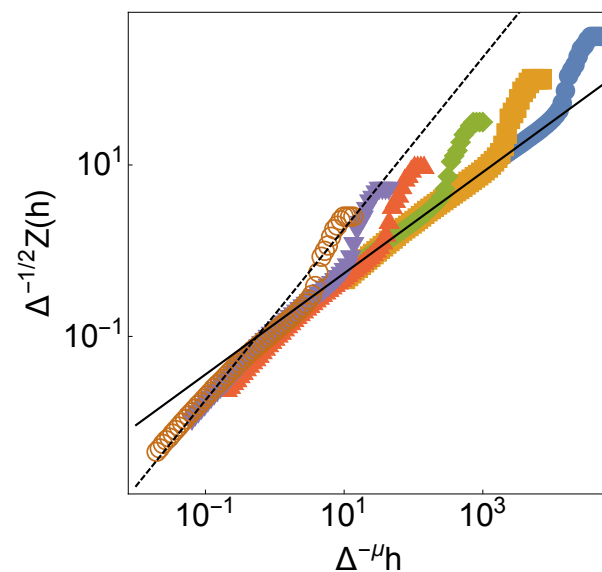
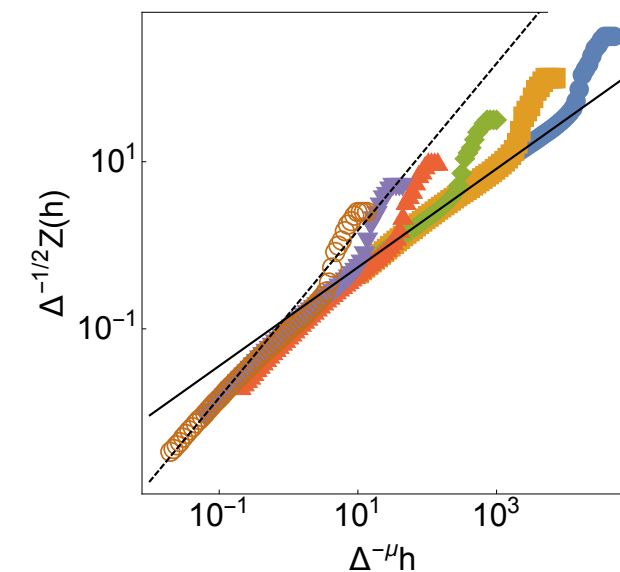
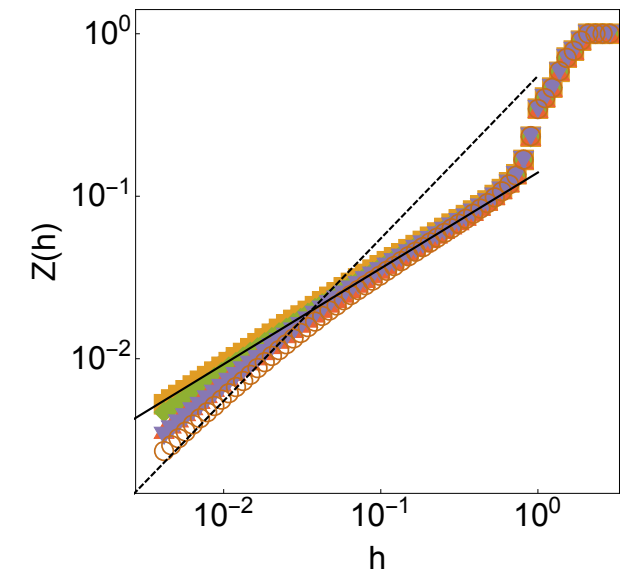
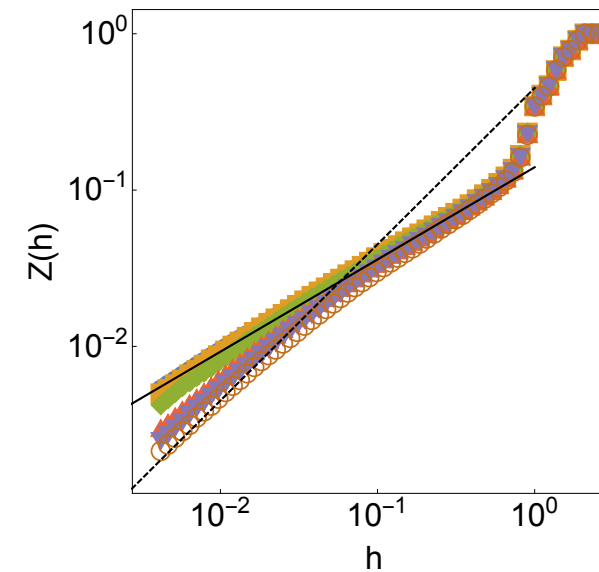
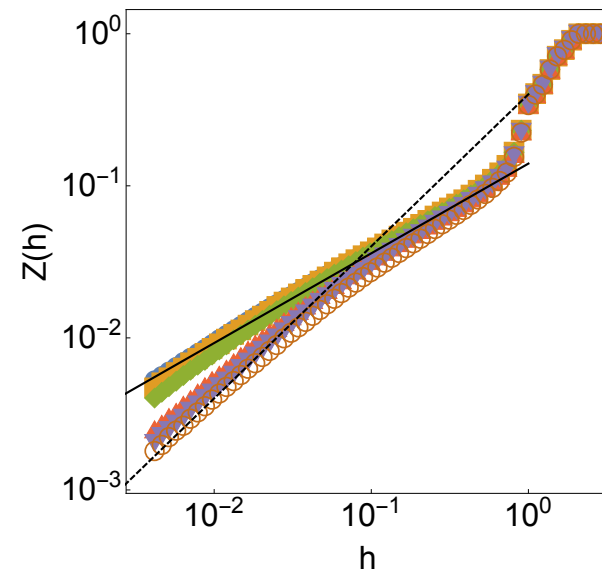
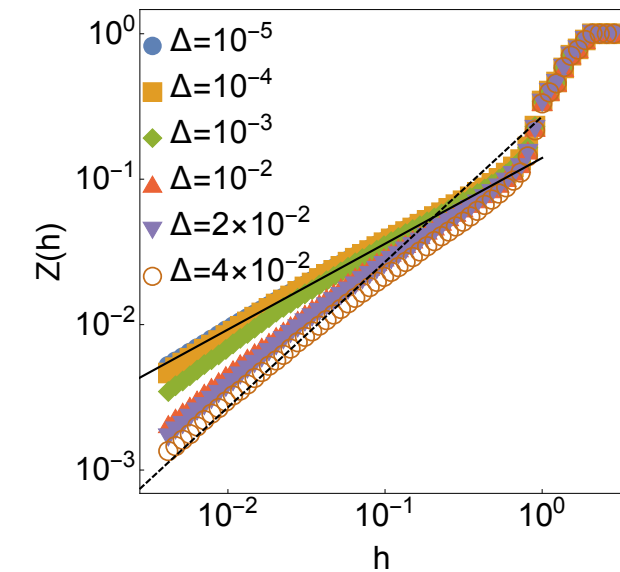
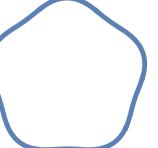
$n=3$



$n=4$



$n=5$



Gap distribution function is finite and regular

$$Z(h) \sim h \rightarrow g(h) \sim \text{const}.$$

Discussion

Spherical vs non-spherical

Spherical particles
(near jamming)

$$z(\varphi) - 2d \sim \delta\varphi^{1/2}$$

$$g(h) \sim \begin{cases} \delta\varphi^{-\mu\gamma} & (h \ll \delta\varphi^\mu) \\ h^{-\gamma} & (h \gg \delta\varphi^\mu) \end{cases}$$

Non-spherical
(at jamming)

$$z_J - 2d \sim \Delta^{1/2}$$

$$g(h) \sim \begin{cases} \Delta^{-\mu\gamma} & (h \ll \Delta^\mu) \\ h^{-\gamma} & (h \gg \Delta^\mu) \end{cases}$$

Both the distance from the jamming transition point and deviation from spheres controls the scaling behaviors.

Non-spherical particles

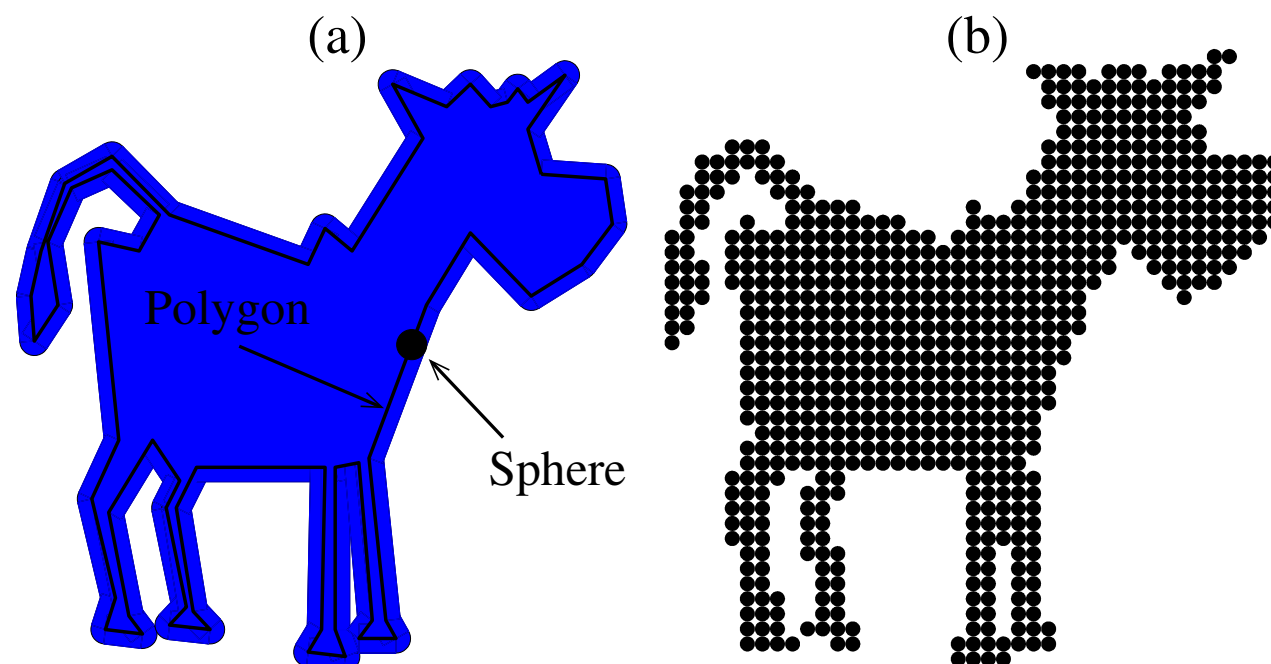
Model

Harmonic potential

$$V_N = \sum_{i < j} \frac{h_{ij}^2}{2} \theta(-h_{ij})$$

Gap function: the minimal distance between two particles

For general shapes of non-spherical particles, the calculation of the gap function is not so easy.



F. Alonso-Marroquin (2008)

Non-spherical particles

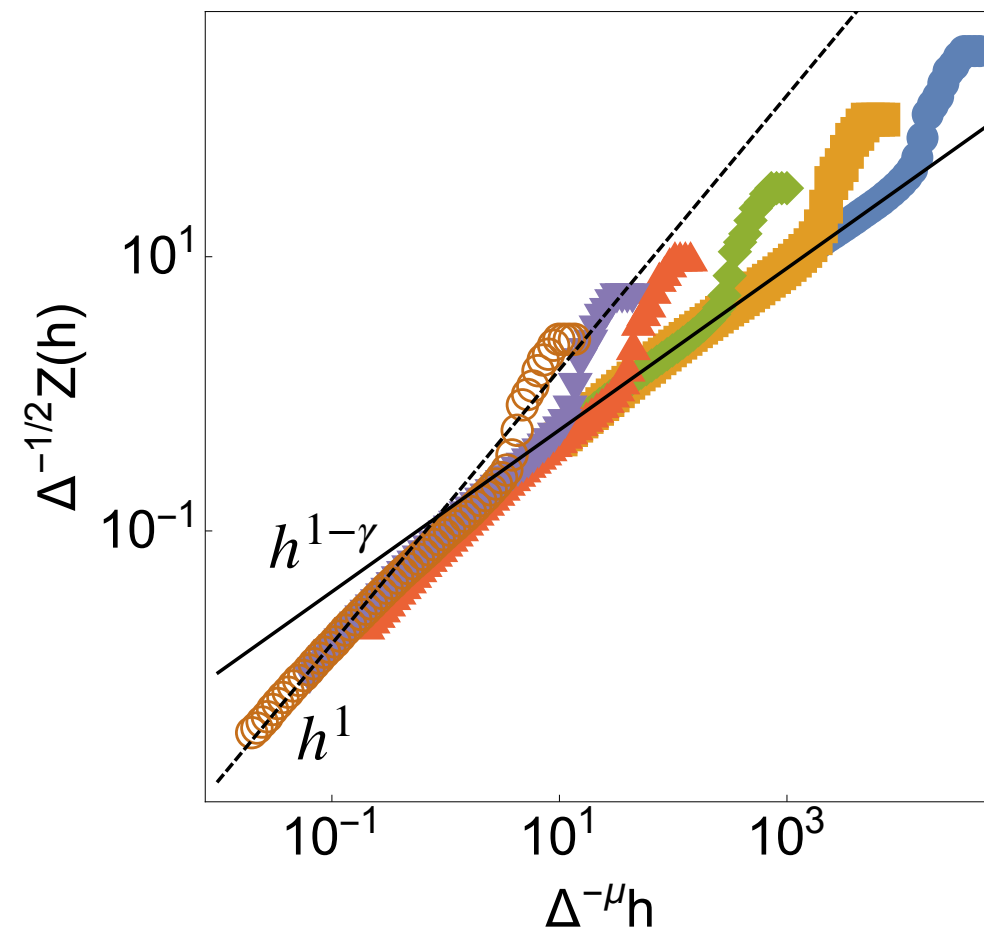
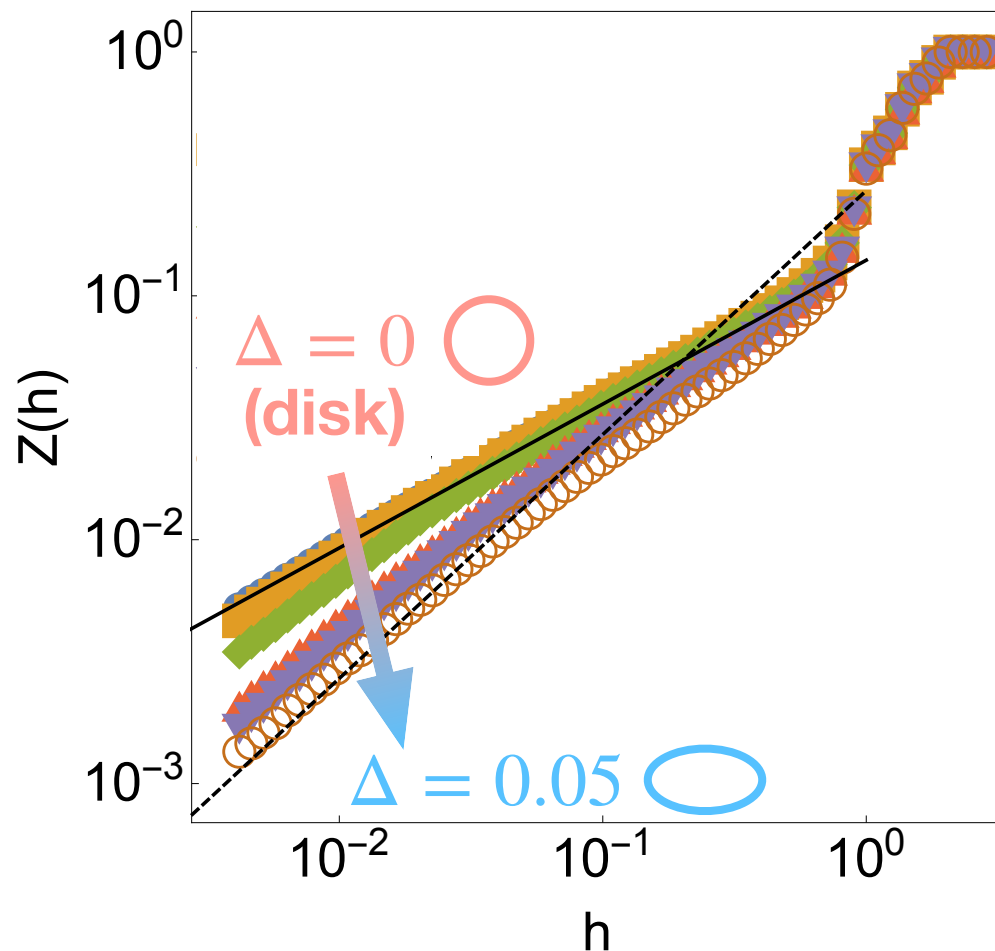
Gap distribution

Cumulative distribution function

$$Z(h) = \int_0^h dh g(h), \quad (g(h) \sim h^{-\gamma} \rightarrow Z(h) \sim h^{1-\gamma})$$

$n=2$

$n=2$ (scaling plot)

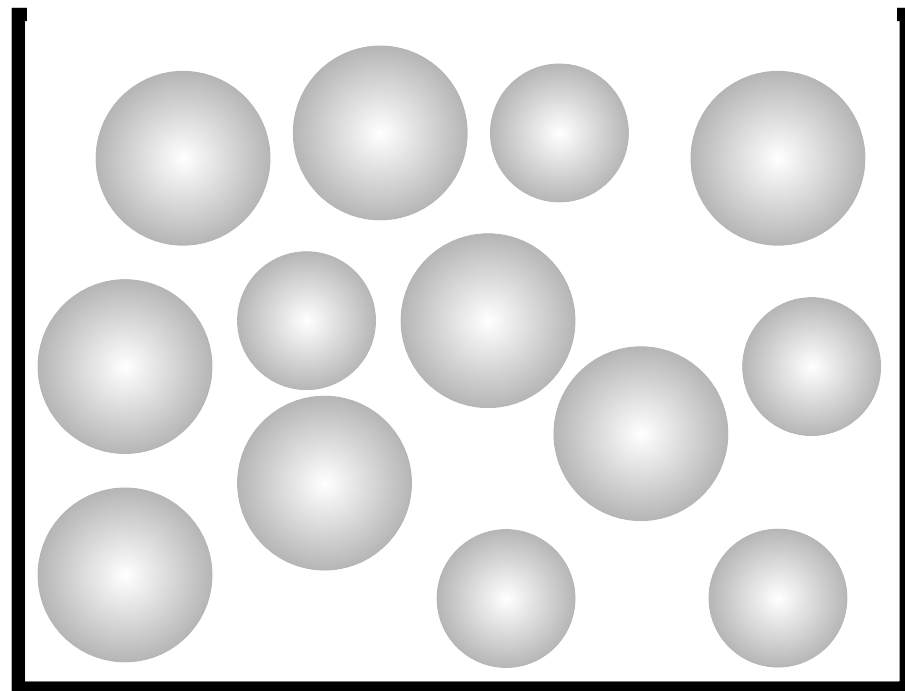


Gap distribution function is finite and regular

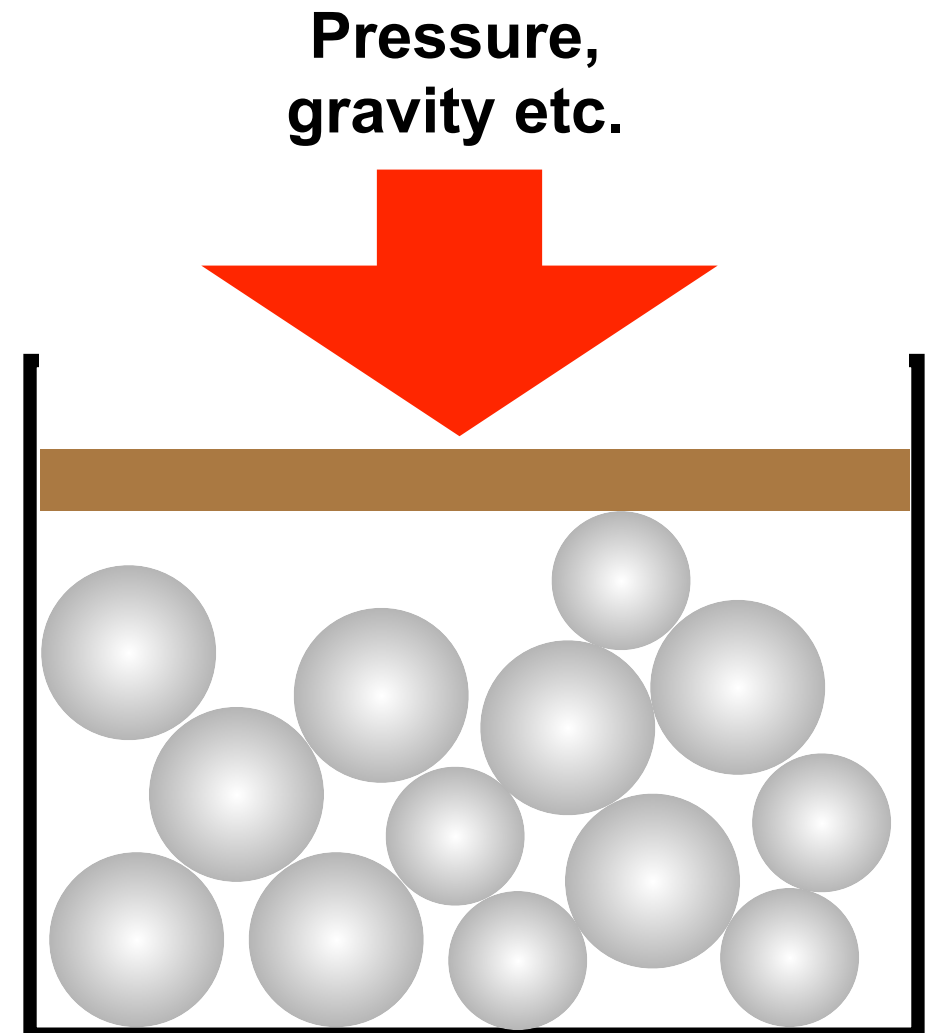
$$Z(h) \sim h \rightarrow g(h) \sim \text{const}.$$

Introduction

What is the jamming transition?



fluid like



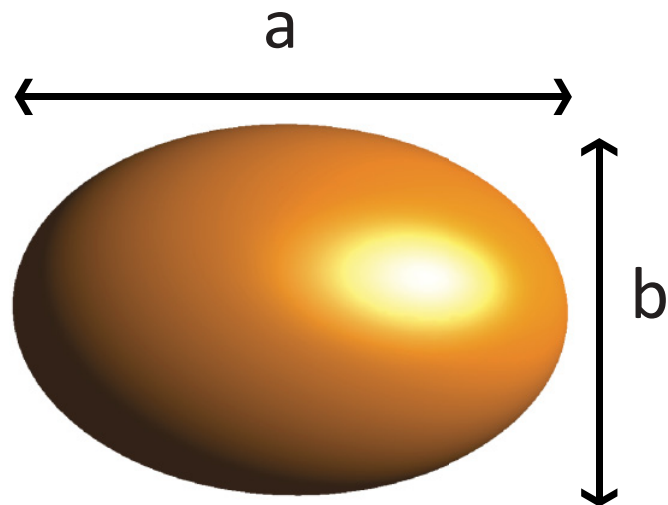
solid like

Jamming transition = fluid-solid phase transition
at zero temperature

Non-spherical particles

Previous numerical studies

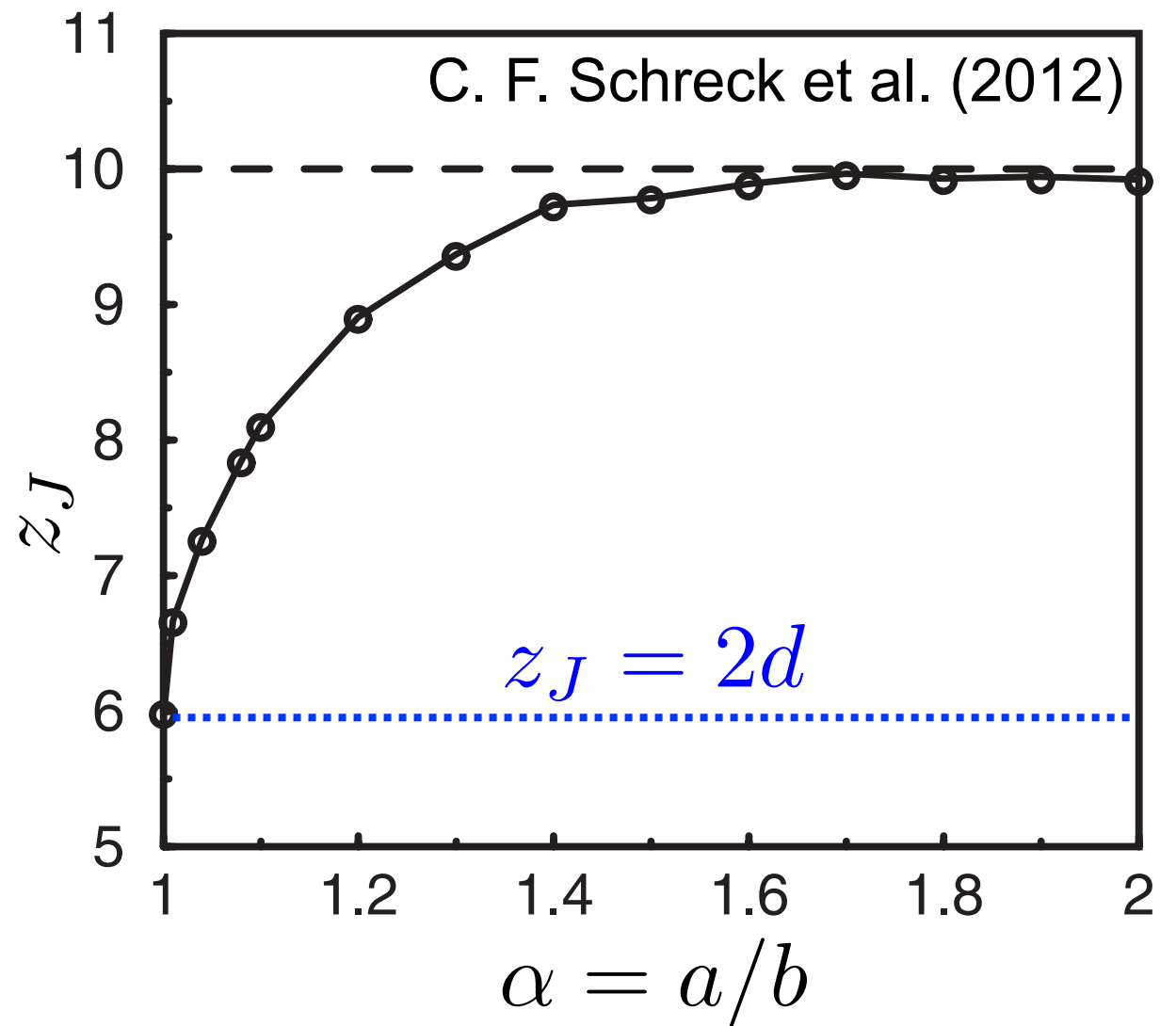
Ellipsoid



Control parameter

$$\alpha = a/b$$

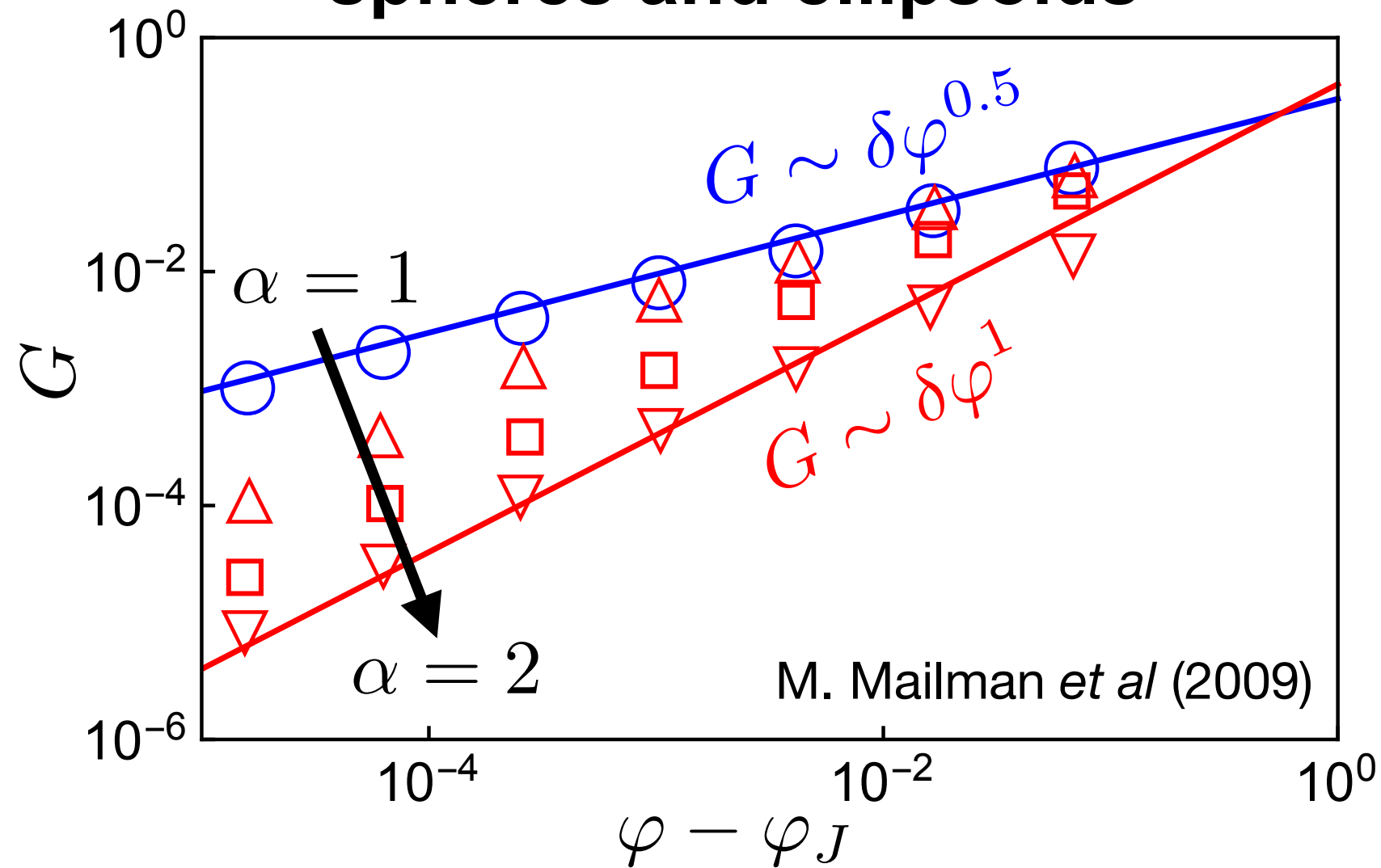
Contact number at the jamming point



Non-spherical particles

Previous numerical studies

Shear modulus of spheres and ellipsoids

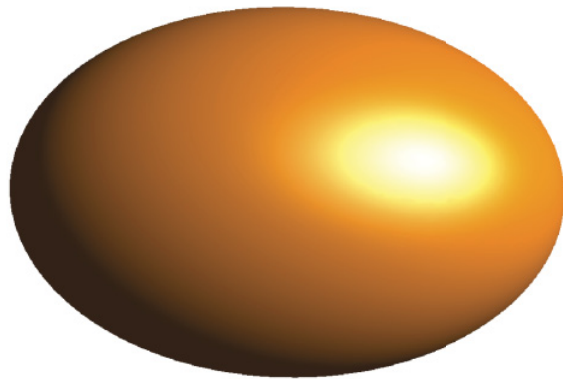


Spheres : $G \sim \delta\varphi^{1/2}$ Ellipsoids : $G \sim \delta\varphi^1$

Non-spherical particles

Previous numerical studies

Ellipsoid



M. Mailman *et al* (2009)

$$G \sim (\varphi - \varphi_J)^1$$

C. F. Schreck *et al.* (2012)

$$z - z_J \sim (\varphi - \varphi_J)^1$$

$$g(r) = ???$$

Other shapes

???

Purpose of this study

Systematic investigation of various shapers of particles

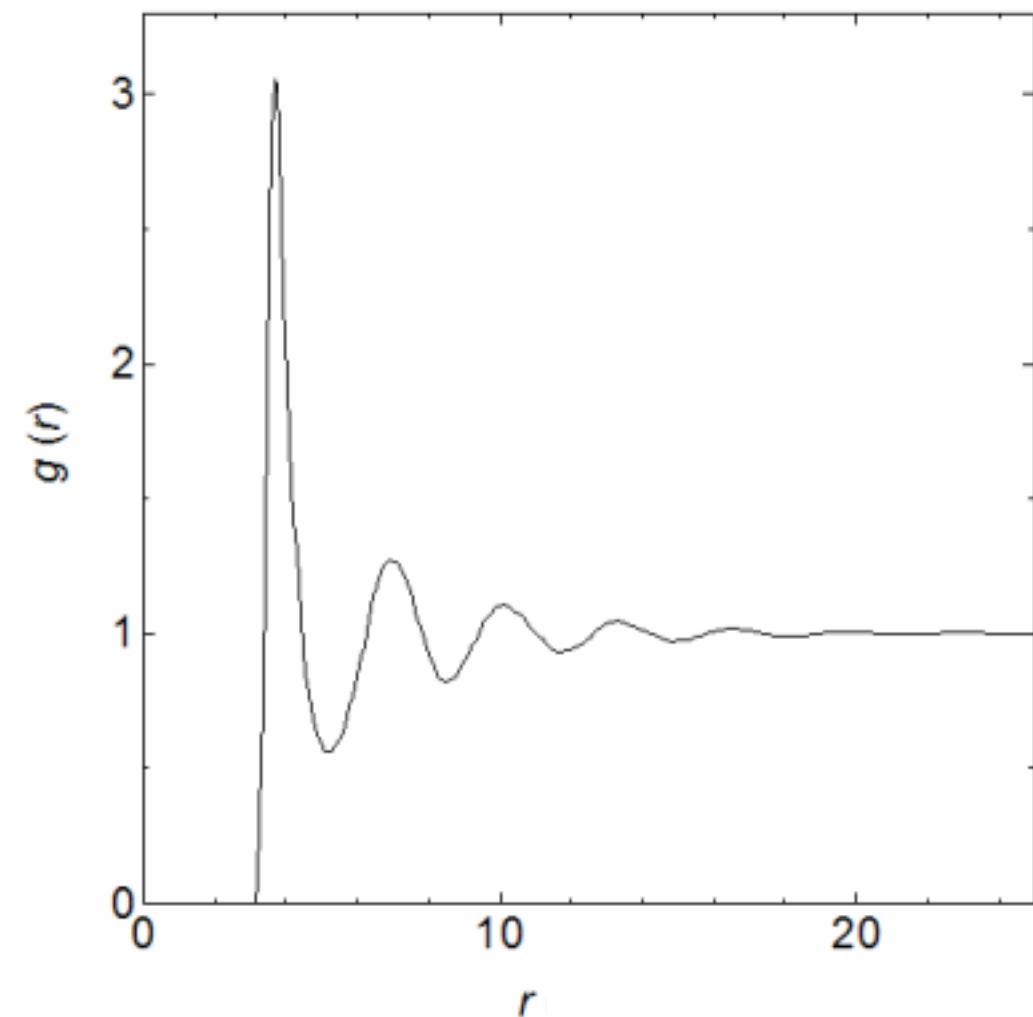
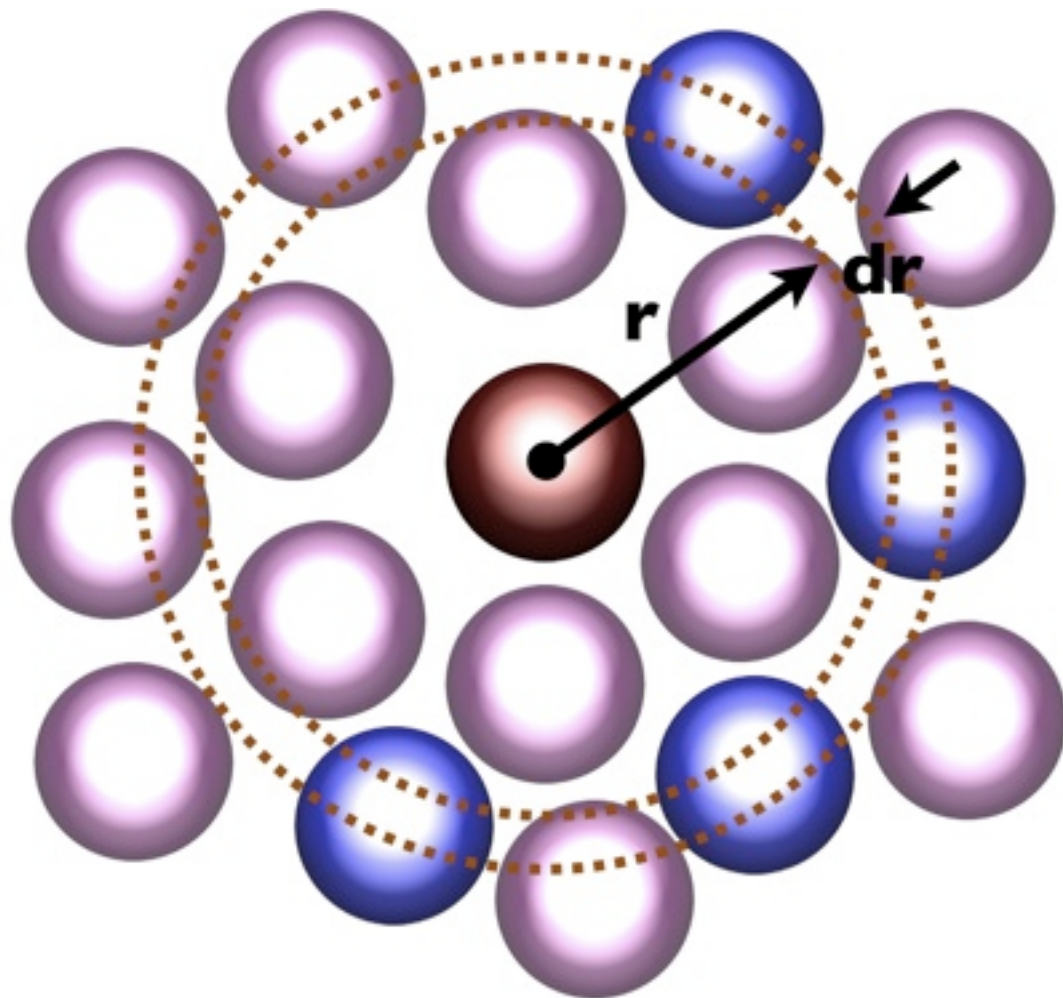
Spherical particles

Power-law

Two point correlation function

$$g(r) = \frac{1}{N} \sum_{i \neq j} \delta(r - |\vec{x}_i - \vec{x}_j|)$$

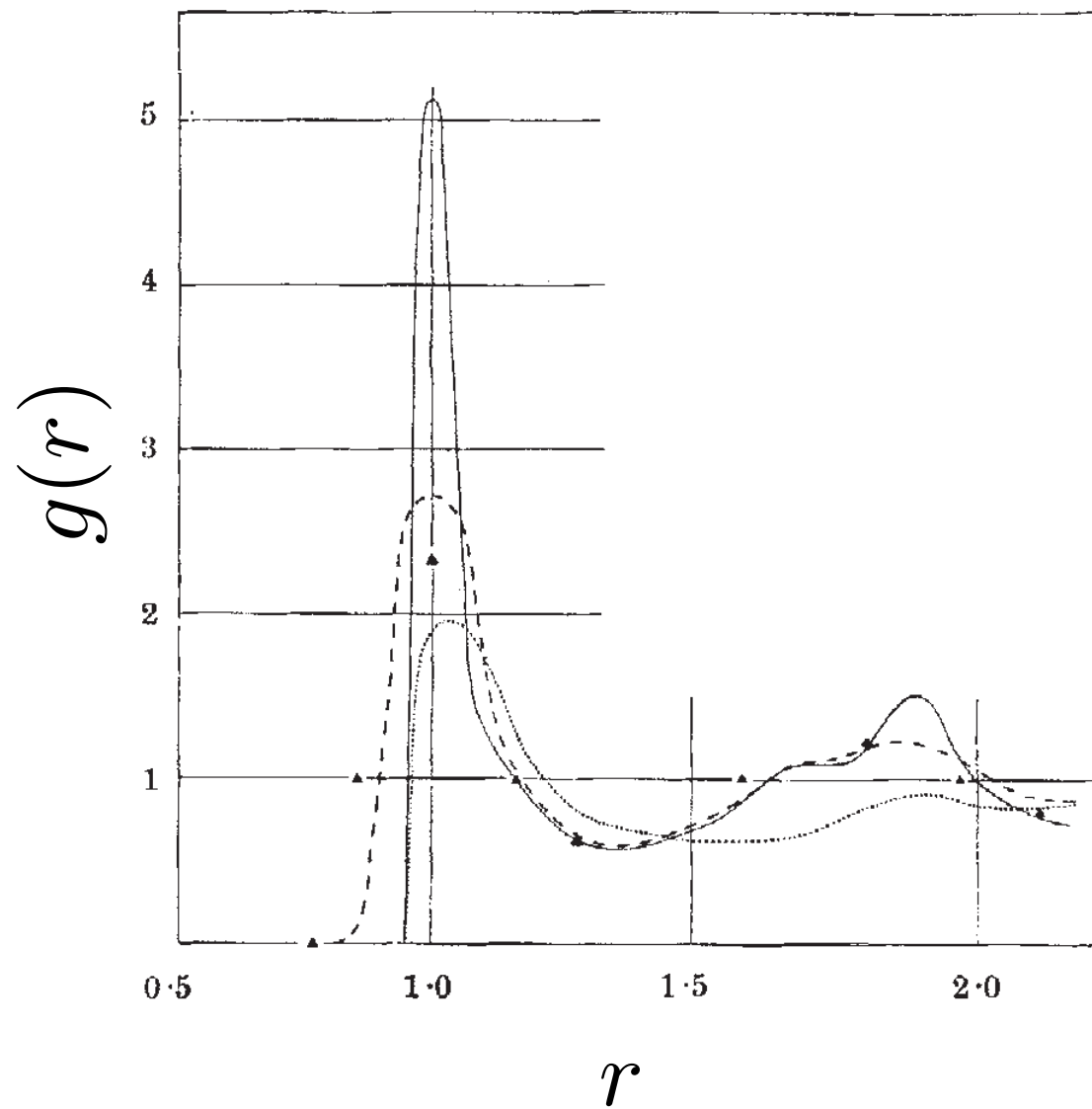
$$\int dr 4\pi r^2 \rho g(r) = N - 1$$



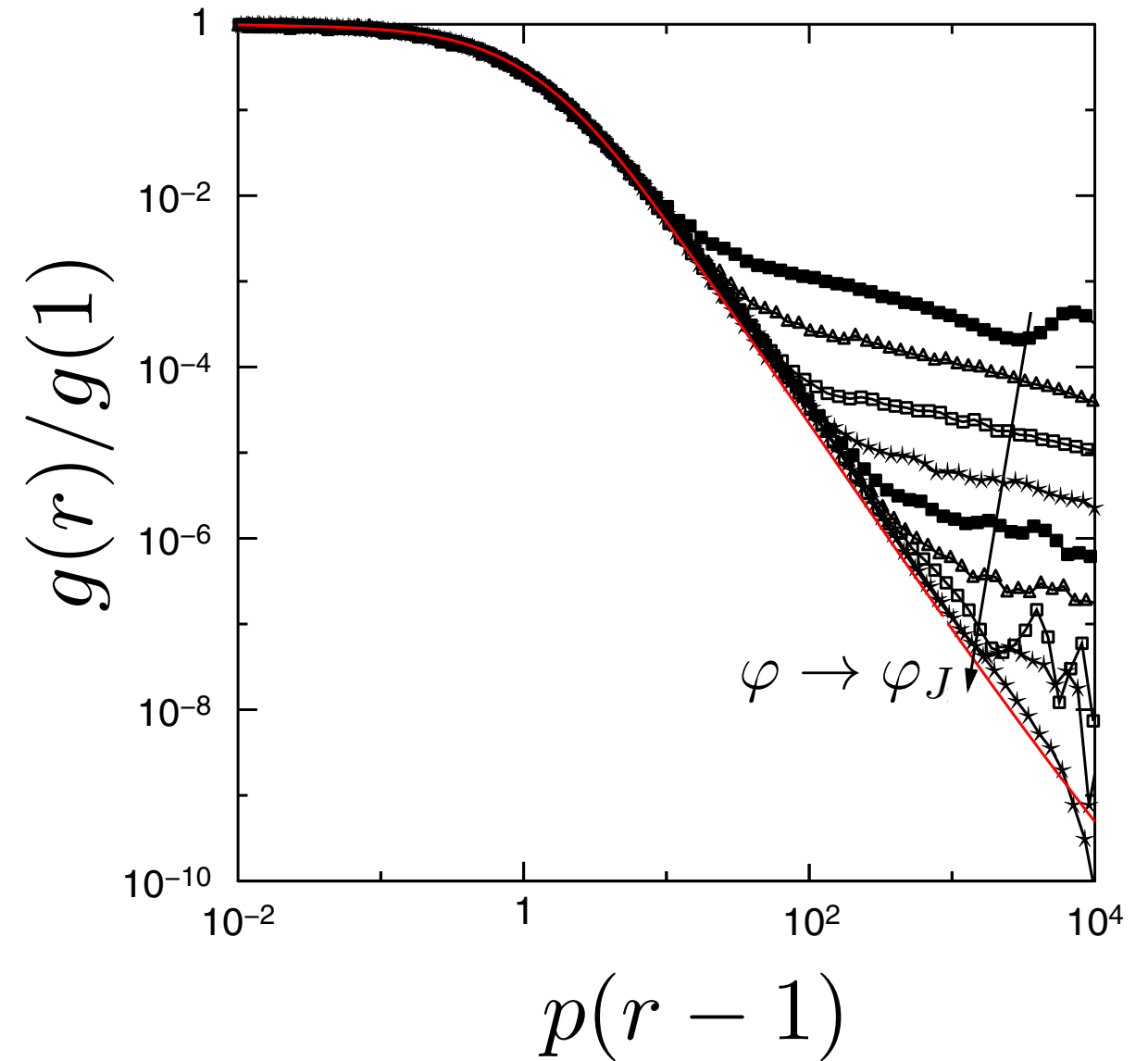
Spherical particles

Power-law

J. D. Bernal (1960)



P. Charbonneau *et al.* (2014)



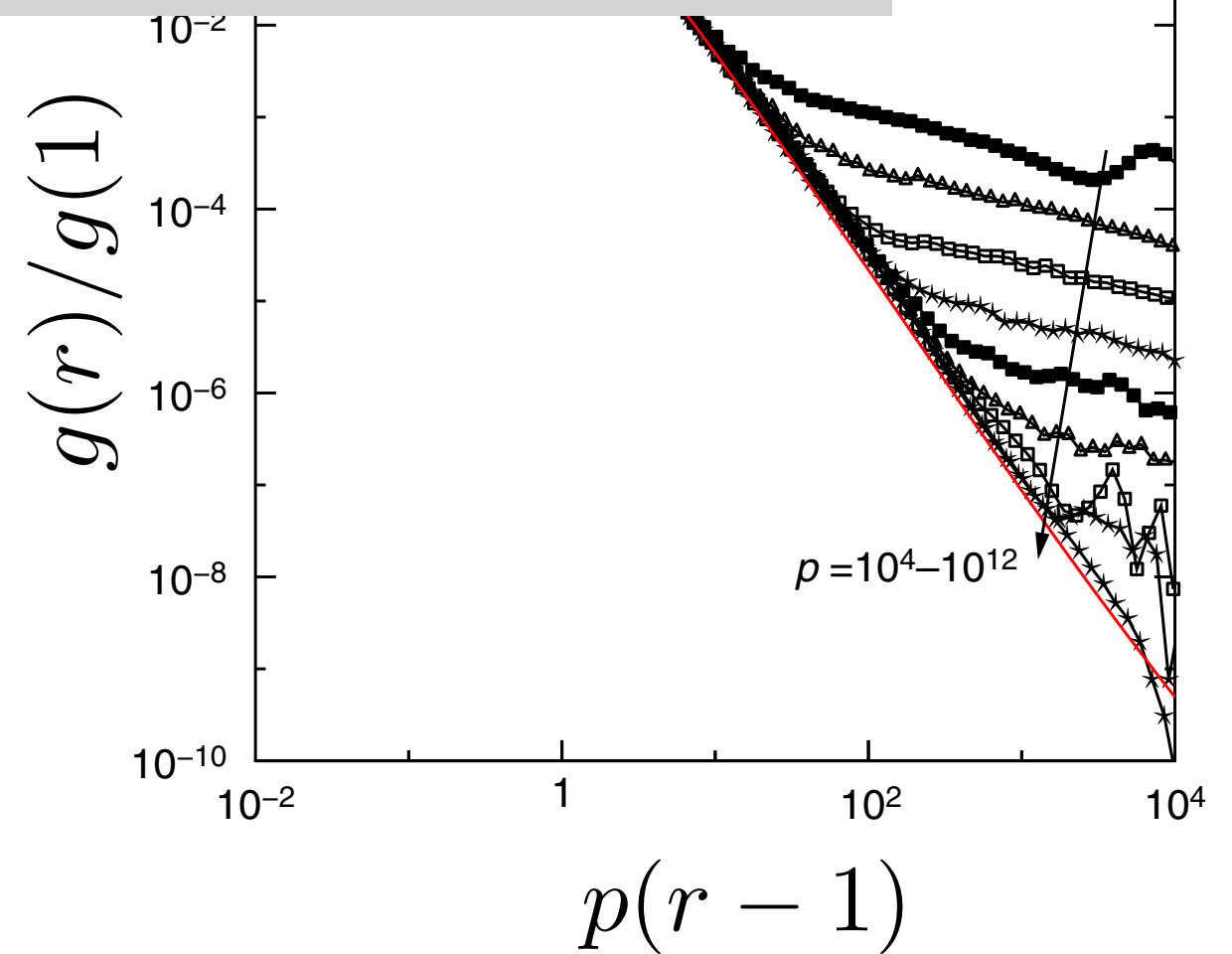
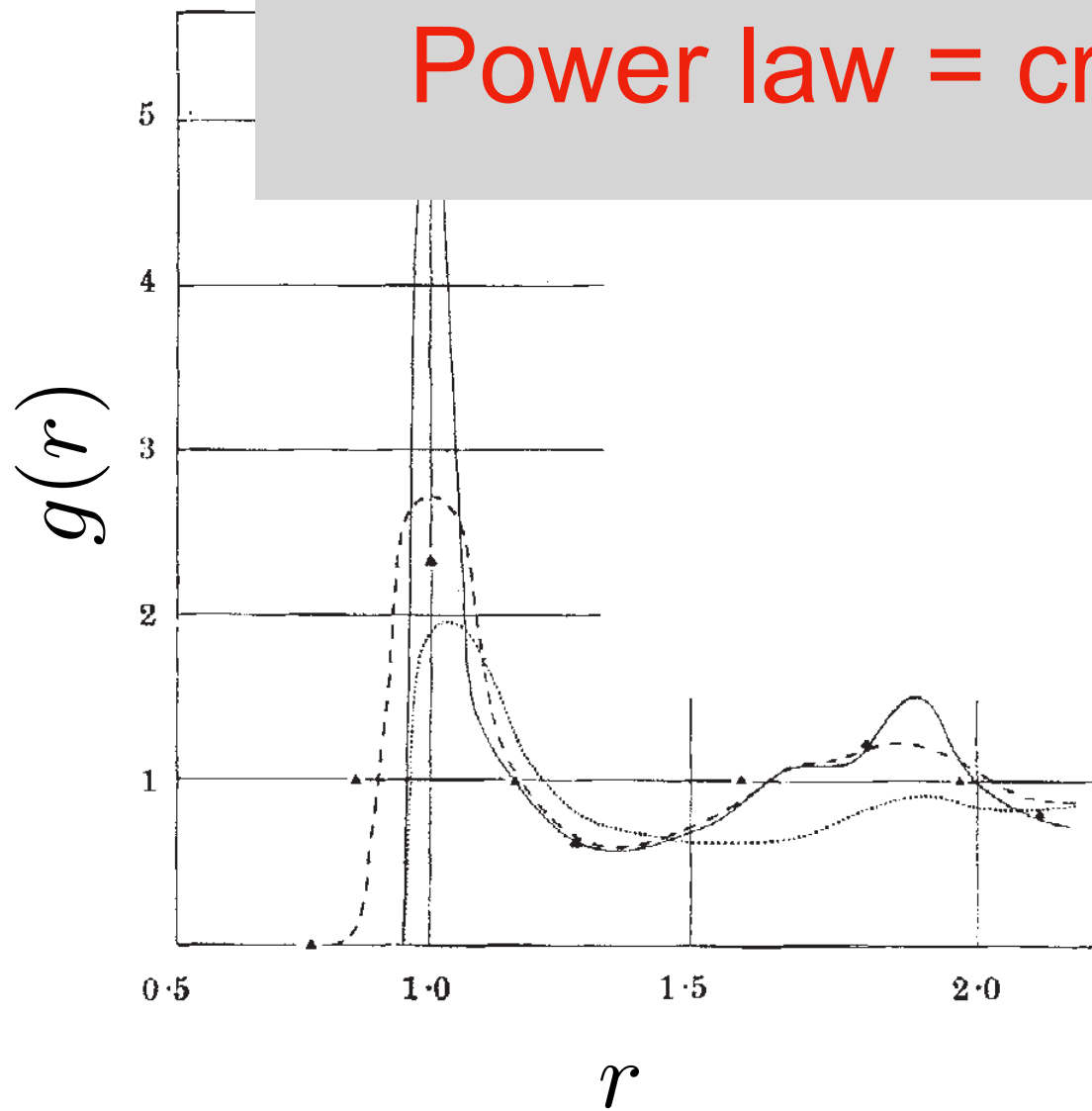
Spherical particles

Power-law

L. D. D. (1999)

D. G. et al. (2014)

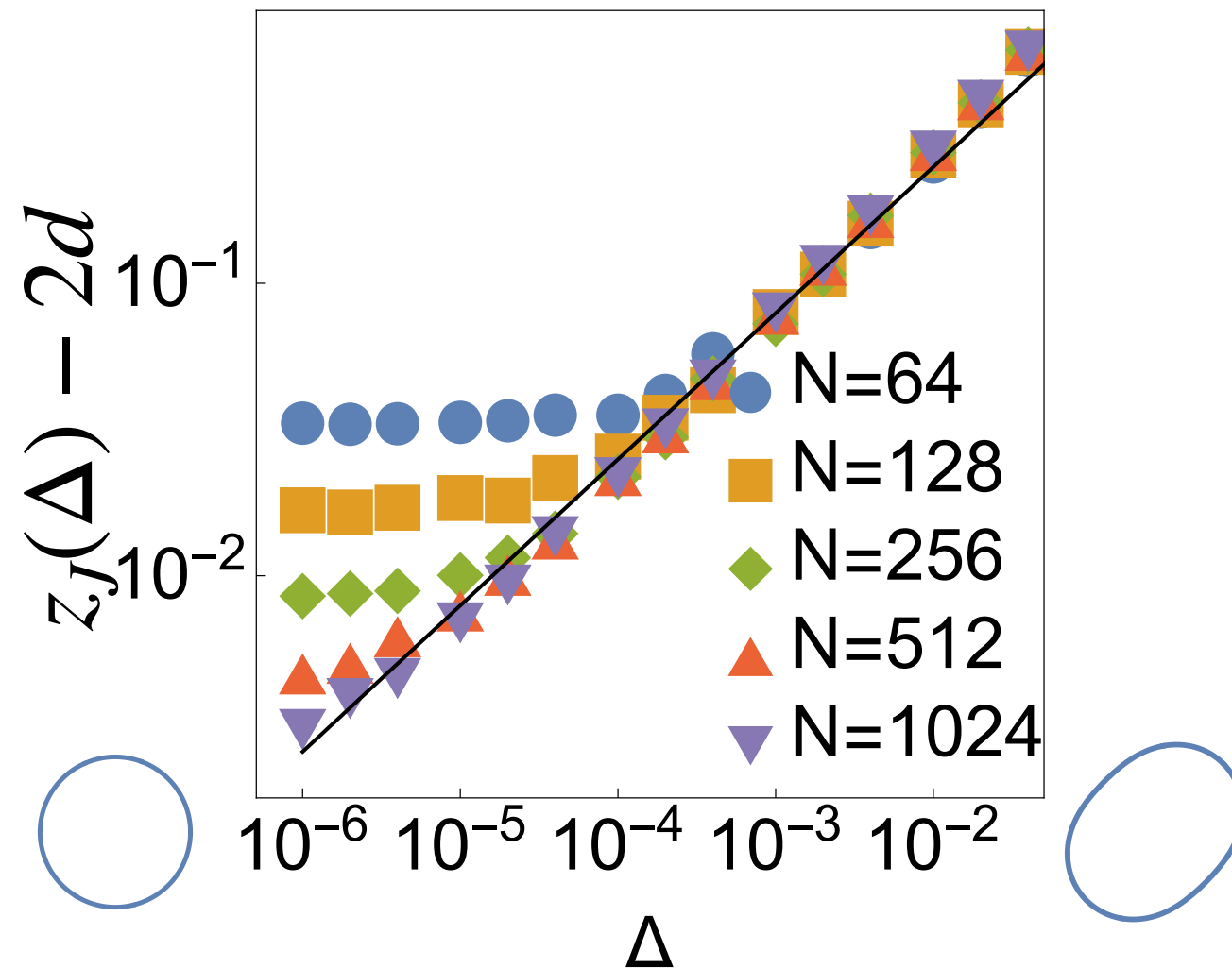
Power law = critical phenomena!!!



Non-spherical particles

Contact number

Contact number for $n=2$



The contact number increases with the aspect ratio as

$$z_J(\Delta) - z_J(0) \sim \Delta^{1/2}$$

Spherical particles

Motivation

To simplify the problem, we consider spherical particles



wikipedia.org/wiki/Spherical_cow